

Study of the $e^+e^- \rightarrow K^+K^-$ reaction in the energy range from 2.6 to 8.0 GeV

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The $e^+e^- \rightarrow K^+K^-$ cross section and charged-kaon electromagnetic form factor are measured in the e^+e^- center-of-mass energy range (E) from 2.6 to 8.0 GeV using the initial-state radiation technique with an undetected photon. The study is performed using 469 fb^{-1} of data collected with the BABAR detector at the PEP-II e^+e^- collider at center-of-mass energies near 10.6 GeV. The form factor is found to decrease with energy faster than $1/E^2$ and approaches the asymptotic QCD prediction. Production of the K^+K^- final state through the J/ψ and $\psi(2S)$ intermediate states is observed. The results for the kaon form factor are used together with data from other experiments to perform a model-independent determination of the relative phases between electromagnetic (single-photon) and strong amplitudes in J/ψ and $\psi(2S) \rightarrow K^+K^-$ decays. The values of the branching fractions measured in the reaction $e^+e^- \rightarrow K^+K^-$ are shifted relative to their true values due to interference between resonant and nonresonant amplitudes. The values of these shifts are determined to be about $\pm 5\%$ for the J/ψ meson and $\pm 15\%$ for the $\psi(2S)$ meson.

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I. INTRODUCTION

The timelike charged-kaon form factor F_K has been measured precisely in the threshold/ ϕ -meson region [1–3] and by several experiments [3–7] in the center-of-mass (c.m.) energy range 1.1–2.4 GeV, where substantial structure is evident. At higher energies, there are precise measurements at 3.671, [8], 3.772, and 4.170 GeV [9], and there is a scan that extends to 5 GeV [3]. The energy dependence of these higher-energy data is consistent with the asymptotic form predicted by perturbative quantum chromodynamics (pQCD), but their magnitude is about a factor of 4 higher than the predicted asymptotic value [10]

$$M_{K^+K^-}^2 |F_K(M_{K^+K^-})| = 8\pi\alpha_s f_K^2, \quad (1)$$

where $M_{K^+K^-}$ is the K^+K^- invariant mass, α_s is the strong coupling constant, and $f_K = 156.2 \pm 0.7 \text{ MeV}$ [[11] p. 1027] is the charged-kaon decay constant.¹ It is expected that the difference between the data and the asymptotic QCD prediction will decrease with increasing energy. Precise measurements at higher energies are needed to test this expectation.

In this paper we analyze the initial-state radiation (ISR) process $e^+e^- \rightarrow K^+K^-\gamma$. The K^+K^- mass spectrum measured in this process is related to the cross section of the nonradiative process $e^+e^- \rightarrow K^+K^-$. Our previous measurement of F_K [3] used the “large-angle” (LA) ISR technique, in which the radiated photon is detected and the $e^+e^- \rightarrow K^+K^-\gamma$ event is fully reconstructed. This gives good precision near threshold, but the cross section decreases rapidly with increasing energy, limiting that measurement to energies below 5 GeV. In this paper we

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utilize small-angle (SA) ISR events, in which the ISR photon is emitted close to the e^+e^- collision axis and so is undetected. This allows us to perform an independent and complementary measurement of the charged-kaon form factor, which has better precision in the range 2.6–5 GeV, and extends the measurements up to 8 GeV.

The Born cross section for the ISR process integrated over the kaon momenta and the photon polar angle is

$$\frac{d\sigma_{K^+K^-\gamma}(M_{K^+K^-})}{dM_{K^+K^-}} = \frac{2M_{K^+K^-}}{s} W(s, x) \sigma_{K^+K^-}(M_{K^+K^-}), \quad (2)$$

where s is the e^+e^- c.m. energy squared, $x \equiv 2E_\gamma^*/\sqrt{s} = 1 - M_{K^+K^-}^2/s$, and E_γ^* is the ISR photon energy in the e^+e^- c.m. frame.² The function $W(s, x)$, describing the probability for single ISR emission at lowest order in quantum electrodynamics, is known to an accuracy better than 0.5% [12–14]. The $e^+e^- \rightarrow K^+K^-$ cross section is given in terms of the form factor by

$$\sigma_{K^+K^-}(M_{K^+K^-}) = \frac{\pi\alpha^2\beta^3 C}{3M_{K^+K^-}^2} |F_K(M_{K^+K^-})|^2, \quad (3)$$

where α is the fine-structure constant, $\beta = \sqrt{1 - 4m_K^2/M_{K^+K^-}^2}$, and C is the final-state correction, which, in particular, takes into account extra photon radiation from the final state (see, e.g., Ref. [15]). In the mass region under study the factor C is close to unity, and varies from 1.008 at 2.6 GeV/ c^2 to 1.007 at 8 GeV/ c^2 .

In addition to the form factor, we measure the branching fractions for the decays $J/\psi \rightarrow K^+K^-$ and $\psi(2S) \rightarrow K^+K^-$. For the latter we study the interference between the resonant and nonresonant $e^+e^- \rightarrow K^+K^-$ amplitudes, and between the single-photon and strong $\psi \rightarrow K\bar{K}$ amplitudes (with $\psi = J/\psi, \psi(2S)$). As a result, we extract the interference corrections to the $J/\psi \rightarrow K^+K^-$ and $\psi(2S) \rightarrow K^+K^-$ branching fractions, which were not taken into account in previous measurements and determine the values of the phase difference between the single-photon and strong amplitudes in $J/\psi \rightarrow K\bar{K}$ and $\psi(2S) \rightarrow K\bar{K}$ decays. In contrast to previous determinations of this phase [16–18], we use a model-independent approach, calculating the single-photon decay amplitudes from our data on the charged-kaon form factor.

II. THE BABAR DETECTOR, DATA, AND SIMULATED SAMPLES

We analyze a data sample corresponding to an integrated luminosity of 469 fb⁻¹ [19] recorded with the BABAR detector at the SLAC PEP-II asymmetric-energy (9-GeV

e^- and 3.1-GeV e^+) collider. About 90% of the data were collected at an e^+e^- c.m. energy of 10.58 GeV (the $\Upsilon(4S)$ mass) and 10% at 10.54 GeV.

The BABAR detector is described in detail elsewhere [20]. Charged-particle tracking is provided by a five-layer silicon vertex tracker (SVT) and a 40-layer drift chamber (DCH), operating in the 1.5 T magnetic field of a superconducting solenoid. The position and energy of a photon-produced cluster are measured with a CsI(Tl) electromagnetic calorimeter (EMC). Charged-particle identification (PID) is provided by specific ionization (dE/dx) measurements in the SVT and DCH, and by an internally reflecting ring-imaging Cherenkov detector (DIRC). Muons are identified in the solenoid's instrumented flux return (IFR).

Simulated samples of signal events, and background $e^+e^- \rightarrow \pi^+\pi^-\gamma$ and $\mu^+\mu^-\gamma$ events, are generated with the Phokhara [21] Monte Carlo (MC) event generator, which takes into account next-to-leading-order radiative corrections. To obtain realistic estimates for the pion and kaon cross sections, the experimental values of the pion and kaon electromagnetic form factors measured in the CLEO experiment at $\sqrt{s} = 3.67$ GeV [8] are used in the event generator. The mass dependence of the form factors is assumed to be $1/m^2$, as predicted by asymptotic QCD [10]. The process $e^+e^- \rightarrow e^+e^-\gamma$ is simulated with the BHWIDE event generator [22].

Two-photon background from the process $e^+e^- \rightarrow e^+e^-K^+K^-$ is simulated with the GamGam event generator [23]. Background contributions from $e^+e^- \rightarrow q\bar{q}(\gamma_{\text{ISR}})$, where q represents a u, d, s or c quark, are simulated with the JETSET event generator [24].

The detector response is simulated using the Geant4 package [25]. The simulation takes into account the variations in the detector and beam-background conditions over the running period of the experiment.

III. EVENT SELECTION

We select events with two tracks of opposite charge originating from the interaction region. The tracks must lie in the polar angle range $25.8^\circ < \theta < 137.5^\circ$ and be identified as kaons. The selected kaon candidates are fitted to a common vertex with a beam-spot constraint. The χ^2 probability for this fit is required to be greater than 0.1%.

Conditions on the K^+K^- transverse momentum (p_{T,K^+K^-}) and the missing-mass squared (M_{miss}^2) recoiling against the K^+K^- system are used for further selection. The p_{T,K^+K^-} distribution for simulated $e^+e^- \rightarrow K^+K^-\gamma$ events is shown in Fig. 1. The peak near zero corresponds to ISR photons emitted along the collision axis, while the long tail is due to photons emitted at large angles. We apply the condition $p_{T,K^+K^-} < 0.15$ GeV/ c , which removes large-angle ISR and suppresses backgrounds from $e^+e^- \rightarrow K^+K^-\pi^0$ and ISR processes with extra π^0 mesons.

²Throughout this paper, an asterisk denotes a quantity that is evaluated in the e^+e^- c.m. frame, while quantities without asterisks are evaluated in the laboratory frame.

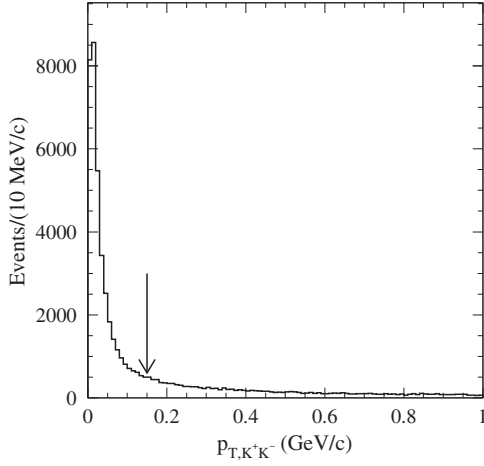


FIG. 1. The p_{T,K^+K^-} distribution for simulated $e^+e^- \rightarrow K^+K^- \gamma$ events. The arrow indicates $p_{T,K^+K^-} = 0.15$ GeV/c.

