

1 **RATIONAL APPROXIMATION TO THE FRACTIONAL**
2 **LAPLACIAN OPERATOR IN REACTION-DIFFUSION PROBLEMS***

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4 **Abstract.** This paper provides a new numerical strategy to solve fractional in space reaction-
5 diffusion equations on bounded domains under homogeneous Dirichlet boundary conditions. Using
6 the matrix transfer technique the fractional Laplacian operator is replaced by a matrix which, in
7 general, is dense. The approach here presented is based on the approximation of this matrix by the
8 product of two suitable banded matrices. This leads to a semi-linear initial value problem in which
9 the matrices involved are sparse. Numerical results are presented to verify the effectiveness of the
10 proposed solution strategy.

11 **Key words.** fractional Laplacian operator, matrix functions, Gauss-Jacobi rule

12 **AMS subject classifications.** 65F60, 35R11, 65D32

13 **1. Introduction.** Fractional-order in space mathematical models, in which an
14 integer-order differential operator is replaced by a corresponding fractional one, are
15 becoming increasingly used since they provide an adequate description of many pro-
16 cesses that exhibit anomalous diffusion. This is due to the fact that the non-local
17 nature of the fractional operators **enables** to capture the spatial heterogeneity that
18 characterizes these processes.

19 There are however some challenges when facing fractional models. First of all,
20 there is no unique way to define fractional in space derivatives and, in general, these
21 definitions are not equivalent especially when more than one spatial dimension is
22 considered [17]. In addition, considering that the value of the solution at a given
23 point depends on the solution behavior on the entire domain, it is intuitive to un-
24 derstand that the boundary conditions deserve a particular attention and should be
25 appropriately chosen in order to model the phenomenon properly.

26 In this paper we consider the following fractional in space reaction-diffusion dif-
27 ferential equation

28 (1)
$$\frac{\partial u(\mathbf{x}, t)}{\partial t} = -\kappa_\alpha (-\Delta)^{\alpha/2} u(\mathbf{x}, t) + f(\mathbf{x}, t, u), \quad \mathbf{x} \in \Omega \subset \mathbb{R}^n, \quad t \in (0, T),$$

29 subject to homogeneous Dirichlet boundary conditions

30 (2)
$$u(\mathbf{x}, t)|_{\hat{\Omega}} = 0, \quad \hat{\Omega} = \mathbb{R}^n \setminus \Omega,$$

31 and the initial condition

32 (3)
$$u(\mathbf{x}, 0) = u_0(\mathbf{x}),$$

33 where κ_α represents the diffusion coefficient and the forcing term f and u_0 are suf-
34 ficiently smooth functions. The symmetric space fractional derivative $-(-\Delta)^{\alpha/2}$ of
35 order α ($1 < \alpha \leq 2$) is defined through the spectral decomposition of the homogeneous
36 Dirichlet Laplace operator $(-\Delta)$, [9, Definition 2]. Assuming that Ω is a Lipschitz

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37 domain, the spectrum of $(-\Delta)$ is discrete and positive, and accumulate at infinity.
 38 Thus,

$$39 \quad (4) \quad -(-\Delta)^{\alpha/2}u = \sum_{s=1}^{\infty} \mu_s^{\alpha/2} c_s \varphi_s,$$

40 where $c_s = \int_{\Omega} u \varphi_s$ are the Fourier coefficients of u , and $\{\mu_s\}$, $\{\varphi_s\}$ are the eigenvalues
 41 and the eigenvectors of $(-\Delta)$, respectively.

42 We remark that the fractional power of the Laplace operator is alternatively defined
 43 in the literature using the Fourier transform on an infinite domain [12], with
 44 a natural extension to finite domain when the function u vanishes on and outside
 45 the boundary of the spatial domain. In this case, in fact, it is possible to consider
 46 non-local problems on bounded domains by simply assuming that the solution of frac-
 47 tional problem is equal to zero everywhere outside the domain of interest. Using such
 48 definition and assuming to work with homogeneous Dirichlet boundary conditions,
 49 in [15, Lemma 1] it has been proved that the one-dimensional fractional Laplacian
 50 operator $-(-\Delta)^{\alpha/2}$ as defined in (4) is equivalent to the Riesz fractional derivative.
 51 Hence, it can be approximated by the ‘fractional centered derivative’ introduced by
 52 Ortigueira in [11]. Çelik and Duman in [5] have used such a method for solving a
 53 fractional diffusion equation with the Riesz fractional derivative in a finite domain.
 54 Moreover, by exploiting the decay of the coefficients characterizing the method, in
 55 [10] a ‘short memory’ version of the scheme has been implemented. However, both
 56 the original and the modified methods only work for one-dimensional problems.

57 A mainstay in the numerical treatment of partial differential problems of type
 58 (1)–(3) is to apply the method of lines. Discretizing in space with a uniform mesh of
 59 stepsize h in each domain direction and using the matrix transfer technique proposed
 60 in [8, 9] by Ilić et al., we obtain

$$61 \quad -(-\Delta)^{\alpha/2}u \approx -\frac{1}{h^{\alpha}} L^{\alpha/2}u,$$

62 where $h^{-2}L$ is the approximate matrix representation of the standard Laplacian ob-
 63 tained by using any finite difference method. Consequently, (1) is transformed into a
 64 system of ordinary differential equations

$$65 \quad (5) \quad \frac{d\mathbf{u}}{dt} = -\frac{\kappa_{\alpha}}{h^{\alpha}} L^{\alpha/2}\mathbf{u} + \mathbf{f}(t, \mathbf{u}),$$

66 where \mathbf{u} and $\mathbf{f}(t, \mathbf{u})$ denote the vectors of node values of u and f , respectively. The
 67 matrix L raised to the fractional power $\alpha/2$ is, in general, a dense matrix which could
 68 be also very large depending on the numbers of mesh points used for the spatial
 69 discretization. Therefore, the computational effort for solving (5) could be really
 70 heavy, independently of the integrator used. Recently, some authors have developed
 71 techniques for reducing this cost. In particular, an approach which can be equally
 72 applicable to fractional-in-space problems in two or three spatial dimensions has been
 73 considered in [4]. The key point of this approach is the efficient computation of the
 74 fractional power of a matrix times a vector.

