

Intermolecular Interactions in Eumelanins: a Computational Bottom-Up Approach. I. Small Building Blocks.

Giacomo Prampolini^{a,*}, Ivo Cacelli^{a,b}, and Alessandro Ferretti^a

^a*Istituto di Chimica dei Composti OrganoMetallici (ICCOM-CNR),
Area della Ricerca, via G. Moruzzi 1, I-56124 Pisa, Italy*

^b*Dipartimento di Chimica e Chimica Industriale, Università di Pisa,
Via G. Moruzzi 3, I-56124 Pisa, Italy*

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*Corresponding author

Abstract

The non-covalent interactions between pairs of the smallest eumelanins building blocks, as 5,6-dihydroxy-indole (DHI) and its redox derivatives, are subjected to a systematic theoretical investigation, elucidating their nature and commenting some of their possible effects on eumelanin layered structure. An accurate yet feasible protocol, based on second order perturbation theory, was set up and validated herein, and thereafter used to sample the intermolecular potential energy surfaces of several DHI related dimers. From the analysis of the resulting local minima, the crucial role of stacking interactions is assessed, evidencing strong effects on the geometrical arrangement of the dimer. Furthermore, the absorption spectra of the considered dimers in their most stable arrangements are computed, and discussed in relation to the well known eumelanin broadband features. The present findings may help in elucidating several eumelanin features, supporting the recently proposed geometrical order/disorder model (Chen *et al.*, *Nat. Commun.* 2014, **5**, 3859).

1 Introduction

As a consequence of their unique features, eumelanins have attracted an increasing attention, which has stimulated a continuous research in many fields, as testified in several recent reviews dedicated to these materials.¹⁻⁷ Besides their peculiar physical-chemical polymer-like properties,² the two more and well known characteristics of both natural and synthetic eumelanins are the intrinsic free radical character and the broadband UV/vis absorption, fundamental in its photo-protective function.⁸ Thanks to this fascinating ensemble of different properties, eumelanins have found an astonishing wealth of applications as bio-compatible multifunctional platforms in several fields as bio-interfaces,⁹ polymer stabilization,¹⁰ organic electronics^{4,11} or, more recently, hybrid materials.¹² A further boost has come from the multiple applications of a eumelanin like material, namely polydopamine, whose features have been recently reviewed as strongly related to those of eumelanins.⁷ In fact, since its first application in 2007 as an efficient coating material,¹³ polydopamine properties have been exploited in many technological fields, ranging from nanomedicine, to bio-sensors and energy production.¹⁴⁻¹⁶

Notwithstanding the many theoretical¹⁷⁻³⁸ and experimental^{2,5,13,33,39-50} studies, the structure-properties relationships subsisting the above mentioned multipurpose applications are still far to be clearly understood, and the patterns of supramolecular assembly or even the chemical nature of eumelanin constituting molecules are still argument of debate.^{7,35,36} The reason for this has been clearly pointed out in a very recent review by D'Ischia *et al.*,⁷ who traced back the difficulties on establishing a detailed picture of eumelanin structure to the fact that "*unlike the vast majority of natural pigments, eumelanins cannot be described in terms of a single well-defined structure, and it is not possible to provide an accurate picture beyond a statistical description of main units and functional groups*". More specifically, since the original amorphous semiconductor model³⁹ was definitively abandoned in the light of experimental evidence of a certain degree of supramolecular structure,^{40,41,43,49,51,52} many of eumelanin's peculiar characteristics (and in particular the broadband absorption spectra) were explained through the so-called chemical disorder model,^{3,4,27} where eumelanins are constituted by a *plethora* of distinct species, which may differ in structure, redox state or geometrical arrangement. In this