The region of low K^+K^- invariant mass cannot be studied with small-angle ISR due to limited detector acceptance. A K^+K^- pair with $p_{T,K^+K^-} < 0.15$ GeV/c is detected in *BABAR* when its invariant mass is larger than 2.5 (4.2) GeV/c² for an ISR photon emitted along the electron (positron) beam direction. The average values of the kaon momentum for the two photon directions are about 2.5 and 5 GeV/c, respectively. Since the probability for particle misidentification increases strongly with increasing momentum, we reject events with an ISR photon along the positron direction.

The M_{miss}^2 distribution for simulated signal events is shown in Fig. 2. The signal distribution is peaked at zero, while the background distributions are shifted to negative values for $e^+e^- \rightarrow e^+e^- \gamma$ and $\mu^+\mu^- \gamma$ events and to positive values for $p\bar{p}\gamma$, two-photon and other ISR events. The condition $|M_{\text{miss}}^2| < 1$ GeV²/c⁴ is applied to suppress background. Sideband regions in M_{miss}^2 and p_{T,K^+K^-} are

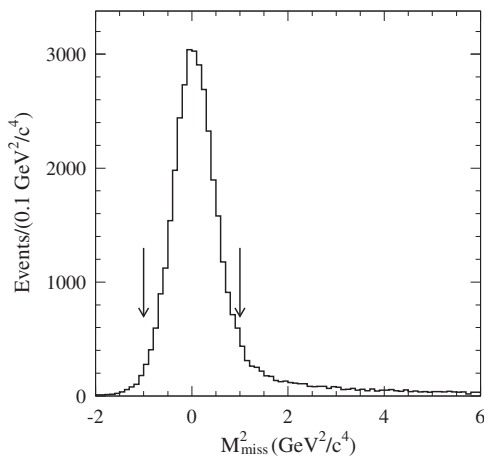


FIG. 2. The M_{miss}^2 distribution for simulated $e^+e^- \rightarrow K^+K^- \gamma$ events. The arrows indicate $|M_{\text{miss}}^2| = 1$ GeV²/c⁴.

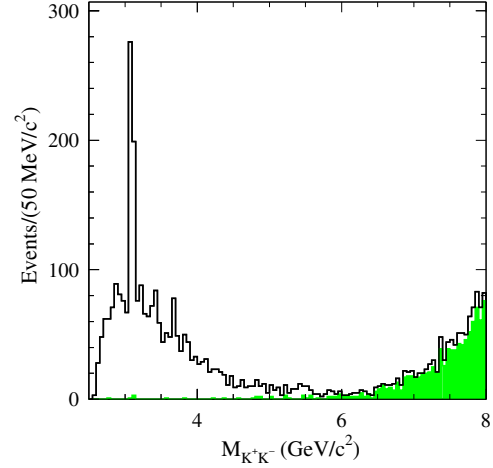


FIG. 3 (color online). The K^+K^- mass spectrum for selected $K^+K^- \gamma$ candidates. The peaks near 3.1, 3.4, and 3.7 GeV/c² are from J/ψ , χ_{c0} , and $\psi(2S)$ decays to K^+K^- , respectively. The shaded histogram shows events with at least one identified muon candidate.

used to estimate the remaining background from these sources, as described in Sec. IV.

The K^+K^- invariant-mass spectrum for events selected with the criteria described above is shown in Fig. 3. A clear J/ψ signal is seen in the spectrum, and there are also indications of small $\psi(2S)$ and χ_{c0} peaks. The χ_{c0} mesons are produced in the reaction $e^+e^- \rightarrow \psi(2S)\gamma \rightarrow \chi_{c0}\gamma\gamma$. The increase in the number of events for $M_{K^+K^-} > 6$ GeV/c² is due to background from the $e^+e^- \rightarrow \mu^+\mu^- \gamma$ process. To suppress the muon background, we apply the additional condition that neither kaon candidate be identified as a muon. Muon identification is based mainly on IFR information, and does not make use of the DIRC and dE/dx measurements used for charged-kaon PID. For $\mu^+\mu^- \gamma$ background events, the probability for at least one of the charged particles to be identified as a muon is about 88% (see subsection IV A). The shaded histogram in Fig. 3 shows events with at least one identified muon candidate. The large muon background for larger values of $M_{K^+K^-}$ prevents us from providing results for $M_{K^+K^-} > 8$ GeV/c².

The mass spectrum with finer binning in the region of the charmonium resonances (3.0–4.5 GeV/c²) is presented in the Supplemental Material [26], together with the mass resolution functions obtained from MC simulation with $M_{K^+K^-}$ near the J/ψ and $\psi(2S)$ masses.

IV. BACKGROUND ESTIMATION AND SUBTRACTION

Sources of background in the selected sample include the following: other two-body ISR processes $e^+e^- \rightarrow e^+e^- \gamma$, $\mu^+\mu^- \gamma$, $\pi^+\pi^- \gamma$, and $p\bar{p}\gamma$; ISR processes containing additional neutral particles, e.g., $e^+e^- \rightarrow K^+K^- \pi^0 \gamma$ and $e^+e^- \rightarrow \psi(2S)\gamma \rightarrow \chi_{cJ}\gamma\gamma \rightarrow K^+K^- \gamma\gamma$; the two-photon

process $e^+e^- \rightarrow e^+e^-K^+K^-$; and nonradiative $e^+e^- \rightarrow q\bar{q}$ events containing a K^+K^- pair plus neutrals, e.g., $e^+e^- \rightarrow K^+K^-\pi^0$. The background from the process $e^+e^- \rightarrow K^+K^-\pi^0$, which was dominant in our LA analysis [3], is strongly suppressed by the requirement on p_{T,K^+K^-} , and is found to be negligible in the SA analysis. The cross section for $e^+e^- \rightarrow p\bar{p}\gamma$ [27,28] is smaller than that for $e^+e^- \rightarrow K^+K^-\gamma$ in the mass region of interest, and this background is reduced to a negligible level by the requirement $|M_{\text{miss}}^2| < 1 \text{ GeV}^2/c^4$. The other categories of background are discussed in the following subsections.

A. Background from $e^+e^- \rightarrow e^+e^-\gamma, \mu^+\mu^-\gamma$, and $\pi^+\pi^-\gamma$

To be selected and, thus, to represent background for this analysis, both final-state charged tracks in $e^+e^- \rightarrow e^+e^-\gamma$, $\mu^+\mu^-\gamma$, and $\pi^+\pi^-\gamma$ events must be misidentified as kaons, and the missing-mass squared must be poorly determined.

The probability to misidentify a pion as a kaon has been measured as a function of charge, momentum, and polar angle using a control sample of pions from $K_S \rightarrow \pi^+\pi^-$ decays. Using the measured misidentification probabilities, we calculate weights for simulated $e^+e^- \rightarrow \pi^+\pi^-\gamma$ events (see Sec. II) to be identified as $K^+K^-\gamma$ events and estimate a $\pi^+\pi^-\gamma$ background rate relative to the signal $K^+K^-\gamma$ rate

ranging from 5×10^{-5} at $3 \text{ GeV}/c^2$ to about 5×10^{-3} at $7.5 \text{ GeV}/c^2$.

A similar approach is used to estimate the $e^+e^- \rightarrow e^+e^-\gamma$ background. The electron misidentification rate has been measured using $e^+e^- \rightarrow e^+e^-\gamma$ events with the photon detected at large angles. From MC simulation we estimate the electron contamination to be at most 0.5%. The PID requirements suppress $e^+e^- \rightarrow e^+e^-\gamma$ events by a factor of about 10^8 . We have verified this suppression by analyzing a sample of LA $K^+K^-\gamma$ candidates with the photon detected in the EMC. In this data sample, surviving $e^+e^- \rightarrow e^+e^-\gamma$ events can be identified by requiring a small opening angle between the photon direction and that of one of the charged-particle tracks.

In the subsequent analysis, we disregard possible backgrounds from $e^+e^-\gamma$ and $\pi^+\pi^-\gamma$ events since their contributions are expected to be negligible.

The $e^+e^- \rightarrow \mu^+\mu^-\gamma$ background is non-negligible for large values of $M_{K^+K^-}$. For $M_{K^+K^-} > 5.5 \text{ GeV}/c^2$, we estimate the numbers of signal and background events in each of the five mass intervals listed in Table I by fitting the M_{miss}^2 distributions in the range $[-2, +1] \text{ GeV}^2/c^4$, as shown in Fig. 4, using three components: signal events, the $\mu^+\mu^-\gamma$ background, and the ISR + two-photon background. The M_{miss}^2 interval is extended to negative values to

TABLE I. The number of selected K^+K^- candidates (N_{data}), number of signal events (N_{sig}), and estimated numbers of background events from $e^+e^- \rightarrow \mu^+\mu^-\gamma$ ($N_{\mu\mu\gamma}$), from the two-photon process $e^+e^- \rightarrow e^+e^-K^+K^-$ ($N_{\gamma\gamma}$), and from ISR processes with extra neutral particle(s) such as $e^+e^- \rightarrow K^+K^-\pi^0\gamma$ and $K^+K^-2\pi^0\gamma$ (N_{ISR}). In the last column, $N_{\psi,\chi}$ refers to the background from $J/\psi \rightarrow K^+K^-$ events for $3.0 < M_{K^+K^-} < 3.2 \text{ GeV}/c^2$, from $\psi(2S) \rightarrow K^+K^-$ events for $3.6 < M_{K^+K^-} < 3.8 \text{ GeV}/c^2$, and from $\psi(2S) \rightarrow \chi_{cJ}\gamma \rightarrow K^+K^-\gamma$ events for $3.2 < M_{K^+K^-} < 3.4$ and $3.4 < M_{K^+K^-} < 3.6 \text{ GeV}/c^2$ (see Fig. 6). Events with $M_{K^+K^-} > 5.5 \text{ GeV}/c^2$ are selected with the looser condition $-2 < M_{\text{miss}}^2 < 1 \text{ GeV}^2/c^4$. For N_{sig} , the first uncertainty is statistical and the second is systematic. For the numbers of background events, the combined uncertainty is quoted.