75 In this paper, we propose a solution strategy based on a suitable approximation
 76 of $L^{\alpha/2}$. In particular, we look for a decomposition of the type

$$77 \quad L^{\alpha/2} \approx M^{-1}K,$$

78 where M and K are both banded matrices arising from a rational approximation of the
 79 function $z^{\alpha/2-1}$, based on the Gauss-Jacobi rule applied to the integral representation
 80 of $L^{\alpha/2}$, cf. [6]. The poles of the formula depends on a continuous parameter whose
 81 choice is crucial for a fast and accurate approximation. The above factorization **allows**
 82 to approximate the solution of (5) by solving

$$83 \quad (6) \quad M \frac{d\mathbf{v}}{dt} = -\frac{\kappa_\alpha}{h^\alpha} K\mathbf{v} + M\mathbf{f}(t, \mathbf{v}).$$

84 By virtue of the structure of the matrices M and K the numerical solution of (6)
 85 may be computed in a more efficient way with respect to the one of (5). We remark
 86 that the approach is independent of the Laplacian working dimension.
 87

88 The paper is organized as follows. In **section 2**, the main results about the matrix
 89 transfer technique are recalled. **Section 3** is devoted to the construction of the rational
 90 approximation together with the analysis of the asymptotically optimal choices of the
 91 poles. In **section 4** a theoretical error analysis is presented. Numerical experiments
 92 are carried out in **section 5**, and the conclusions follow in **section 6**.

93 **2. Background on the matrix transfer technique.** For an independent
 94 reading, in this section the basic facts concerning the matrix transfer technique pro-
 95 posed by Ilić et al. in [8, 9] to discretize the one-dimensional fractional Laplacian
 96 operator are recalled. In addition, since in this work we also lead with problems in
 97 two spatial dimensions, we refer to the results given in [16] as well.
 98

99 Working with the basic assumption that the fractional Laplacian operator with
 100 Dirichlet boundary conditions can be defined as the fractional power of the standard
 101 Laplacian, the matrix transfer technique simply consists in approximating the oper-
 102 ator $(-\Delta)^{\alpha/2}$ through the matrix $-h^{-\alpha}L^{\alpha/2}$, where $h^{-2}L$ is any finite-difference
 103 approximation of $(-\Delta)$ on a uniform mesh of size h . The only important requirement
 104 is that the matrix L is positive definite so that its fractional power is well defined. This
 105 requirement is fulfilled by the existing standard central difference schemes. Working
 106 like that, the original problem (1)–(3) is then transformed into the semi-linear initial
 107 value problem

$$108 \quad (7) \quad \begin{aligned} \frac{d\mathbf{u}}{dt} &= -\frac{\kappa_\alpha}{h^\alpha} L^{\alpha/2}\mathbf{u} + \mathbf{f}(t, \mathbf{u}), & t \in (0, T), \\ \mathbf{u}(0) &= \mathbf{u}_0, \end{aligned}$$

110 where \mathbf{u}_0 denotes the vector of node values of u_0 .

111 It is important to remark that while L is typically sparse, when $\alpha \neq 2$, the ma-
 112 trix $L^{\alpha/2}$ loses its sparsity and becomes dense. Observe moreover that the stiffness
 113 property of (7) for $\alpha = 2$ is essentially inherited by the fractional counterpart so that
 114 an implicit scheme or an exponential integrator is generally needed for solving this
 115 initial value problem. In both cases the density of $L^{\alpha/2}$ may lead to a computational
 116 demanding integrator when the discretization is sharp. In order to overcome the limi-
 117 tations in terms of computational efficiency, we propose a strategy based on a suitable
 118 approximate factorization of $L^{\alpha/2}$. In the next section we focus on the construction
 119 of such approximation.

120 **3. Approximation to the matrix fractional power.** From the theory of
 121 matrix functions (see [7] for a survey) we know that the fractional power of a generic

122 matrix A can be written as a contour integral

$$123 \quad A^\beta = \frac{A}{2\pi i} \int_{\Gamma} z^{\beta-1} (zI - A)^{-1} dz,$$

124 where Γ is a suitable closed contour enclosing the spectrum of A , $\sigma(A)$, in its interior.
 125 The following known result (see, e.g., [2]) expresses A^β in terms of a real integral.
 126 The proof is based on a particular choice of Γ and a subsequent change of variable.

127 **PROPOSITION 1.** *Let $A \in \mathbb{R}^{m \times m}$ be such that $\sigma(A) \subset \mathbb{C} \setminus (-\infty, 0]$. For $0 < \beta < 1$
 128 the following representation holds*

$$129 \quad A^\beta = \frac{A \sin(\beta\pi)}{\beta\pi} \int_0^\infty (\rho^{1/\beta} I + A)^{-1} d\rho.$$

130 In order to confine the dependence of β to a weight function, we consider the
 131 change of variable

$$132 \quad (8) \quad \rho^{1/\beta} = \tau \frac{1-t}{1+t}, \quad \tau > 0,$$

133 which yields

$$\begin{aligned} 134 \quad & \frac{1}{\beta} \int_0^\infty (\rho^{1/\beta} I + A)^{-1} d\rho \\ 135 \quad & = 2 \int_{-1}^1 \left(\tau \frac{1-t}{1+t} \right)^{\beta-1} \left(\tau \frac{1-t}{1+t} I + A \right)^{-1} \frac{\tau}{(1+t)^2} dt \\ 136 \quad & = 2\tau^\beta \int_{-1}^1 (1-t)^{\beta-1} (1+t)^{-\beta} (\tau(1-t)I + (1+t)A)^{-1} dt, \\ 137 \end{aligned}$$

138 and hence

$$139 \quad (9) \quad A^\beta = A \frac{\sin(\beta\pi)}{\pi} 2\tau^\beta \int_{-1}^1 (1-t)^{\beta-1} (1+t)^{-\beta} (\tau(1-t)I + (1+t)A)^{-1} dt.$$

140 The above formula naturally suggests the use of the k -point Gauss-Jacobi rule and
 141 consequently a rational approximation of the type

$$142 \quad (10) \quad A^\beta \approx R_k(A) := A \sum_{j=1}^k \gamma_j (\eta_j I + A)^{-1},$$

143 where the coefficients γ_j and η_j are given by

$$144 \quad \gamma_j = \frac{2 \sin(\beta\pi) \tau^\beta}{\pi} \frac{w_j}{1 + \vartheta_j}, \quad \eta_j = \frac{\tau(1 - \vartheta_j)}{1 + \vartheta_j},$$

145 in which w_j and ϑ_j are, respectively, the weights and nodes of the Gauss-Jacobi
 146 quadrature rule with weight function $(1-t)^{\beta-1}(1+t)^{-\beta}$. Of course, the above ap-
 147 proximation can be used in our case with $\beta = \alpha/2$ whenever $A = h^{-2}L$ represents
 148 the discrete Laplacian operator with Dirichlet boundary conditions, whose spectrum
 149 is contained in \mathbb{R}^+ . At this point, denoting by $z P_{k-1}(z)$ and $Q_k(z)$ the polynomials
 150 of degree k such that $R_k(z) = (z P_{k-1}(z))/Q_k(z)$, we can approximate the solution
 151 of (7) by solving (6) with $M = Q_k(L)$ and $K = LP_{k-1}(L)$. We remark that the use

152 of the Gauss-Jacobi rule ensures that $\gamma_j > 0$ and $\eta_j > 0$ for each j , and hence it is
 153 immediate to verify that the spectrum of $R_k(L)$ is strictly contained in the positive
 154 real axis. This condition is fundamental to preserve the stability properties of (7)
 155 whenever $L^{\alpha/2}$ is replaced by $R_k(L)$.