scenario, most of the resulting eumelanin properties can be rationalized as a consequence of such disorder: the broadband absorption, for instance, might be interpreted as a superposition of the spectra of eumelanins multiple constituents. The main lack of the chemical disorder model is undoubtedly neglecting the non covalent interactions occurring among the various eumelanin constituents nor the effect that the supramolecular motifs that originate from these interactions might have on eumelanin optical and mechanical properties, as reported by many investigations.^{18,23,36-38,48} Very recently Chen *et al.*^{36,38} have complemented the chemical disorder model with the geometrical order/disorder picture, where the existence of stacked layers containing a number of eumelanin forming units (protomolecules) and the conformational and configurational disorder within and among them is also taken into account (see also Scheme 3 in Ref. [7]). Despite the elusive nature of this description, several points seem now to be shared by the scientific community:¹⁻⁷ i) the basic² unit components (see also Figure 1) of natural and synthetic eumelanins are 5,6-dihydroxy-indole 2-carboxylic acid (DHICA), 5,6-dihydroxy-indole (DHI), its redox form indole-quinone (IQ) and its tautomer quinone-imine (MQ); ii) the above units are linked through covalent C-C bonds into small oligomers (protomolecules), ranging from tetramers to octamers, although a monomeric based picture has been also recently proposed;³⁴ iii) to date, the most probable and investigated oligomeric structures are those of a cyclic, porphyrin-like planar tetramer and linear or cross linked oligomers. It might be interesting to note that the former structure were mainly deduced from computational results,²⁶ whereas the latter were hypothesized, based on experimental measures;^{46,47} iv) whatever their structure, the protomolecules are arranged into stacked layers, with a typical separation of 3-4 Å; v) the dimensions of these stacked aggregates are nanometric and no long range orientational or positional order was found among the layers. Additionally, it is most probable that, within each layer, protomolecules can adopt a variety of geometrical conformations and/or redox state.

Based on these few points, it appears as it is now of fundamental importance to deeply understand and characterize the non covalent interactions among protomolecules, responsible for both the formation of the stacked layers and the tertiary structure arising from their spatial disposition. In this context, the potentiality of a computational approach can

be exploited to unravel the inherent complexity of eumelanin supramolecular structure. Surprisingly, up to our knowledge, only few attempts to thoroughly investigate non covalent interactions in eumelanins have been to date reported in literature. Indeed, apart from two recent works, where quantum mechanical (QM) methods as density functional theory (DFT)³¹ and its tight binding extension combined with self-consistent charge technique (SCC-DFTB)³⁵ were employed, most of the computational studies dealing with interaction between protomolecules were based on empirical, transferable force-fields (FFs),^{23,36,38} whose accuracy in handling non covalent forces has been recently questioned.^{53–55} Moreover, all previous studies dealt only with few spatial arrangements of the protomolecules. On the contrary, considering the crucial role of the non covalent interactions which may occur between eumelanin building blocks, the dependence from their relative orientation as well as the possible balance between Van der Waals (stacked) and hydrogen bond (HB) forces should be also taken into account. Therefore a more systematic though accurate investigation appears to be necessary.

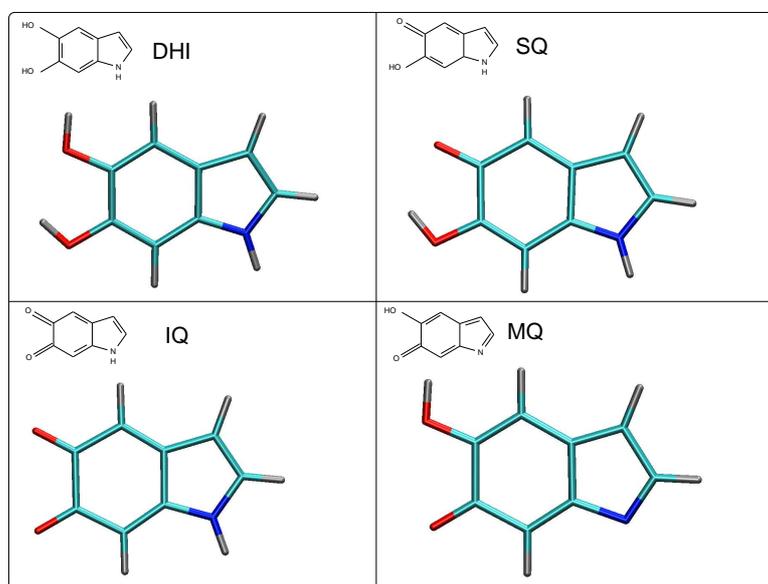


Figure 1: Monomeric species involved in eumelanins^{2,7} and considered in the present work: 5,6-dihydroxy-indole (DHI), its redox forms semi-quinone (SQ) and indole-quinone (IQ) and its tautomer quinone-imine (MQ).