$M_{K^+K^-}$ (GeV/ c^2)	N_{data}	N_{sig}	$N_{\mu\mu\gamma}$	$N_{\gamma\gamma}$	N_{ISR}	$N_{\psi,\chi}$
2.6–2.7	76	$75 \pm 9 \pm 2$	< 0.1	< 2	0.6 ± 0.5	–
2.7–2.8	123	$121 \pm 11 \pm 2$	< 0.1	< 2	1.6 ± 1.0	–
2.8–2.9	160	$157 \pm 13 \pm 2$	< 0.1	2.6 ± 1.9	0.9 ± 0.7	–
2.9–3.0	157	$152 \pm 13 \pm 2$	< 0.1	3.7 ± 2.1	1.3 ± 0.9	–
3.0–3.2	614	$297 \pm 22 \pm 3$	< 0.1	7.3 ± 2.8	2.3 ± 1.6	307.1 ± 21.3
3.2–3.4	290	$279 \pm 17 \pm 2$	< 0.1	5.1 ± 2.1	1.8 ± 1.3	4.6 ± 1.6
3.4–3.6	237	$194 \pm 16 \pm 2$	< 0.1	6.1 ± 1.8	3.1 ± 2.0	33.7 ± 13.8
3.6–3.8	212	$162 \pm 16 \pm 1$	< 0.1	3.2 ± 0.9	1.5 ± 1.0	45.8 ± 11.0
3.8–4.0	156	$152 \pm 13 \pm 1$	< 0.1	2.6 ± 0.6	1.4 ± 1.0	–
4.0–4.2	108	$105 \pm 11 \pm 1$	< 0.1	2.8 ± 0.5	0.3 ± 0.4	–
4.2–4.4	84	$81 \pm 9 \pm 1$	0.2 ± 0.1	1.2 ± 0.2	1.7 ± 1.0	–
4.4–4.6	47	$44.7 \pm 6.9 \pm 0.6$	0.1 ± 0.1	1.2 ± 0.2	1.0 ± 0.7	–
4.6–4.8	43	$41.2 \pm 6.6 \pm 0.3$	0.1 ± 0.1	1.5 ± 0.3	0.2 ± 0.3	–
4.8–5.0	38	$36.2 \pm 6.2 \pm 0.5$	0.5 ± 0.3	0.8 ± 0.2	0.5 ± 0.4	–
5.0–5.2	28	$26.8 \pm 5.3 \pm 0.3$	0.2 ± 0.1	0.6 ± 0.1	0.4 ± 0.4	–
5.2–5.5	47	$45.2 \pm 6.9 \pm 0.6$	0.9 ± 0.5	0.6 ± 0.2	0.3 ± 0.3	–
5.5–6.0	42	$35.3 \pm 6.7 \pm 0.7$	6.0 ± 3.7	0.4 ± 0.3	0.7 ± 0.5	–
6.0–6.5	25	$10.9 \pm 4.6 \pm 1.2$	11.4 ± 4.3	0.3 ± 0.2	2.0 ± 1.1	–
6.5–7.0	34	$13.8 \pm 5.4 \pm 0.7$	18.5 ± 5.6	< 0.3	0.8 ± 0.6	–
7.0–7.5	44	$7.5 \pm 5.3 \pm 1.9$	33.3 ± 6.9	< 0.5	3.4 ± 1.8	–
7.5–8.0	91	$0.0 \pm 7.0 \pm 2.0$	87.6 ± 10.5	< 0.5	3.5 ± 1.9	–

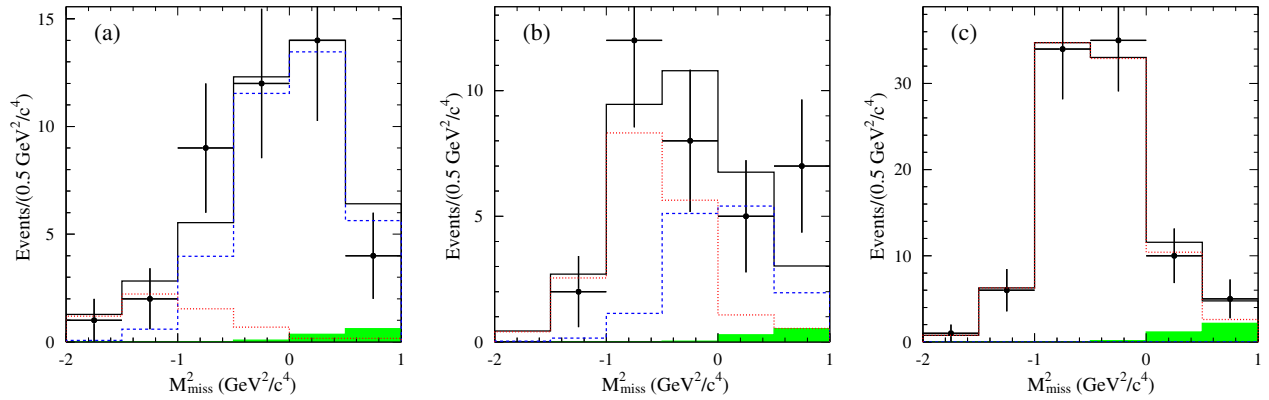


FIG. 4 (color online). The M_{miss}^2 distributions for data (points with error bars) in three $M_{K^+K^-}$ intervals: (a) 5.5–6.0 GeV/c^2 , (b) 6.5–7.0 GeV/c^2 , and (c) 7.5–8.0 GeV/c^2 . The solid histogram is the result of the fit described in the text. The dashed blue, dotted red, and shaded histograms show the contributions of signal, muon background, and ISR + two-photon background, respectively.

increase the sensitivity to $e^+e^- \rightarrow \mu^+\mu^-\gamma$ background and thus to better determine its contribution. The distribution for signal events is taken from simulation and is centered at zero. The distribution for the $\mu^+\mu^-\gamma$ background is obtained using data events with at least one identified muon, and is shifted to negative M_{miss}^2 values because of the muon-kaon mass difference. We also include the small contributions from ISR and two-photon events estimated as described below in Secs. IV B and IV C. The fitted parameters are the numbers of signal (N_{sig}) and muon-background ($N_{\mu\mu\gamma}$) events.

The results of the fits are listed in the last five rows of Table I and are shown in Fig. 4 for three representative intervals of $M_{K^+K^-}$. The first uncertainty in N_{sig} is statistical, while the second, systematic, uncertainty, accounts for the uncertainty in the numbers of ISR and two-photon background events. The N_{sig} results in Table I for $M_{K^+K^-} > 5.5 \text{ GeV}/c^2$ are obtained with the condition $-2 < M_{\text{miss}}^2 < 1 \text{ GeV}^2/c^4$. They can be scaled into our standard selection $|M_{\text{miss}}^2| < 1 \text{ GeV}^2/c^4$ by multiplying the results in the 5.5–6.5 GeV/c^2 and 6.5–7.5 GeV/c^2 mass ranges by 0.98 and 0.99, respectively. For $M_{K^+K^-} > 7.5 \text{ GeV}/c^2$ the scaling factor is consistent with 1.0. The scale factors are determined using simulated signal events.

Below 5.5 GeV/c^2 , where the muon background is small, we adopt a simpler approach and estimate the number of $\mu^+\mu^-\gamma$ background events in each mass interval using the number of selected events $N_{1\mu}$ with at least one charged track identified as a muon. The number of background events is estimated as

$$N_{\mu\mu\gamma} = C_\mu(N_{1\mu} - k_{1\mu}N_{\text{data}}), \quad (4)$$

where C_μ , evaluated as described below, is the ratio of the number of selected $\mu^+\mu^-\gamma$ events with no identified muon to the number of events with at least one identified muon, $k_{1\mu}$ is the fraction of selected $K^+K^-\gamma$ events with at least

one identified muon, and N_{data} is the number of events in the respective $M_{K^+K^-}$ interval. The value of $k_{1\mu}$ is taken from simulated signal events and varies from 0.006 at $M_{K^+K^-} = 2.6 \text{ GeV}/c^2$ to 0.01 at $M_{K^+K^-} = 5.5 \text{ GeV}/c^2$.

Very few simulated $\mu^+\mu^-\gamma$ events have both tracks identified as kaons and neither identified as a muon, so C_μ is studied as a function of $M_{K^+K^-}$ using the probability for an individual muon to be identified as both a kaon and a muon, assuming the probabilities for the two tracks to be independent. We find that C_μ does not exhibit a significant dependence on mass within the range of our measurements, 2.6–8.0 GeV/c^2 . Therefore, C_μ used in Eq. (4) is estimated from the fitted numbers of $\mu^+\mu^-\gamma$ events above 5.5 GeV/c^2 . We find $C_\mu = 0.14 \pm 0.01 \pm 0.08$, where the first uncertainty is from the fits and the second accounts for the full range of values in different mass intervals in data and simulation (for purposes of information, the MC result is $C_\mu = 0.11$). The resulting estimated numbers of $\mu^+\mu^-\gamma$ background events are listed in Table I. For masses below 4.2 GeV/c^2 , $(N_{1\mu} - k_{1\mu}N_{\text{data}})$ is consistent with zero, and we take 0.1 as both an upper limit and uncertainty.