156 **We need to mention that in the field of fractional calculus the approximation (10)**
 157 **has already been used in [1] for the approximation of the Caputo's fractional deriva-**
 158 **tive. Here, however, the definition of τ in (8), and the subsequent error analysis will**
 159 **be completely different because of the spectral properties of the Laplacian operator**
 160 **with respect to the ones of the first-derivative.**

161 **3.1. Choice of τ .** The choice of the parameter τ in the change of variable (8)
 162 is crucial for the quality of the approximation attainable by (10). Assuming that the
 163 generic matrix A is symmetric and positive definite, let λ_{\min} and λ_{\max} be its smallest
 164 and largest eigenvalue, respectively. Let moreover $\Lambda = [\lambda_{\min}, \lambda_{\max}]$. It is well known
 165 that

166 (11)
$$\|A^\beta - R_k(A)\|_2 \leq \max_{\lambda \in \Lambda} |\lambda^\beta - R_k(\lambda)|.$$

167 In this view, looking at (9), a good choice of τ is the one that minimizes, uniformly
 168 with respect to $\lambda \in \Lambda$, the error of the Gauss-Jacobi formula when applied to the
 169 computation of

170
$$\int_{-1}^1 (1-t)^{\beta-1} (1+t)^{-\beta} (\tau(1-t) + (1+t)\lambda)^{-1} dt, \quad \lambda \in \Lambda.$$

171 From the theory of best uniform polynomial approximation and its application to
 172 the analysis of the Gauss quadrature rules (see e.g. [14] for a recent study) it is known
 173 that the position of the poles of the integrand function with respect to the interval of
 174 integration defines the quality of the approximation. In our case, we observe that for
 175 each $\tau \in \Lambda$ the poles of the integrand function

176
$$f_{\tau,\lambda}(t) = (\tau(1-t) + (1+t)\lambda)^{-1},$$

177 are functions of λ defined by

178
$$p_\tau(\lambda) = \frac{\tau + \lambda}{\tau - \lambda},$$

179 and we clearly have $p_\tau(\lambda) > 1$ for $\lambda < \tau$, and $p_\tau(\lambda) < -1$ for $\lambda > \tau$. Our aim is to
 180 define τ in order to maximize the distance of the set

181
$$Q_\tau = \{p_\tau(\lambda), \lambda \in \Lambda\}$$

182 from the interval of integration $[-1, 1] \subset \mathbb{R} \setminus Q_\tau$. We observe that for $\lambda_{\min} < \tau < \lambda_{\max}$
 183 the worst case is given by $\lambda = \lambda_{\min}$ or $\lambda = \lambda_{\max}$ since we have respectively

184
$$\min_{\lambda \in \Lambda} \text{dist}(p_\tau(\lambda), [-1, 1]) = p_\tau(\lambda_{\min}) - 1,$$

185 or

186
$$\min_{\lambda \in \Lambda} \text{dist}(p_\tau(\lambda), [-1, 1]) = -p_\tau(\lambda_{\max}) - 1.$$

187 As consequence, the idea is to set τ such that

188
$$p_\tau(\lambda_{\min}) - 1 = -p_\tau(\lambda_{\max}) - 1,$$

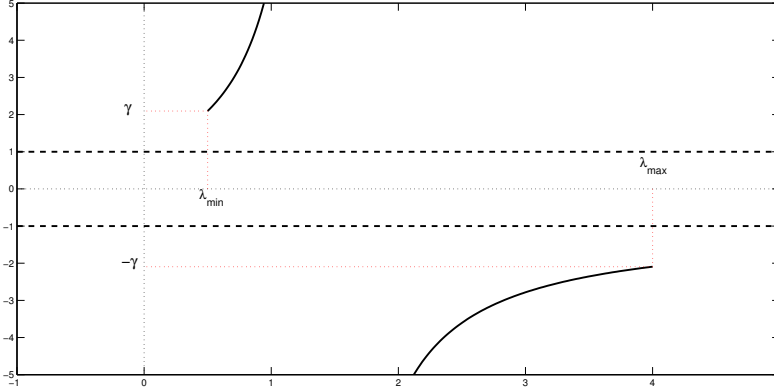


FIG. 1. Example of function $p_\tau(\lambda)$ for $\lambda_{\min} = 0.5$, $\lambda_{\max} = 4$. The choice of τ as in (12) ensures the symmetry of the set Q_τ . The minimum distance of the curve $p_\tau(\lambda)$ from the set $[-1, 1]$ is given by $\gamma - 1$ and is attained either in $\lambda = \lambda_{\min}$ or $\lambda = \lambda_{\max}$.

189 that leads directly to the equation

$$190 \quad \frac{\tau + \lambda_{\min}}{\tau - \lambda_{\min}} = -\frac{\tau + \lambda_{\max}}{\tau - \lambda_{\max}},$$

191 whose solution is

$$192 \quad (12) \quad \tau_{opt} = \sqrt{\lambda_{\min} \lambda_{\max}}.$$

193 Formally, τ_{opt} is given by

$$194 \quad \tau_{opt} = \arg \max_{\lambda_{\min} < \tau < \lambda_{\max}} \min_{\lambda \in \Lambda} |p_\tau(\lambda)|.$$

195 In this way, the set $Q_{\tau_{opt}}$ is symmetric with respect to the origin, that is $Q_{\tau_{opt}} =$
 196 $(-\infty, -\gamma) \cup (\gamma, +\infty)$, where

$$197 \quad (13) \quad \gamma = \frac{\sqrt{\kappa(A)} + 1}{\sqrt{\kappa(A)} - 1},$$

198 in which $\kappa(A)$ denotes the spectral condition number of A . This situation is summa-
 199 rized in an example reported in [Figure 1](#).