B. Multibody ISR background

Background ISR events containing a K^+K^- pair and one or more π^0 and/or η mesons are distinguishable by their nonzero values of M_{miss}^2 and p_{T,K^+K^-} , but some events with a small number of neutral particles still can enter the selected data sample. Figure 5(a) shows the two-dimensional distribution of M_{miss}^2 versus p_{T,K^+K^-} for data events before the requirements on these two variables, indicated by the lines, are applied. The bottom left rectangle is the signal region. The same distribution for simulated signal events is shown in Fig. 5(b), and is similar in structure except for a deficit in the upper right rectangle, which we take as a sideband region. The distribution for ISR background events produced by JETSET is shown in Fig. 5(c). It should be noted that most (98%) simulated

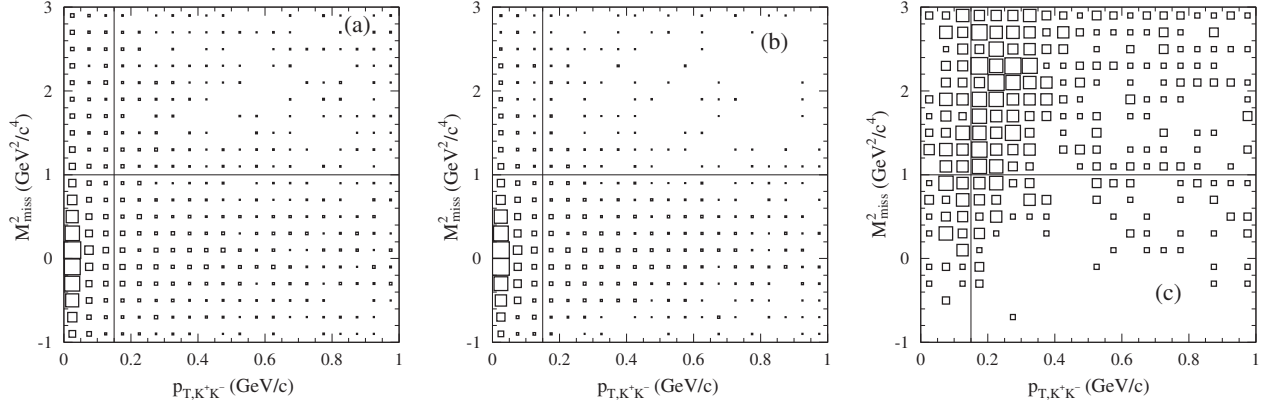


FIG. 5. The distributions of M_{miss}^2 versus p_{T,K^+K^-} for (a) data events, (b) simulated signal events, and (c) simulated ISR background events. Events in the J/ψ and χ_{c0} mass regions, $3.05 < M_{K^+K^-} < 3.15 \text{ GeV}/c^2$ and $3.38 < M_{K^+K^-} < 3.46 \text{ GeV}/c^2$, are excluded from the distributions; regions near the χ_{c2} and $\psi(2S)$ are not excluded, since their signal content is quite small. The lines indicate the boundaries of the signal region (bottom left rectangle) and the sideband region (top right rectangle).

background events in the signal region are from the process $e^+e^- \rightarrow K^+K^-\pi^0\gamma$, while the fraction in the sideband region is about 80%.

The number of data events in the sideband region N_2 is used to estimate the ISR background in the signal region using

$$N_{\text{ISR}} = \frac{N_2 - \beta_{\text{sig}}N_1}{\beta_{\text{bkg}} - \beta_{\text{sig}}}, \quad (5)$$

where N_1 is the number of data events in the signal region, and β_{sig} and β_{bkg} are the N_2/N_1 ratios from signal and background simulation, respectively. The coefficient β_{sig} increases linearly from 0.046 ± 0.005 at $M_{K^+K^-} = 2.6 \text{ GeV}/c^2$ to 0.074 ± 0.005 at $8.0 \text{ GeV}/c^2$, where the uncertainty is statistical, whereas $\beta_{\text{bkg}} = 7.6 \pm 1.0 \pm 4.0$ is independent of $M_{K^+K^-}$. The first uncertainty in β_{bkg} is statistical, and the second is systematic. The latter takes into

account possible differences between data and simulation in the background composition, and in the kinematic distributions of $e^+e^- \rightarrow K^+K^-\pi^0\gamma$ events.

The regions $3.0\text{--}3.2$ and $3.6\text{--}3.8 \text{ GeV}/c^2$ contain resonant contributions from the decays $J/\psi \rightarrow K^+K^-$ and $\psi(2S) \rightarrow K^+K^-$, respectively. The resonant and nonresonant contributions are determined by the fits described in Sec. VII, and such fits are also applied to the sideband regions. The resulting numbers of nonresonant events, N_1 and N_2 , are used to estimate the ISR background.

Similarly, the $M_{K^+K^-}$ regions $3.2\text{--}3.4$ and $3.4\text{--}3.6 \text{ GeV}/c^2$ contain χ_{c0} and χ_{c2} decays, as seen in Fig. 6. The χ_{cJ} states are produced in the reaction $e^+e^- \rightarrow \psi(2S)\gamma$, followed by $\psi(2S) \rightarrow \chi_{cJ}\gamma$. A similar set of fits is used to determine N_1 , N_2 , and the background contribution from the χ_{cJ} states, and the fit results are shown in Fig. 6. The estimated numbers of ISR background events are listed in Table I along with the fitted numbers of ψ and χ_{cJ} decays in the relevant mass intervals.

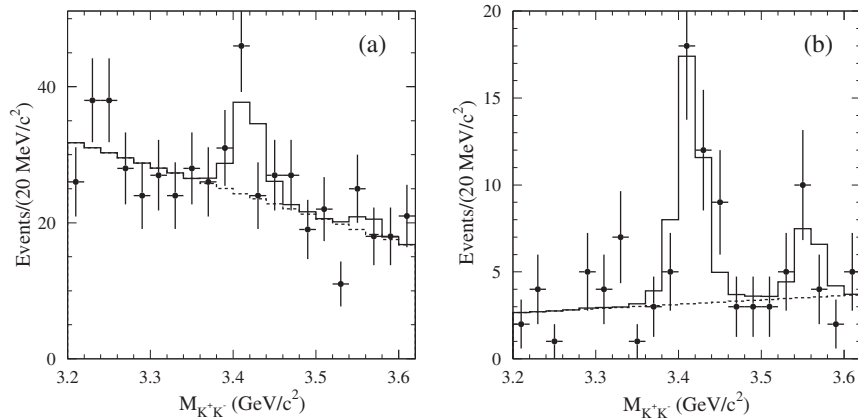


FIG. 6. The $M_{K^+K^-}$ spectra for data (points with error bars) in the vicinity of the χ_{c0} and χ_{c2} resonances for the $M_{\text{miss}}^2 - p_{T,K^+K^-}$ (a) signal region and (b) sideband region. The solid histograms result from the fits described in the text. The dashed histograms represent the nonresonant contributions.

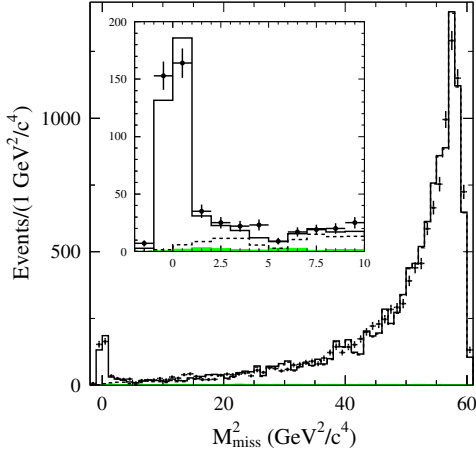


FIG. 7 (color online). The M_{miss}^2 distribution for data (points with error bars) with $2.8 < M_{K^+K^-} < 3.0 \text{ GeV}/c^2$ selected with all the criteria described in Sec. III except for the requirement $|M_{\text{miss}}^2| < 1 \text{ GeV}^2/c^4$. The solid histogram is a sum of signal and background distributions obtained from MC simulation. The dashed histogram shows the distribution for two-photon events, and the shaded (almost invisible) histogram shows the small contribution of all other background processes. The inset shows an enlarged view of the region $-2 < M_{\text{miss}}^2 < 10 \text{ GeV}^2/c^4$. Above $10 \text{ GeV}^2/c^4$ the solid and dashed histograms are indistinguishable.

C. Two-photon background

Two-photon events corresponding to the process $e^+e^- \rightarrow e^+e^-\gamma^*\gamma^* \rightarrow e^+e^-K^+K^-$ are distinguished by their larger values of M_{miss}^2 . Figure 7 shows the M_{miss}^2 distribution for data events in the range $2.8 < M_{K^+K^-} < 3.0 \text{ GeV}/c^2$ that satisfy all the criteria in Sec. III except for that on M_{miss}^2 . The two-photon events, which dominate the large $|M_{\text{miss}}^2|$ region, are generally seen to be well separated from signal events but to nonetheless have a tail that extends into the signal region $|M_{\text{miss}}^2| < 1 \text{ GeV}^2/c^4$. The exact shape of this tail depends on the unknown kaon angular distribution. Therefore, we reweight our simulation (generated with a uniform distribution) to reproduce the $\cos \theta_K$ distribution observed in the data in each $M_{K^+K^-}$ interval; here θ_K is the angle between the K^+ momentum in the K^+K^- rest frame and the e^- beam direction in the e^+e^- c.m. frame. The data and reweighted simulated $\cos \theta_K$ distributions are compared in Fig. 8. The simulated M_{miss}^2 distribution is shown in Fig. 7, where it is seen to reproduce the data well.

The two-photon background in each $M_{K^+K^-}$ interval is estimated from the number of data events with $M_{\text{miss}}^2 > d$ and a scale factor from the simulation. The M_{miss}^2 distribution changes with $M_{K^+K^-}$, and the value of d is $20 \text{ GeV}^2/c^4$ for $M_{K^+K^-} < 4.4 \text{ GeV}/c^2$, $10 \text{ GeV}^2/c^4$ for $M_{K^+K^-} > 6.5 \text{ GeV}/c^2$, and varies linearly in-between. The scale factor ranges from 10^{-4} in the $2.6\text{--}2.7 \text{ GeV}/c^2$ interval to about 10^{-2} in the $7.0\text{--}7.5 \text{ GeV}/c^2$ interval. However, the number of two-photon events decreases with increasing $M_{K^+K^-}$. The estimated background event contributions are listed in Table I.

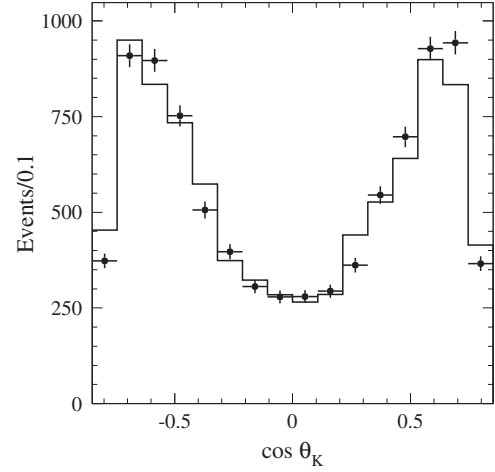


FIG. 8. The $\cos \theta_K$ distribution for data (points with error bars) and reweighted simulated events (histogram) with $M_{\text{miss}}^2 > 20 \text{ GeV}^2/c^4$ from the K^+K^- mass range $2.8\text{--}3.0 \text{ GeV}/c^2$.

The numbers of signal events obtained after background subtraction are listed in Table I. The first uncertainty in N_{sig} is statistical and the second is systematic. The systematic term accounts for the uncertainties in the numbers of $e^+e^- \rightarrow \mu^+\mu^-\gamma$ and two-photon background events, and the uncertainties in the coefficients β_{sig} and β_{bkg} in the ISR background subtraction procedure.

V. DETECTION EFFICIENCY

The detection efficiency, ε_{MC} , determined using MC simulation, is shown in Fig. 9 as a function of $M_{K^+K^-}$. The nonmonotonic behavior observed for $M_{K^+K^-} > 5.5 \text{ GeV}/c^2$ is introduced by filters designed to reduce background before the event-reconstruction stage.

Corrections are applied to ε_{MC} to account for data-MC simulation differences in detector response

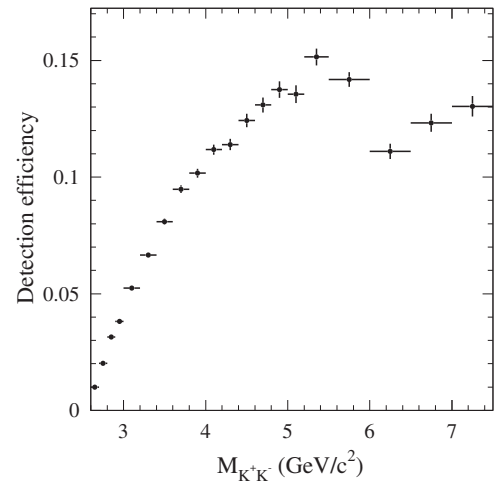


FIG. 9. The K^+K^- mass dependence of the detection efficiency for simulated $e^+e^- \rightarrow K^+K^-\gamma$ events.

TABLE II. Values of the efficiency corrections, δ_i . The three values in the rows “PID” and “total” correspond to $M_{K^+K^-} = 2.6, 6.0, \text{ and } 7.5 \text{ GeV}/c^2$, respectively.

Effect	δ_i (%)
Trigger	-1.0 ± 0.5
PID	$-2.0 \pm 0.4 / -4.0 \pm 0.5 / -10.0 \pm 1.7$
Track reconstruction	0.0 ± 1.6
Requirements on p_{T,K^+K^-} and M_{miss}^2	0.0 ± 1.5
Total	$-3.0 \pm 2.3 / -5.0 \pm 2.3 / -11.0 \pm 2.8$

$$\varepsilon = \varepsilon_{\text{MC}} \prod_{i=1}^4 (1 + \delta_i), \quad (6)$$

where the δ_i terms are the efficiency corrections listed in Table II. The difference between data and simulation in trigger efficiency is studied using the overlap of the samples of events satisfying two independent sets of trigger criteria based on signals from the EMC and DCH. The correction for trigger inefficiency is found to be $(-1.0 \pm 0.5)\%$.

The track-reconstruction efficiency for charged kaons has been studied in events with similar topology [3] and we use the correction derived therein. The charged-kaon identification efficiency is studied as a function of the track momentum and polar angle using a control sample of

kaons from the decay chain $D^{*+} \rightarrow \pi^+ D^0, D^0 \rightarrow K^- \pi^+$. The ratio of the efficiencies is then used to reweight simulated signal events, resulting in an overall correction that varies slowly from from -2% at $2.6 \text{ GeV}/c^2$ to -4% at $6 \text{ GeV}/c^2$, and then falls to about -10% at $7.5 \text{ GeV}/c^2$. The statistical uncertainty in the correction term defines the systematic uncertainty in this correction.

The remaining criteria are based on M_{miss}^2 and p_{T,K^+K^-} , which we believe to be well simulated. The track momentum and angular resolutions have been studied in, e.g., Ref. [27] for $e^+e^- \rightarrow \mu^+\mu^-\gamma$ events with a detected photon. Based on this and similar variables in our previous ISR studies, we make no correction, and assign a conservative systematic uncertainty of 1.5% to cover these remaining factors. The corrections to the detection efficiency are listed in Table II.

VI. THE $e^+e^- \rightarrow K^+K^-$ CROSS SECTION AND THE CHARGED-KAON FORM FACTOR

The $e^+e^- \rightarrow K^+K^-$ cross section in each K^+K^- mass interval i is calculated as

$$\sigma_{K^+K^-,i} = \frac{N_{\text{sig},i}}{\varepsilon_i L_i}. \quad (7)$$

The number of selected events ($N_{\text{sig},i}$) for each K^+K^- mass interval after background subtraction is listed in Table III.

TABLE III. The K^+K^- invariant-mass interval ($M_{K^+K^-}$), number of selected events (N_{sig}) after background subtraction, detection efficiency (ε), ISR luminosity (L), measured $e^+e^- \rightarrow K^+K^-$ cross section ($\sigma_{K^+K^-}$), and the charged-kaon form factor ($|F_K|$). For the number of events and cross section, the first uncertainty is statistical and the second is systematic. For the form factor, we quote the combined uncertainty. For the mass interval $7.5\text{--}8.0 \text{ GeV}/c^2$, the 90% C.L. upper limits for the cross section and form factor are listed.

$M_{K^+K^-}$ (GeV/ c^2)	N_{sig}	ε (%)	L (pb $^{-1}$)	$\sigma_{K^+K^-}$ (pb)	$ F_K \times 100$
2.6–2.7	$75 \pm 9 \pm 2$	0.96	113	$69.3 \pm 8.1 \pm 3.7$	16.7 ± 1.1
2.7–2.8	$121 \pm 11 \pm 2$	1.94	118	$52.9 \pm 4.9 \pm 2.4$	15.0 ± 0.8
2.8–2.9	$157 \pm 13 \pm 2$	3.03	122	$42.0 \pm 3.4 \pm 1.6$	13.7 ± 0.6
2.9–3.0	$152 \pm 13 \pm 2$	3.69	127	$32.2 \pm 2.7 \pm 1.2$	12.4 ± 0.6
3.0–3.2	$297 \pm 22 \pm 3$	5.07	271	$21.7 \pm 1.6 \pm 0.6$	10.6 ± 0.4
3.2–3.4	$279 \pm 17 \pm 2$	6.43	292	$14.8 \pm 0.9 \pm 0.4$	9.2 ± 0.3
3.4–3.6	$194 \pm 16 \pm 2$	7.81	313	$7.92 \pm 0.63 \pm 0.24$	7.1 ± 0.3
3.6–3.8	$162 \pm 16 \pm 1$	9.15	336	$5.26 \pm 0.51 \pm 0.16$	6.1 ± 0.3
3.8–4.0	$152 \pm 13 \pm 1$	9.80	361	$4.30 \pm 0.36 \pm 0.13$	5.7 ± 0.3
4.0–4.2	$105 \pm 11 \pm 1$	10.8	386	$2.52 \pm 0.25 \pm 0.08$	4.60 ± 0.25
4.2–4.4	$81 \pm 9 \pm 1$	11.0	413	$1.79 \pm 0.20 \pm 0.06$	4.05 ± 0.25
4.4–4.6	$44.7 \pm 6.9 \pm 0.6$	11.9	442	$0.85 \pm 0.13 \pm 0.03$	2.91 ± 0.23
4.6–4.8	$41.2 \pm 6.6 \pm 0.3$	12.5	473	$0.70 \pm 0.11 \pm 0.03$	2.74 ± 0.23
4.8–5.0	$36.2 \pm 6.2 \pm 0.5$	13.1	507	$0.55 \pm 0.09 \pm 0.02$	2.52 ± 0.22
5.0–5.2	$26.8 \pm 5.3 \pm 0.3$	12.9	543	$0.38 \pm 0.08 \pm 0.02$	2.19 ± 0.22
5.2–5.5	$45.2 \pm 6.9 \pm 0.6$	14.4	888	$0.35 \pm 0.05 \pm 0.01$	2.21 ± 0.18
5.5–6.0	$34.6 \pm 6.6 \pm 0.7$	13.5	1710	$0.150 \pm 0.029 \pm 0.006$	1.54 ± 0.15
6.0–6.5	$10.7 \pm 4.5 \pm 1.2$	10.6	2062	$0.049 \pm 0.021 \pm 0.006$	0.95 ± 0.22
6.5–7.0	$13.6 \pm 5.3 \pm 0.7$	11.6	2523	$0.047 \pm 0.018 \pm 0.004$	1.00 ± 0.20
7.0–7.5	$7.4 \pm 5.3 \pm 1.9$	11.9	3144	$0.020 \pm 0.014 \pm 0.004$	$0.70_{-0.36}^{+0.23}$
7.5–8.0	$0.0 \pm 6.9 \pm 2.0$	9.36	4015	< 0.024	< 0.9

The N_{sig} values for $M_{K^+K^-} > 5.5 \text{ GeV}/c^2$ differ from the corresponding values in Table I. They are corrected to correspond to the nominal selection $|M_{\text{miss}}^2| < 1 \text{ GeV}^2/c^4$ as described in Sec. IV A. The first uncertainty in N_{sig} is statistical; the second is systematic due to background subtraction. The value of the ISR luminosity L_i is obtained by integrating $W(s, x)$ from Refs. [12,13] over mass interval i and is listed in Table III. The formulas from Refs. [12,13] include higher-order radiative corrections. However, we do not include in $W(s, x)$ corrections for leptonic and hadronic vacuum polarization in the photon propagator. Cross sections obtained in this way are referred to as “dressed.”