200 **4. Error analysis.** In this section we analyze the error of the rational approx-
 201 imation (10) with the choice of $\tau = \tau_{opt}$ in (8). We start with the following result,
 202 whose proof is given in [14, Theorems 4.3 and 4.4].

203 **THEOREM 2.** *Let $g(z)$ be a function analytic in an open subset of the complex*
 204 *plane containing the ellipse*

$$205 \quad \Gamma_\rho = \left\{ z = \frac{1}{2} \left(\rho e^{i\theta} + \frac{1}{\rho e^{i\theta}} \right), \rho > 1, \theta \in [0, 2\pi) \right\}.$$

206 *Let moreover $p_k^*[g]$ be the polynomial of degree $\leq k$ of best uniform approximation of*
 207 *g in $[-1, 1]$ and*

$$208 \quad E_k^*[g] = \max_{t \in [-1, 1]} |g(t) - p_k^*[g](t)|.$$

209 Then

$$210 \quad (14) \quad E_k^*[g] \leq \frac{2M(\rho)}{(\rho-1)\rho^k},$$

211 where

$$212 \quad M(\rho) = \max_{z \in \Gamma_\rho} |g(z)|.$$

213 **THEOREM 3.** *Let A be a symmetric positive definite matrix and $0 < \beta < 1$. Then*
 214 *for k sufficiently large, the error of the rational approximation (10), generated by the*
 215 *Gauss-Jacobi rule applied to the integral (9) for $\tau = \tau_{opt}$, is given by*

$$216 \quad \|A^\beta - R_k(A)\|_2 \leq C \|A\|_2 \tau^\beta \frac{(\rho_M + 1)}{(\rho_M - 1)(\rho_M - \gamma)} \frac{k}{\rho_M^{2k}},$$

217 where C is a constant independent of k , and

$$218 \quad \rho_M = \gamma + \sqrt{\gamma^2 - 1}.$$

219 *Proof.* For $\lambda \in \Lambda$ let

$$220 \quad f_\lambda(t) = (\tau_{opt}(1-t) + (1+t)\lambda)^{-1},$$

221 and

$$222 \quad I(f_\lambda) = \int_{-1}^1 (1-t)^{\beta-1} (1+t)^{-\beta} f_\lambda(t) dt.$$

223 Let moreover $I_k(f_\lambda)$ be the corresponding k -point Gauss-Jacobi approximation with
 224 weights w_j , $j = 1, \dots, k$. By standard arguments we have that

$$225 \quad |I(f_\lambda) - I_k(f_\lambda)| \leq |I(f_\lambda - p_{2k-1}^*[f_\lambda])| + |I_k(f_\lambda - p_{2k-1}^*[f_\lambda])| \\ 226 \quad (15) \quad \leq 2C_\beta E_{2k-1}^*[f_\lambda],$$

228 where, since $w_j > 0$,

$$229 \quad C_\beta = \sum_{j=1}^k |w_j| = \sum_{j=1}^k w_j = \int_{-1}^1 (1-t)^{\beta-1} (1+t)^{-\beta} dt.$$

230 Now, independently of $\lambda \in \Lambda$, the choice of $\tau = \tau_{opt}$ makes possible to use the bound
 231 (14) for each $1 < \rho < \rho_M$ where ρ_M solves

$$232 \quad \frac{1}{2} \left(\rho_M + \frac{1}{\rho_M} \right) = \gamma,$$

233 since $Q_{\tau_{opt}} = (-\infty, -\gamma) \cup (\gamma, +\infty)$. Thus by (15), (14) and using

$$234 \quad M(\rho) = \max_{z \in \Gamma_\rho} |f_\lambda(z)| \leq \frac{1}{\gamma - \frac{1}{2} \left(\rho + \frac{1}{\rho} \right)},$$

235 we obtain

$$236 \quad (16) \quad |I(f_\lambda) - I_k(f_\lambda)| \leq \frac{4C_\beta}{(\rho-1)\rho^{2k-1} \left(\gamma - \frac{1}{2} \left(\rho + \frac{1}{\rho} \right) \right)}, \quad 1 < \rho < \rho_M.$$

237 Now, neglecting the factor $1/(\rho - 1)$ and then minimizing with respect to ρ yields

$$238 \quad \bar{\rho} = \frac{2k-1}{2k} \left(\gamma + \sqrt{\gamma^2 - 1 + \frac{1}{(2k-1)^2}} \right)$$

$$239 \quad \approx \frac{2k-1}{2k} \rho_M =: \rho^*.$$

241 Hence, for k large enough (we need $\rho^* > 1$), we can use ρ^* in (16), obtaining

$$242 \quad (17) \quad |I(f_\lambda) - I_k(f_\lambda)| \leq \frac{8k\epsilon C_\beta (\rho_M + 1)}{(\rho_M - 1) \rho_M^{2k} (\rho_M - \gamma)}.$$

243 Indeed, defining k^* such that

$$244 \quad \frac{2k-1}{2k} \geq \frac{2}{\rho_M + 1} \quad \text{for } k \geq k^*$$

245 we have

$$246 \quad \frac{1}{\frac{2k-1}{2k} \rho_M - 1} \leq \frac{\rho_M + 1}{\rho_M - 1}.$$

247 Moreover, in (17) we have used the inequalities

$$248 \quad \frac{1}{\left(\frac{2k-1}{2k} \rho_M\right)^{2k-1}} \leq \frac{e}{\rho_M^{2k-1}},$$

$$249 \quad \gamma - \frac{1}{2} \left(\frac{2k-1}{2k} \rho_M + \frac{2k}{2k-1} \frac{1}{\rho_M} \right) \geq \frac{\rho_M (\rho_M - \gamma)}{2k}.$$

251 Finally, since by (9)

$$252 \quad \|A^\beta - R_k(A)\|_2 \leq \frac{\|A\|_2 \sin(\beta\pi)}{\pi} 2\tau^\beta \max_{\lambda \in \Lambda} |I(f_\lambda) - I_k(f_\lambda)|,$$

253 using (17) we obtain the result. \square

254 **COROLLARY 4.** *The asymptotic convergence factor **fulfills***

$$255 \quad \overline{\lim}_{k \rightarrow \infty} \|A^\beta - R_k(A)\|_2^{1/k} \leq \left(\frac{\sqrt[4]{\kappa(A)} - 1}{\sqrt[4]{\kappa(A)} + 1} \right)^2.$$

256 *Proof.* By (13)

$$257 \quad \rho_M = \gamma + \sqrt{\gamma^2 - 1} = \frac{\sqrt[4]{\kappa(A)} + 1}{\sqrt[4]{\kappa(A)} - 1}. \quad \square$$