The values obtained for the $e^+e^- \rightarrow K^+K^-$ cross section are listed in Table III. For the mass intervals $3.0\text{--}3.2 \text{ GeV}/c^2$ and $3.6\text{--}3.8 \text{ GeV}/c^2$ we quote the non-resonant cross sections with the respective J/ψ and $\psi(2S)$ contributions excluded. The quoted uncertainties are statistical and systematic, respectively. The statistical uncertainty results from the statistical uncertainty in the number of selected $K^+K^- \gamma$ events. The systematic uncertainty includes the systematic uncertainty in the number of events, the statistical uncertainty in the detection efficiency (1.5%–4.0%), and the uncertainties in the efficiency correction (2.6%–5.1%), the integrated luminosity (0.5%) [19], and the ISR luminosity (0.5%) [12,13]. For the mass interval $7.5\text{--}8.0 \text{ GeV}/c^2$, the 90% confidence level (C.L.) upper limit on the cross section is listed. The measured cross section is shown in Fig. 10.

It is more convenient to perform comparisons with previous measurements and theoretical predictions in terms of the form factor. The values of the charged-kaon electromagnetic form factor obtained using Eq. (3) are listed in Table III. The form factor is plotted in Fig. 11 as a function of $M_{K^+K^-}$ over the range $2.6\text{--}5.0 \text{ GeV}/c^2$, together with all

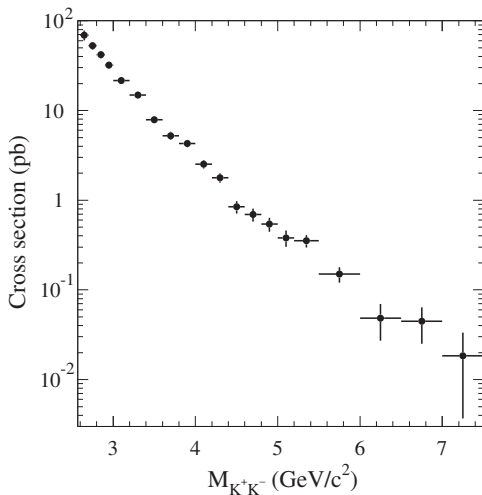


FIG. 10. The $e^+e^- \rightarrow K^+K^-$ cross section measured in this work.

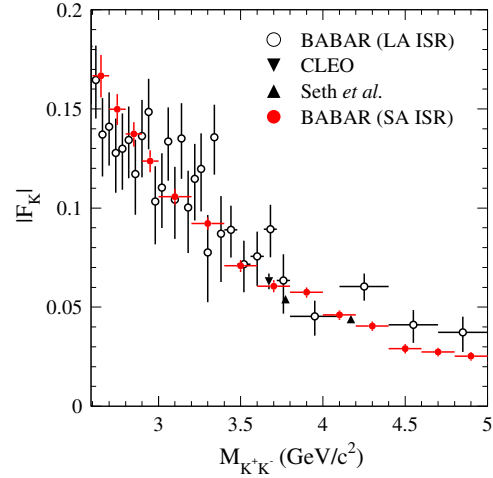


FIG. 11 (color online). The charged-kaon electromagnetic form factor measured in this analysis [BABAR (SA ISR)] in comparison with previous measurements: CLEO [8], Seth *et al.* [9], and BABAR (LA ISR) [3] in the mass region $2.6\text{--}5.0 \text{ GeV}/c^2$. Only statistical uncertainties are shown.

other measurements [3,8,9]. The present measurement is consistent within the uncertainties with our previous, independent LA ISR result [3], which used a data sample corresponding only to 232 fb^{-1} , and provides a much better constraint on the mass dependence in this region.

To compare our results with the precise measurements of Refs. [8,9], we plot in Fig. 12 the scaled form factor $M_{K^+K^-}^2 |F_K(M_{K^+K^-})|$ and fit our data with a smooth

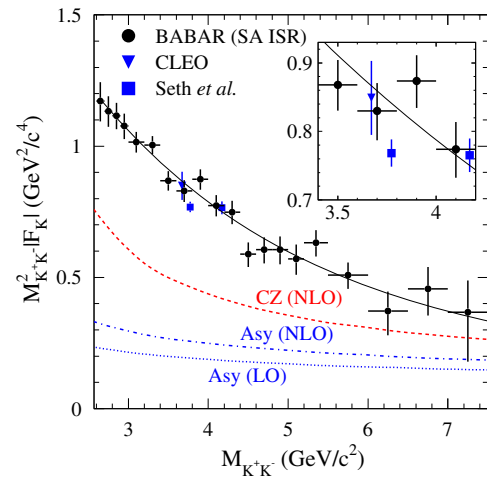


FIG. 12 (color online). The scaled charged-kaon electromagnetic form factor measured in this analysis [BABAR (SA ISR)] and in other experiments: CLEO [8] and Seth *et al.* [9]. Our data are approximated by a smooth curve. The dotted blue curve shows the LO pQCD prediction for the form factor obtained with the asymptotic kaon distribution amplitude. The dot-dashed blue and dashed red curves are the NLO pQCD predictions obtained with the asymptotic and Chernyak-Zhitnitsky distribution amplitudes, respectively. The inset shows an enlarged version of the mass region $3.4\text{--}4.2 \text{ GeV}/c^2$.

function $x^2|F_K(x)| = A/(x^\gamma + B)$, where A , B , and γ are fitted parameters. As seen in the inset, the CLEO and Seth *et al.* points at 3.67 and 4.17 GeV/ c^2 , respectively, are consistent with this function, whereas the Seth *et al.* point at 3.772 GeV/ c^2 lies about three standard deviations below. Since this data point is obtained at the maximum of the $\psi(3770)$ resonance, the deviation may be a result of interference between the resonant and nonresonant amplitudes of the $e^+e^- \rightarrow K^+K^-$ reaction, which we discuss in the next section.

The dotted curve in Fig. 12 represents the leading-order (LO) asymptotic pQCD prediction of Eq. (1), calculated with $\alpha_s(M_{K^+K^-}^2/4)$ [29]. It lies well below most of the data, which might be explained by higher-order pQCD, power corrections, and a deviation of the kaon distribution amplitude (DA), which describes the quark relative momentum distribution inside the meson, from its asymptotic shape. The dot-dashed curve represents the leading-twist, next-to-leading-order (NLO) prediction using the asymptotic DA, obtained by multiplying the pion form factor from Ref. [29] by a factor $f_K^2/f_\pi^2 \approx 1.45$. The NLO correction leads to an increase of about 20%, nearly independent of mass. The dashed curve represents the NLO prediction using the Chernyak-Zhitnitsky (CZ) DA, obtained by multiplying the result in Ref. [29] by a factor of 0.95 [30].

Our data lie well above all predictions. However, they decrease faster than $\alpha_s/M_{K^+K^-}^2$ as $M_{K^+K^-}$ increases, and are consistent with an approach to the pQCD prediction at higher mass. In particular, the ratio of the measured form factor to the asymptotic pQCD prediction (curve Asy(LO) in Fig. 12) changes from about 5.3 at 3 GeV/ c^2 to about 2.6 at 7 GeV/ c^2 .

VII. J/ψ AND $\psi(2S)$ DECAYS INTO K^+K^-

To study the production of K^+K^- pairs through the J/ψ and $\psi(2S)$ resonances, we increase the detection efficiency by selecting events with the looser requirements $p_{T,K^+K^-} < 1$ GeV/ c and $-2 < M_{\text{miss}}^2 < 3$ GeV $^2/c^4$. The resulting

K^+K^- mass spectra in the J/ψ and $\psi(2S)$ mass regions are shown in Fig. 13. Each of these spectra is fitted with the sum of a signal probability density function (PDF) and a linear background. The signal PDF is a Breit-Wigner (BW) function convolved with a double-Gaussian function describing signal resolution. In each fit, the BW mass and width are fixed to their known values [11] and the nominal resolution parameters are taken from simulation. In order to account for deficiencies in the simulation, a mass shift ΔM is allowed, and an increase in both Gaussian widths by a term σ_G added in quadrature is introduced. The free parameters in the J/ψ fit are the numbers of signal and background events, the slope of the background function, ΔM , and σ_G . In the $\psi(2S)$ fit, σ_G and ΔM are fixed to the values obtained for the J/ψ .

The fitted curves are shown in Fig. 13. The numbers of J/ψ and $\psi(2S)$ events are found to be 462 ± 28 and 66 ± 13 , respectively. Results for other fitted parameters are $\sigma_G = 3 \pm 3$ MeV/ c^2 and $\Delta M = M_{J/\psi} - M_{J/\psi}^{\text{MC}} = (0.0 \pm 0.9)$ MeV/ c^2 . The σ_G value corresponds to a difference of about 4% in mass resolution (11 MeV/ c^2 at the J/ψ) between data and simulation. The detection efficiencies, corrected for the data-MC simulation difference in detector response, are $(7.6 \pm 0.2)\%$ for the J/ψ and $(13.4 \pm 0.3)\%$ for the $\psi(2S)$.

The total cross sections for the processes $e^+e^- \rightarrow \psi\gamma \rightarrow K^+K^-\gamma$, where ψ is a narrow resonance like the J/ψ or $\psi(2S)$, is proportional to the electronic width of the resonance and its branching fraction into K^+K^- , i.e., $\sigma_{\psi\gamma} = a_\psi \Gamma(\psi \rightarrow e^+e^-) \mathcal{B}(\psi \rightarrow K^+K^-)$. The coefficient a_ψ can be calculated by integrating Eq. (2) with $\sigma_{K^+K^-}$ set to the appropriate BW function. Using $W(s, x)$ from Refs. [12,13], we obtain $a_{J/\psi} = 6.91$ nb/keV and $a_{\psi(2S)} = 6.07$ nb/keV.

From the measured values of the cross sections, $\sigma_{\psi\gamma}^{\text{exp}} = N_\psi/(\epsilon L)$, we obtain the measured values of the products $\Gamma(\psi \rightarrow e^+e^-) \mathcal{B}(\psi \rightarrow K^+K^-)$ listed in Table IV. The term ‘‘measured value’’ is used because the value of the product obtained in this way may differ from the true value

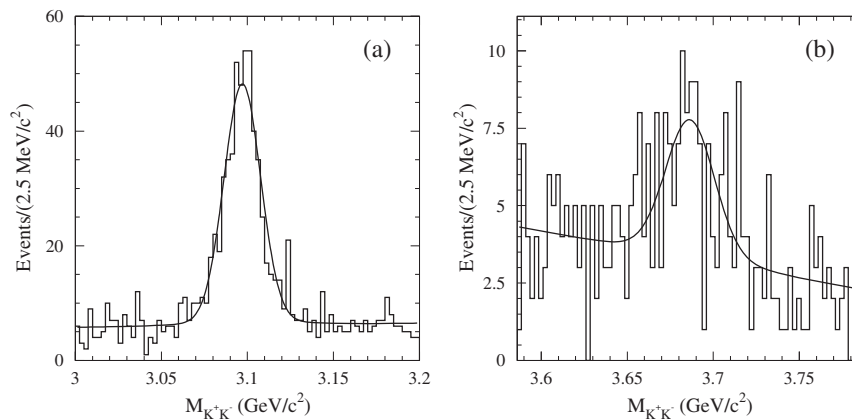


FIG. 13. The K^+K^- mass spectra in the regions near the J/ψ (a) and $\psi(2S)$ (b) resonances. The curves exhibit the results of the fits described in the text.

TABLE IV. The products $\Gamma(\psi \rightarrow e^+e^-)\mathcal{B}(\psi \rightarrow K^+K^-)$ and the branching fractions $\mathcal{B}(\psi \rightarrow K^+K^-)$ obtained in this work for the J/ψ and $\psi(2S)$ resonances. The directly measured values are shown in the rows labeled “measured values.” In the rows marked as “corrected,” the values of the products and branching fractions corrected for the shift due to interference between resonant and nonresonant amplitudes are listed for opposite signs of $\sin\varphi$. In the row marked “ $e^+e^- \rightarrow K^+K^-$ average,” we give the average of the values of the branching fractions measured in the reaction $e^+e^- \rightarrow K^+K^-$. In the row “ $\psi(2S) \rightarrow J/\psi\pi^+\pi^-$, $J/\psi \rightarrow K^+K^-$ ”, the result for the $J/\psi \rightarrow K^+K^-$ branching fraction obtained in Ref. [18] is reported.

$\Gamma(\psi \rightarrow e^+e^-)\mathcal{B}(\psi \rightarrow K^+K^-)$ (eV)	J/ψ	$\psi(2S)$
Measured value	$1.86 \pm 0.11 \pm 0.05$	$0.173 \pm 0.035 \pm 0.005$
Corrected with $\sin\varphi > 0$	$1.78 \pm 0.11 \pm 0.05$	$0.147 \pm 0.035 \pm 0.005$
Corrected with $\sin\varphi < 0$	$1.94 \pm 0.11 \pm 0.05$	$0.197 \pm 0.035 \pm 0.005$
$\mathcal{B}(\psi \rightarrow K^+K^-) \times 10^4$	J/ψ	$\psi(2S)$
Measured value	$3.36 \pm 0.20 \pm 0.12$	$0.73 \pm 0.15 \pm 0.02$
Corrected with $\sin\varphi > 0$	$3.22 \pm 0.20 \pm 0.12$	$0.62 \pm 0.15 \pm 0.02$
Corrected with $\sin\varphi < 0$	$3.50 \pm 0.20 \pm 0.12$	$0.83 \pm 0.15 \pm 0.02$
$e^+e^- \rightarrow K^+K^-$ average	2.43 ± 0.26 [3,31,32]	0.71 ± 0.05 [11]
$\psi(2S) \rightarrow J/\psi\pi^+\pi^-$, $J/\psi \rightarrow K^+K^-$ [18]	2.86 ± 0.21	

due to interference with the nonresonant process $e^+e^- \rightarrow K^+K^-$, as discussed below. The quoted systematic uncertainty includes the uncertainties in the detection efficiency, the integrated luminosity (0.5%), and the theoretical uncertainty in the ISR luminosity (0.5%). Using the nominal values of the electronic widths [11], $\Gamma(J/\psi \rightarrow e^+e^-) = 5.55 \pm 0.14$ keV and $\Gamma(\psi(2S) \rightarrow e^+e^-) = 2.37 \pm 0.04$ keV, we calculate the measured values of the $\psi \rightarrow K^+K^-$ branching fractions listed in Table IV. Since the decay $\psi(2S) \rightarrow K^+K^-$ was studied previously only in the reaction $e^+e^- \rightarrow K^+K^-$, our measurement of $\mathcal{B}(\psi(2S) \rightarrow K^+K^-)$ can be directly compared with the PDG value [11], $(0.71 \pm 0.05) \times 10^{-4}$. Although it is less precise, our measured value agrees well with that from Ref. [11]. For $J/\psi \rightarrow K^+K^-$ there are several measurements [3,31,32] in the $e^+e^- \rightarrow K^+K^-$ reaction, and one measurement [18] in which J/ψ 's were produced in the $\psi(2S) \rightarrow J/\psi\pi^+\pi^-$ decay. To compare with the e^+e^- measurements, we calculate the average of the results [3,31,32], and obtain the value $(2.43 \pm 0.26) \times 10^{-4}$. Our result is larger than this average by 2.7 standard deviations. A comparison with the measurement of Ref. [18] is presented below after applying a correction for interference.

To estimate the effect of interference, we represent the c.m. energy (E) dependence of the $e^+e^- \rightarrow K^+K^-$ cross section near the ψ resonance as [33]

$$\begin{aligned}
\sigma_{K^+K^-}(E) &= \sigma_0 \left| 1 - \sqrt{\frac{\sigma_\psi}{\sigma_0}} (A_\gamma + A_s e^{i\varphi}) \frac{m\Gamma}{D} \right|^2 \\
&= \sigma_0 + \sigma_\psi \left[\mathcal{B}(\psi \rightarrow K^+K^-) + 2\sqrt{\frac{\sigma_0}{\sigma_\psi}} A_s \sin\varphi \right] \\
&\quad \times \frac{m^2\Gamma^2}{|D|^2} - 2\sqrt{\sigma_0\sigma_\psi} (A_\gamma + A_s \cos\varphi) \frac{m\Gamma(m^2 - E^2)}{|D|^2},
\end{aligned} \tag{8}$$

where σ_0 is the nonresonant cross section [Eq. (3)], $\sigma_\psi = (12\pi/m^2)\mathcal{B}(\psi \rightarrow e^+e^-)$, A_γ and A_s are the moduli of the single-photon and strong ψ decay amplitudes, respectively, φ is their relative phase, $D = m^2 - E^2 - im\Gamma$, and m and Γ are the resonance mass and width, respectively. The decay amplitudes are defined such that $\mathcal{B}(\psi \rightarrow K^+K^-) = |A_\gamma + A_s e^{i\varphi}|^2$. The value of the single-photon contribution is related to the kaon form factor through

$$A_\gamma^2 = \mathcal{B}(\psi \rightarrow e^+e^-) \frac{|F_K(m)|^2}{4} \beta^3(m), \tag{9}$$

where β is the phase-space factor from Eq. (3).

For narrow resonances, the interference term proportional to $m^2 - E^2$ integrates to zero due to the beam energy spread in direct e^+e^- experiments and detector resolution in ISR measurements. The remaining interference term has a BW shape and causes a shift of the measured $\mathcal{B}(\psi \rightarrow K^+K^-)$ relative to its true value by

$$\delta\mathcal{B} = 2\sqrt{\frac{\sigma_0}{\sigma_\psi}} A_s \sin\varphi. \tag{10}$$

The values of $\cos\varphi$ and A_s can be obtained from a combined analysis of the $\psi \rightarrow K^+K^-$ and $\psi \rightarrow K_S K_L$ decays, whose branching fractions depend on the same strong amplitude [16,17]

$$\begin{aligned}
\mathcal{B}(\psi \rightarrow K^+K^-) &= |A_\gamma^{K^+K^-} + A_s e^{i\varphi}|^2, \\
\mathcal{B}(\psi \rightarrow K_S K_L) &= |\kappa A_\gamma^{K^+K^-} + A_s e^{i\varphi}|^2,
\end{aligned} \tag{11}$$

where κ is the ratio of the single-photon amplitudes for the $\psi \rightarrow K_S K_L$ and $\psi \rightarrow K^+K^-$ decays, and $|\kappa| = A_\gamma^{K_S K_L} / A_\gamma^{K^+K^-}$. It is expected that in the energy region under study,

the single-photon amplitudes have the same sign of the real parts, and similar ratios of the imaginary-to-real parts [34]; i.e., κ is a positive real number to a good approximation.