258 **REMARK 5.** *From the above analysis it is easy to observe that for the Laplacian*
 259 *operator $h^{-2}L$, discretized with standard central differences (3-points or 5-points in*
 260 *one or two dimensions, respectively), we have*

$$261 \quad \left(\frac{\sqrt[4]{\kappa(L)} + 1}{\sqrt[4]{\kappa(L)} - 1} \right)^2 \approx 1 + \frac{2\pi}{N},$$

262 *where N represents the number of discretization points in one dimension.*

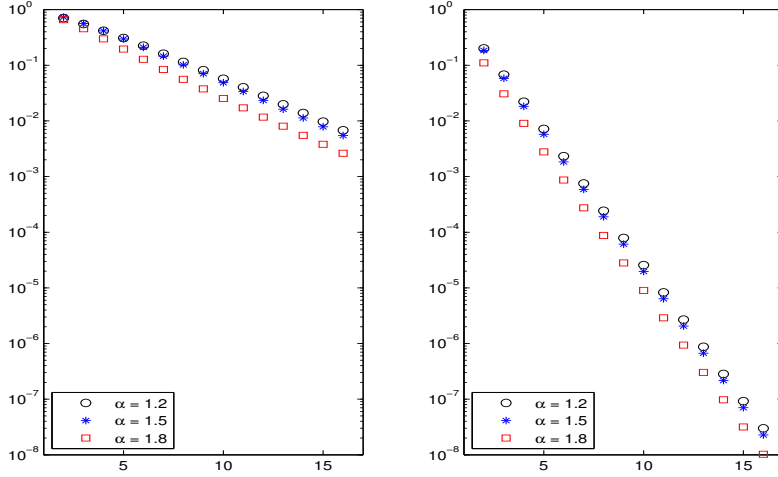


FIG. 2. Relative error of the rational approximation versus k , the number of points of the Gauss-Jacobi rule, for some values of α . The one- and the two-dimensional cases are on the left and on the right, respectively. In the first case the dimension of the problem is 200 and in the second one it is 400.

263 In Figure 2 we plot the relative error for the one- and two-dimensional Laplacian
 264 discretized as in the previous remark for some values of α . The geometric conver-
 265 gence theoretically proved in this section is clear in the pictures, together with the
 266 substantial independence of α , which is absorbed by the weight function. It is also
 267 quite evident that the method is particularly effective for the two-dimensional case;
 268 this represents an important feature since most of the standard techniques for the
 269 discretization of the fractional Laplacian only works in one dimension.

270 **5. Solving fractional in space reaction-diffusion problems.** As already
 271 said in Section 2, if we discretize on a uniform mesh the fractional Laplacian operator
 272 occurring in (1), we obtain the initial value problem

273 (18)
$$\frac{d\mathbf{u}}{dt} = -\frac{\kappa_\alpha}{h^\alpha} L^{\alpha/2} \mathbf{u} + \mathbf{f}(t, \mathbf{u}), \quad \mathbf{u}(0) = \mathbf{u}_0.$$

274

275 Therefore, the application of the rational approximation (10) of $L^{\alpha/2}$, based on the
 276 k -point Gauss-Jacobi rule and given by $R_k(L) \equiv M^{-1}K$, leads to the following initial
 277 value problem

278 (19)
$$\frac{d\mathbf{v}}{dt} = -\frac{\kappa_\alpha}{h^\alpha} R_k(L) \mathbf{v} + \mathbf{f}(t, \mathbf{v}), \quad \mathbf{v}(0) = \mathbf{u}_0.$$

279

280 Denoting by

281
$$E_k = R_k(L) - L^{\alpha/2}$$

282 the error of the rational approximation, (19) can be equivalently written as

283 (20)
$$\frac{d\mathbf{v}}{dt} = -\frac{\kappa_\alpha}{h^\alpha} \left(L^{\alpha/2} + E_k \right) \mathbf{v} + \mathbf{f}(t, \mathbf{v}), \quad \mathbf{v}(0) = \mathbf{u}_0.$$

284

285 The solution $\mathbf{v}(t)$ is therefore the solution of the perturbed version of (18). Setting

286
$$g(t, \mathbf{u}) = -\frac{\kappa_\alpha}{h^\alpha} L^{\alpha/2} \mathbf{u} + \mathbf{f}(t, \mathbf{u}),$$

287 we assume that g satisfies the one-sided Lipschitz condition

$$288 \quad (21) \quad \langle g(t, \mathbf{u}_1) - g(t, \mathbf{u}_2), \mathbf{u}_1 - \mathbf{u}_2 \rangle \leq m \|\mathbf{u}_1 - \mathbf{u}_2\|_2^2, \quad m < 0, \quad \mathbf{u}_1, \mathbf{u}_2 \in \mathbb{R}^N,$$

289 where $\langle \cdot, \cdot \rangle$ denotes the scalar product corresponding to the Euclidean norm $\|\cdot\|_2$. We
 290 are interested in an a priori upper bound on the distance of \mathbf{u} and \mathbf{v} . The difference
 291 $\mathbf{w}(t) = \mathbf{u}(t) - \mathbf{v}(t)$ is the solution of the initial value problem

$$292 \quad \frac{d\mathbf{w}}{dt} = g(t, \mathbf{u}) - g(t, \mathbf{v}) + \frac{\kappa_\alpha}{h^\alpha} E_k \mathbf{v}, \quad t \in (0, T),$$

$$293 \quad \mathbf{w}(0) = \mathbf{0}.$$

295 Calculating the scalar product of $\mathbf{w}(t)$ with both sides of the differential equation, we
 296 find for the left-hand side

$$297 \quad \left\langle \frac{d\mathbf{w}(t)}{dt}, \mathbf{w}(t) \right\rangle = \frac{1}{2} \frac{d}{dt} \langle \mathbf{w}(t), \mathbf{w}(t) \rangle = \frac{1}{2} \frac{d}{dt} \|\mathbf{w}(t)\|_2^2 = \|\mathbf{w}(t)\|_2 \frac{d}{dt} \|\mathbf{w}(t)\|_2.$$

298 From (21) and Schwarz' inequality we obtain for the right-hand side

$$299 \quad \left\langle g(t, \mathbf{u}) - g(t, \mathbf{v}) + \frac{\kappa_\alpha}{h^\alpha} E_k \mathbf{v}, \mathbf{w}(t) \right\rangle = \langle g(t, \mathbf{u}) - g(t, \mathbf{v}), \mathbf{w}(t) \rangle + \frac{\kappa_\alpha}{h^\alpha} \langle E_k \mathbf{v}, \mathbf{w}(t) \rangle$$

$$300 \quad \leq m \|\mathbf{w}\|_2^2 + \frac{\kappa_\alpha}{h^\alpha} \|E_k \mathbf{v}\|_2 \|\mathbf{w}\|_2.$$