For the branching fraction $\mathcal{B}(J/\psi \rightarrow K^+K^-)$ in Eqs. (11), we use the average of the existing measurements [3,18,31,32] and our result. For $\mathcal{B}(J/\psi \rightarrow K_S K_L)$ there are two relatively precise measurements, $(1.82 \pm 0.14) \times 10^{-4}$ [35] and $(2.62 \pm 0.21) \times 10^{-4}$ [18], which are not consistent with each other. We solve Eqs. (11) separately for these two values of $\mathcal{B}(J/\psi \rightarrow K_S K_L)$. The branching fractions measured in the reactions $e^+e^- \rightarrow K^+K^-$ and $e^+e^- \rightarrow K_S K_L$ are corrected as $\mathcal{B} \rightarrow \mathcal{B} - \delta\mathcal{B}$ before averaging, where $\delta\mathcal{B}$ is given by Eq. (10). This correction is not needed for the measurements of Ref. [18], in which the J/ψ mesons are produced in $\psi(2S) \rightarrow J/\psi\pi^+\pi^-$ decays. For $\psi(2S)$ decays we use the branching fraction values from Ref. [11] corrected using Eq. (10).

The coefficient $|\kappa|$ in Eqs. (11) is equal to the ratio of the neutral- and charged-kaon form factors $|F_{K^0}/F_K|$. Data on F_{K^0} above 2 GeV are scarce. There are two measurements [5,36] near 2 GeV, from which we estimate $|\kappa| = 0.28 \pm 0.08$, and there is only one measurement at higher energy, namely $|\kappa| = 0.12 \pm 0.04$ at 4.17 GeV [37]. Using linear interpolation we estimate $|\kappa| = 0.2 \pm 0.1$ at the mass of the J/ψ and 0.15 ± 0.07 at the mass of the $\psi(2S)$. The values of the charged-kaon form factor, $F_K(M_{J/\psi}) = 0.107 \pm 0.002$ and $F_K(M_{\psi(2S)}) = 0.0634 \pm 0.0014$, needed to calculate $A_\gamma^{K^+K^-}$, are taken from the fit to our form factor data shown in Fig. 12.

The values of φ and $\delta\mathcal{B}(\psi \rightarrow K^+K^-)$ obtained using Eqs. (10) and (11) are listed in Tables V and VI. Since Eqs. (11) do not allow us to determine the sign of $\sin\varphi$, the calculations are performed twice, once assuming $\sin\varphi > 0$ and once assuming $\sin\varphi < 0$. For the results in Table V the two upper (bottom) rows marked ‘‘BES’’ (‘‘Seth *et al.*’’) present results obtained using $\mathcal{B}(J/\psi \rightarrow K_S K_L)$ from Ref. [35] (Ref. [18]). We also list the results obtained for $\kappa = 0$, corresponding to the assumption $A_\gamma^{K_S K_L} \ll A_\gamma^{K^+K^-}$ used for most previous determinations of φ . It is

TABLE V. The relative phase (φ) between the single-photon and strong amplitudes for $J/\psi \rightarrow K\bar{K}$ decays calculated with $\kappa = A_\gamma^{K_S K_L}/A_\gamma^{K^+K^-} = 0.2 \pm 0.1$ and $\kappa = 0$, and the correction to the value of $\mathcal{B}(J/\psi \rightarrow K^+K^-)$ measured in the reaction $e^+e^- \rightarrow K^+K^-$. The calculation is performed for the value of $\mathcal{B}(J/\psi \rightarrow K_S K_L)$ obtained in Ref. [35], the value obtained in Ref. [18], and assuming either a positive or negative value for $\sin\varphi$.

$J/\psi \rightarrow K_S K_L$	φ	$\varphi(\kappa = 0)$	$\delta\mathcal{B}(J/\psi \rightarrow K^+K^-) \times 10^4$
BES [35]	$(97 \pm 5)^\circ$	$(98 \pm 4)^\circ$	0.13 ± 0.01
	$-(97 \pm 5)^\circ$	$-(96 \pm 4)^\circ$	-0.13 ± 0.01
Seth <i>et al.</i> [18]	$(111 \pm 5)^\circ$	$(108 \pm 4)^\circ$	0.15 ± 0.01
	$-(109 \pm 5)^\circ$	$-(107 \pm 4)^\circ$	-0.15 ± 0.01

TABLE VI. The relative phase (φ) between the single-photon and strong amplitudes for $\psi(2S) \rightarrow K\bar{K}$ decays calculated with $\kappa = A_\gamma^{K_S K_L}/A_\gamma^{K^+K^-} = 0.15 \pm 0.07$ and $\kappa = 0$, and the correction to the value of $\mathcal{B}(\psi(2S) \rightarrow K^+K^-)$ measured in the reaction $e^+e^- \rightarrow K^+K^-$. The calculation was performed for each sign of $\sin\varphi$.

φ	$\varphi(\kappa = 0)$	$\delta\mathcal{B}(\psi(2S) \rightarrow K^+K^-) \times 10^4$
$(82 \pm 12)^\circ$	$(92 \pm 9)^\circ$	0.11 ± 0.01
$-(58 \pm 14)^\circ$	$-(57 \pm 12)^\circ$	-0.10 ± 0.02

seen that allowing $A_\gamma^{K_S K_L}$ to be nonzero does not lead to a significant change in the results.

For the J/ψ , for which the most precise measurement of $\mathcal{B}(J/\psi \rightarrow K^+K^-)$ was performed in $\psi(2S)$ decay, the result for $\cos\varphi$ is weakly dependent on the sign of $\sin\varphi$. We confirm the conclusion of Refs. [16,17] to the effect that the strong amplitude describing $J/\psi \rightarrow K^+K^-$ decay has a large imaginary part. Using $\mathcal{B}(J/\psi \rightarrow K_S K_L)$ from Ref. [18], a non-negligible real part of the strong amplitude is obtained.

For the $\psi(2S)$, the result on $\cos\varphi$ is strongly dependent on the sign of $\sin\varphi$. Here theoretical arguments may help to choose the sign. The ratio of the strong amplitudes for $\psi(2S) \rightarrow K\bar{K}$ and $J/\psi \rightarrow K\bar{K}$ decays is expected [38] to be

$$\begin{aligned} \frac{A_s^2(\psi(2S) \rightarrow K\bar{K})}{A_s^2(J/\psi \rightarrow K\bar{K})} &\approx \frac{\mathcal{B}(\psi(2S) \rightarrow e^+e^-) \beta^3(M_{\psi(2S)})}{\mathcal{B}(J/\psi \rightarrow e^+e^-) \beta^3(M_{J/\psi})} \\ &= 0.138 \pm 0.003. \end{aligned} \quad (12)$$

The experimental value of this ratio (for J/ψ we used A_s obtained with $\mathcal{B}(J/\psi \rightarrow K_S K_L)$ from Ref. [18]) is 0.192 ± 0.026 for $\sin\varphi < 0$ and 0.170 ± 0.023 for $\sin\varphi > 0$. The result for the positive sign is in slightly better agreement with the prediction.

Using the values of $\delta\mathcal{B}$ given in Tables V and VI, we correct the measured values of the products $\Gamma(\psi \rightarrow e^+e^-)\mathcal{B}(\psi \rightarrow K^+K^-)$ and the branching fractions and list the corrected values in Table IV.

The corrected values of $\mathcal{B}(J/\psi \rightarrow K^+K^-)$ can be compared with the measurement of Ref. [18] $(2.86 \pm 0.21) \times 10^{-4}$. The difference between the two measurements is 2σ for $\sin\varphi < 0$, and 1σ for $\sin\varphi > 0$. Our result for $J/\psi \rightarrow K^+K^-$ thus provides an indication that $\sin\varphi$ is positive. It should be stressed that the shifts we find between the measured and true values of $\mathcal{B}(\psi \rightarrow K^+K^-)$ are significant: about 5% for the J/ψ and about 15% for the $\psi(2S)$. Thus, the interference effect should be taken into account in any future precise measurements of the branching fraction for $J/\psi \rightarrow K^+K^-$ or $\psi(2S) \rightarrow K^+K^-$.

VIII. SUMMARY

The process $e^+e^- \rightarrow K^+K^-\gamma$ has been studied in the K^+K^- invariant mass range from 2.6 to 8 GeV/ c^2 using events in which the photon is emitted close to the collision axis. From the measured K^+K^- mass spectrum we obtain the $e^+e^- \rightarrow K^+K^-$ cross section and determine the charged-kaon electromagnetic form factor (Table III). This is the first measurement of the kaon form factor for K^+K^- invariant masses higher than 5 GeV/ c^2 and the most precise measurement in the range 2.6–5.0 GeV/ c^2 . Our data indicate clearly that the difference between the measured form factor and the leading twist pQCD prediction decreases with increasing K^+K^- invariant mass.

We present measurements of the $J/\psi \rightarrow K^+K^-$ and $\psi(2S) \rightarrow K^+K^-$ branching fractions (Table IV). Using the measured values of the branching fractions and charged-kaon form factors, and data from other experiments on $e^+e^- \rightarrow K_S K_L$ and $\psi \rightarrow K\bar{K}$ decays, we have determined the phase difference φ between the strong and single-photon amplitudes for $J/\psi \rightarrow K\bar{K}$ and $\psi(2S) \rightarrow K\bar{K}$ decays. We have calculated the shifts in the measured values of the branching fractions due to interference between resonant and nonresonant amplitudes in the $e^+e^- \rightarrow K^+K^-$ reaction. The shift has been found to be relatively large, about $\pm 5\%$ for the J/ψ and about $\pm 15\%$ for the $\psi(2S)$, where the sign is determined by the sign of $\sin \varphi$.

It should be noted that the sign of $\sin \varphi$ for J/ψ decays can be determined experimentally from the difference, $\delta\mathcal{B}/\mathcal{B}$, between the $J/\psi \rightarrow K^+K^-$ branching fractions

measured in the reaction $e^+e^- \rightarrow K^+K^-$ and in the decay $\psi(2S) \rightarrow J/\psi\pi^+\pi^-$. We hope that this measurement will be performed in future experiments.

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