302 Combining these results we find for the scalar function $\|\mathbf{w}(t)\|_2$:

$$303 \quad \frac{d}{dt} \|\mathbf{w}(t)\|_2 \leq m \|\mathbf{w}(t)\|_2 + \frac{\kappa_\alpha}{h^\alpha} \|E_k \mathbf{v}(t)\|_2.$$

304 Now, since $\mathbf{w}(0) = \mathbf{0}$, by the Grönwall inequality we obtain

$$305 \quad (22) \quad \|\mathbf{w}(t)\|_2 \leq \frac{\kappa_\alpha}{h^\alpha} \|E_k\|_2 \int_0^t e^{m(t-s)} \|\mathbf{v}(s)\|_2 ds.$$

306 This immediately provides the proof of the following result.

307 **PROPOSITION 6.** *If \mathbf{u} and \mathbf{v} are solutions of (18) and (19), respectively, and m*
 308 *is the one-sided Lipschitz constant introduced in (21), assuming that $\|\mathbf{v}(s)\|_2 \leq c$ for*
 309 *each $s \in [0, t]$, then*

$$310 \quad (23) \quad \|\mathbf{u}(t) - \mathbf{v}(t)\|_2 \leq \frac{\kappa_\alpha}{h^\alpha} \|E_k\|_2 \frac{c}{m} (e^{mt} - 1).$$

311 A more sharper condition estimate may be derived if also the right-hand side of (20)
 312 satisfies a one-sided Lipschitz condition.

313 **PROPOSITION 7.** *Assume that there exists $\bar{m} < 0$ such that for each $\mathbf{u}_1, \mathbf{u}_2 \in \mathbb{R}^N$*

$$314 \quad (24) \quad \left\langle g(t, \mathbf{u}_1) - g(t, \mathbf{u}_2) - \frac{\kappa_\alpha}{h^\alpha} E_k (\mathbf{u}_1 - \mathbf{u}_2), \mathbf{u}_1 - \mathbf{u}_2 \right\rangle \leq \bar{m} \|\mathbf{u}_1 - \mathbf{u}_2\|_2^2.$$

315 *Moreover, assume that the zero solution is an equilibrium point for (18), that is,*
 316 *$\mathbf{f}(t, \mathbf{0}) = \mathbf{0}$. Then there exists $\tilde{m} < 0$ such that*

$$317 \quad (25) \quad \|\mathbf{u}(t) - \mathbf{v}(t)\|_2 \leq \frac{\kappa_\alpha}{h^\alpha} \|E_k\|_2 \|\mathbf{v}(0)\|_2 t e^{\tilde{m}t}.$$

318 *Proof.* Condition (24) ensures that for each $\mathbf{z}(t)$, solution of (19) with initial
 319 condition $\mathbf{z}(0) = \mathbf{z}_0$, we have

320
$$\|\mathbf{v}(t) - \mathbf{z}(t)\|_2 \leq e^{\tilde{m}t} \|\mathbf{v}(0) - \mathbf{z}(0)\|_2.$$

322 Since we have assumed that $\mathbf{z} \equiv \mathbf{0}$ is a solution, we have

323
$$\|\mathbf{v}(t)\|_2 \leq e^{\tilde{m}t} \|\mathbf{v}(0)\|_2.$$

325 Now, let $\tilde{m} := \max(m, \bar{m})$. Replacing m with \tilde{m} in (22) and \bar{m} with \tilde{m} in the above
 326 inequality, (25) immediately follows. \square

327 **5.1. Choice of k .** The above propositions can easily be adopted to select a value
 328 of k that allows to keep the error below a given tolerance. Indeed, for using (23) one
 329 can consider the approximation

330 (26)
$$m \approx -\frac{k_\alpha}{h^\alpha} (\lambda_{\min}(L))^{\alpha/2}$$

331 and define $c := \|\mathbf{v}(0)\|_2$. If $\mathbf{v}(0) = \mathbf{0}$ some information about \mathbf{v} can be obtained
 332 working in small dimension. The approximation (26) can also be used for \tilde{m} whenever
 333 it is possible to employ (25). Finally, since (see (11))

334
$$\|E_k\|_2 \leq \left| R_k(\lambda_{\min}(L)) - (\lambda_{\min}(L))^{\alpha/2} \right|,$$

335 working scalarly we can easily select k such that

336
$$\|\mathbf{u}(t) - \mathbf{v}(t)\|_2 \leq C \text{tol},$$

337 where tol is a given tolerance and the constant C plays the role of a scaling factor
 338 that should be set equal to c or $\|\mathbf{v}(0)\|_2$ when using (23) or (25), respectively (this
 339 to avoid the dependence of k on the magnitude of the solution). Anyway, we remark
 340 that this procedure may provide a conservative value of k , because of the bounds used
 341 for obtaining the estimate.

342 **5.2. Numerical examples.** In this subsection, we provide the results of some
 343 numerical experiments we have conducted to prove the effectiveness of the proposed
 344 approach with respect to the matrix transfer technique. In order to take the advantage
 345 in terms of computational work and memory saving of the sparse structure of the
 346 matrices M and K , instead of (19), we solve the equivalent initial value problem

347 (27)
$$M \frac{d\mathbf{v}}{dt} = -\frac{\kappa_\alpha}{h^\alpha} K\mathbf{v} + M\mathbf{f}(t, \mathbf{v}), \quad \mathbf{v}(0) = \mathbf{u}_0.$$

349 In particular, we first focus on a fractional in space diffusion equation. Then, we con-
 350 sider two reaction-diffusion equations; in the second one the forcing term is dependent
 351 on the solution and we choose an initial data that does not satisfy the boundary condi-
 352 tions. All of these examples are in one spatial dimension. In each case, discretizing the
 353 spatial domain $\Omega = (a, b)$ with a uniform mesh having stepsize $h = (b - a)/(N + 1)$,
 354 we consider the standard 3-points central difference discretization of the Laplacian
 355 $h^{-2}L = h^{-2}\text{tridiag}(-1, 2, -1) \in \mathbb{R}^{N \times N}$. Finally, we also report the results obtained
 356 by applying our approach for the numerical solution of a fractional reaction-diffusion
 357 example in two space dimensions. In this case, we discretize in space the problem via
 358 the 5-points finite difference stencil. The matrix L is therefore a block tridiagonal

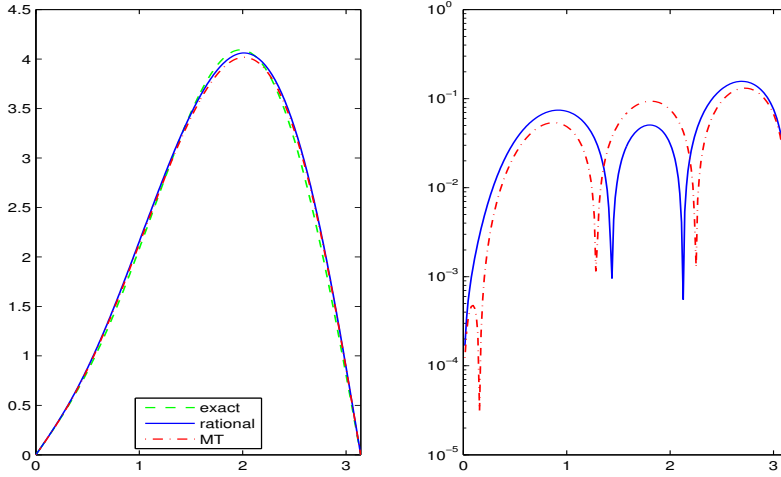


FIG. 3. Comparison of the analytic solution of the problem in *Example 1* with the numerical solutions provided by the rational approach and the matrix transfer technique at $t = 0.4$ (left) and corresponding errors (right).

359 matrix of size N^2 having the following form $L = \text{tridiag}(-I, B, -I)$, with I denoting
 360 the identity matrix of size N and $B = \text{tridiag}(-1, 4, -1) \in \mathbb{R}^{N \times N}$.

361

362 In all examples, we solve (18) and (27) by the MATLAB routine `ode15s`. More-
 363 over, we indicate by ‘*exact*’ the analytical solution, by ‘*MT*’ the solution of the
 364 problem (18), obtained by applying the matrix transfer technique, and by ‘*rational*’
 365 the solution arising from (27).

366 **EXAMPLE 1.** Consider the problem (1) on the spatial domain $\Omega = (0, \pi)$, with
 367 $\kappa_\alpha = 0.25$ and $f = 0$. According to [15, Section 3.1], the analytic solution correspond-
 368 ing to the initial condition $u_0(x) = x^2(\pi - x)$ is given by

$$369 \quad u(x, t) = \sum_{n=1}^{\infty} \frac{8(-1)^{n+1} - 4}{n^3} \sin(nx) \exp(-\kappa_\alpha n^\alpha t).$$

370 Setting $\alpha = 1.8$, at time $t = 0.4$ in the left-hand side of *Figure 3* the exact solution
 371 is compared with the numerical solutions of the semi-discrete problems (18) and (27)
 372 with $h = \pi/201$ (that is $h = 0.0157$) and $k = 4$. On the right picture, the step-by-step
 373 maximum norm of the difference between the analytic solution and the numerical ones
 374 is reported. As one can see, the numerical solution provided by the rational approxi-
 375 mation is in good agreement with the one obtained by the matrix transfer technique.
 376 For this choice of t , the bounds (23) and (25) essentially coincide and an error com-
 377 parable to the one obtained experimentally is predicted by choosing $k = 20$. Of course,
 378 a smaller value of k would be obtained for $t > 0.4$ since the bound (25) rapidly goes
 379 to zero as t grows.

380 **EXAMPLE 2.** Consider the problem (1) on the spatial domain $\Omega = (-1, 1)$, with
 381 $u_0(x) = (1 - x^2)^{1+\alpha/2}$ and the source term fixed such that the exact solution is given
 382 by

$$383 \quad u(x, t) = (t + 1)^\gamma (1 - x^2)^{1+\alpha/2}, \quad \gamma < 0.$$

384 In our experiments, we select the model parameters $\kappa_\alpha = 0.1, \alpha = 1.2, \gamma = -1$ and we

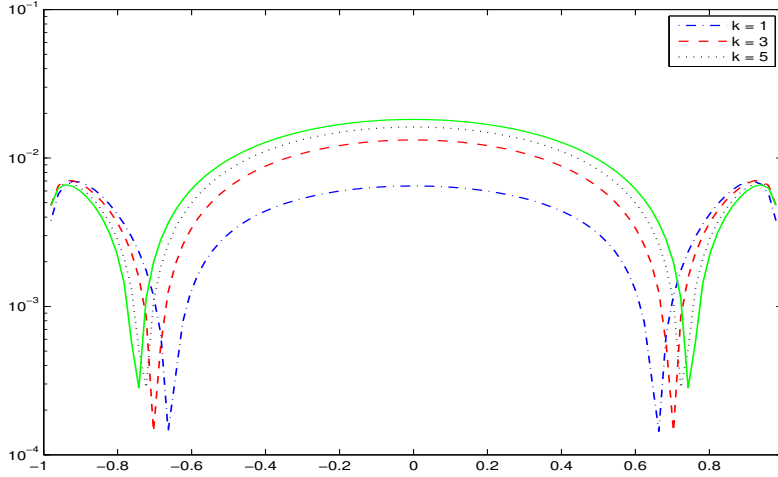


FIG. 4. Comparison of the errors provided by solving the problem of the Example 2 using both rational with $k = 1$ (blue dashed-dot-line), $k = 3$ (red dashed-line) and $k = 5$ (black dot-line) and MT (green solid-line).

385 *discretize the spacial domain using $N = 100$. In Figure 4 we report the step-by-step*
 386 *error provided by the numerical solutions obtained by applying the Gauss-Jacobi rule*
 387 *with $k = 1, 3, 5$ at $t = 0.1$ compared with the one obtained by solving directly (18).*
 388 *As expected, the solution provided by the rational approach is able to mimic, as k*
 389 *increases, the one derived by applying the matrix transfer technique. In particular,*
 390 *when $k = 5$ the corresponding error $\|\mathbf{u} - \mathbf{v}\|_2 = O(10^{-2})$ while, by the theoretical*
 391 *estimate (23), we would obtain this bound for $k = 7$.*

392 **EXAMPLE 3.** Consider now equation (1) in the one-dimensional case with non-
 393 linear source term

$$f(x, t, u) = u(1 - u),$$

395 known in the literature as the fractional Fisher-KPP reaction-diffusion equation. In
 396 this example, we set $\kappa_\alpha = 1$, $\alpha = 1.5$ and $\Omega = (-100, 100)$. We use $N = 500$ and
 397 we compute the numerical solutions provided by the matrix transfer technique and the
 398 rational approximation with $k = 3$. In particular, in the two pictures on the top of
 399 Figure 5 we have drawn the solutions profiles corresponding to the initial condition
 400 $u_0(x) = 0.5$ while, in the ones on the bottom, the solutions profiles corresponding to
 401 $u_0(x) = (100 + x)^2(100 - x)^2$ at time $t = 1$ (left) and $t = 10$ (right), respectively. As
 402 one can see, the two numerical solutions match very well and behave similarly with
 403 the theoretical one as t grows, cf. [13].

404 **EXAMPLE 4.** We solve the fractional reaction-diffusion equation in two space di-
 405 mensions

$$406 \quad \frac{\partial u(x, y, t)}{\partial t} = -\kappa_\alpha (-\Delta)^{\alpha/2} u(x, y, t) + f(x, y, t, u), \quad (x, y) \in (0, 1) \times (0, 1),$$

407 with

$$408 \quad f(x, y, t, u) = t^\alpha \frac{\kappa_\alpha}{16} \sum_{j=1}^4 (1 + \mu_j^{\alpha/2}) v_j + \alpha t^{\alpha-1} \sin^3(\pi x) \sin^3(\pi y) - \kappa_\alpha u,$$

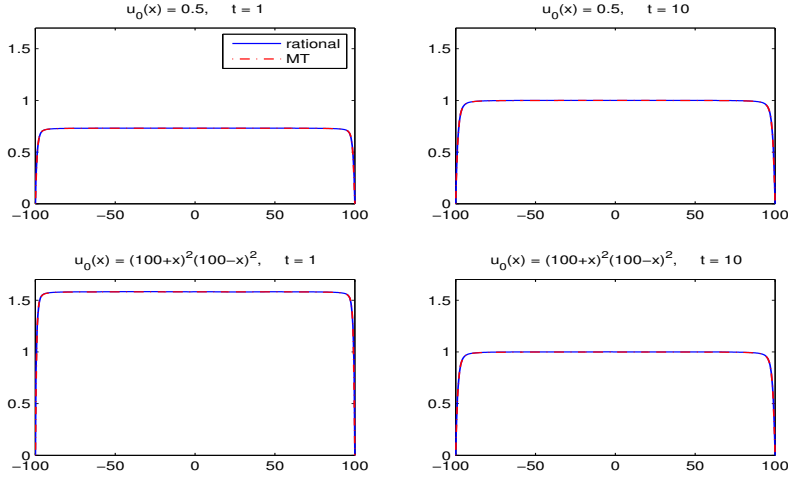


FIG. 5. Comparison of the numerical solutions of the problem in [Example 3](#) provided by MT and rational at $t = 1$ (left) and $t = 10$ (right) for two different choice of u_0 .

409 where

$$\begin{aligned}
 410 \quad v_1 &= 9 \sin(\pi x) \sin(\pi y), & \mu_1 &= 2\pi^2, \\
 411 \quad v_2 &= -3 \sin(\pi x) \sin(3\pi y), & \mu_2 &= 10\pi^2, \\
 412 \quad v_3 &= -3 \sin(3\pi x) \sin(\pi y), & \mu_3 &= 10\pi^2, \\
 413 \quad v_4 &= \sin(3\pi x) \sin(3\pi y), & \mu_4 &= 18\pi^2,
 \end{aligned}$$

415 subject to $u(x, y, 0) = 0$ and homogeneous Dirichlet boundary conditions [3].

416 The exact solution to this problem is

$$417 \quad u(x, y, t) = t^\alpha \sin^3(\pi x) \sin^3(\pi y).$$

418 The numerical solution provided by the rational approach based on the Gauss-
 419 Jacobi rule with $k = 7$ and the matrix transfer technique are drawn at $t = 1$ in
 420 [Figure 6](#) using $\alpha = 1.5$, $\kappa_\alpha = 10$ and $N = 40$ points in each domain direction. It
 421 worth noting that in order to obtain the same accuracy, the matrix transfer technique
 422 costs three times the rational approach.

423 **6. Conclusions.** In this paper we have proposed a rational approximation to
 424 the discrete fractional Laplacian. When applied for solving the reaction-diffusion
 425 equations this leads to a semi-discrete problem which can be solved in an efficient
 426 way due to the band structure of the matrices occurring in the definition of the
 427 approximation. With respect to the existing approaches based on the discretization of
 428 the Riesz derivative the main advantages are the ones of the matrix transfer technique
 429 itself, that is: the approach can be generalized to work in more than one dimension
 430 without modifying the overall solution methodology, it does not require to work with
 431 a uniform grid in space, all linear algebra tasks are with sparse matrices.

432

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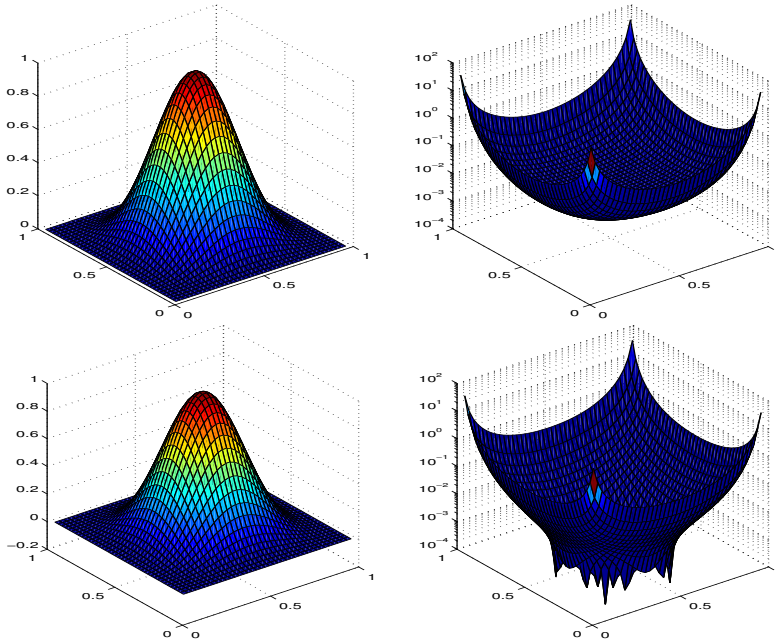


FIG. 6. Comparison of the analytical solution of the problem in *Example 4* with the numerical solution provided at $t = 1$ by rational (top) and MT (bottom) and corresponding relative errors (right) for $\alpha = 1.5$ and $\kappa_\alpha = 10$.

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