The first step towards a deeper comprehension of protomolecules structural motifs

consists in devising a computational method able to accurately investigate the interactions between pairs of the smallest eumelanin building blocks, *i.e.* DHI and its redox (IQ) and tautomeric (MQ) forms. This is the main objective of the present work. Considering the aromatic character of all species considered, and the presence of one or more -OH functional groups, it is not unlikely that the most favorable arrangements of such pairs will be driven from a delicate balance between stacked (π - π) and HB interactions. The most accurate and reliable QM method to handle such kind of non covalent forces is probably the coupled cluster technique, evaluated at complete basis set (CCSD(T)@cbs), often referred as the gold standard of quantum chemistry.⁵⁶ Unfortunately, the prohibitive computational cost of this technique rules out its adoption in sampling a large number of non covalent dimer conformations. Computational convenience would suggest to turn to DFT dispersion corrected calculations but the accuracy of the results is known to depend on the choice of the functional and its empirical correction.^{57,58} In a rather time-consuming procedure, some purposely computed CCSD(T)@cbs data could be used as reference to select the most accurate empirical dispersion correction,^{57,59-61} although it has been recently reported that the performances of a given functional may vary dramatically with the target dimer and even with its spatial arrangement.^{62,63}

In this scenario, an alternative approach was adopted in our group, exploiting the combination of the second order Möller-Plesset perturbation theory (MP2) with purposely modified basis sets (therefore named⁶⁴ MP2mod). Very recently, the MP2mod method was successfully applied by us^{64,65} on a non covalent dimer, similar to those encountered in eumelanins, namely quinhydrone. As a matter of fact, MP2mod results obtained for quinhydrone well agreed with reference high-level CCSD(T)@cbs data, whereas, thanks to the reduced basis set dimensions, the observed computational times were remarkably reduced with respect to the other investigated methods.⁶⁴ Given the similarities (also recently pointed out in Ref. [34]) between quinhydrone's forming monomers and the eumelanin building blocks here considered, the MP2mod is expected to give good performances in this case also. Nonetheless, for a more robust assessment of the method accuracy, it is here again validated against CCSD(T)@cbs data, purposely computed on few selected arrangements. Subsequently, the MP2mod is employed to the calculation of the interaction

energy of different non covalent pairs, arising from the combination of the considered eumelanin minimal building blocks. Finally, the effect of the resulting dimer conformations on the shape of the absorption spectrum is also considered and discussed in some detail.

2 Computational Details

The initial geometry of each monomer was obtained by optimizing its energy at DFT level, using the B3LYP functional with the Dunning’s correlation consistent cc-pVTZ basis set.

The interaction potential energy surface (IPES) of each investigated pair was first sampled by computing the interaction energy in several dimer arrangements, constructed without altering the internal (B3LYP/cc-pVTz) geometry of each monomer. In this case the interaction (ΔE^{inter}) and the binding (ΔE^{bind}) energies between monomers A and B have the same expression:

$$\Delta E^{inter} = E_{AB} - (E_A^0 + E_B^0) = \Delta E^{bind} \quad (1)$$

where E_A^0 , E_B^0 are the energies of the two isolated monomers in their equilibrium geometry, whereas E_{AB} is the energy of the whole dimer AB in the chosen conformation.

Next, several local minima of each considered IPES were further explored by optimizing the dimer geometries corresponding to the most representative arrangements found in the previous sampling. In this case, since the monomer geometries may change during the full optimization, ΔE^{inter} and ΔE^{bind} may differ, as

$$\Delta E^{inter} = E_{AB} - (E_A + E_B) \quad (2)$$

where E_A and E_B are the energies of the two monomers, computed for each monomer in the geometries assumed within the AB dimer. In this case, the binding energy is

$$\Delta E^{bind} = E_{AB} - (E_A^0 + E_B^0) = \Delta E^{inter} + (\Delta E_A + \Delta E_B) \quad (3)$$

where

$$\Delta E_M = E_M - E_M^0 \quad ; \quad M = A, B \quad (4)$$

is the relaxation/distortion energy of monomer A (or B), due to the AB complex formation. To evaluate this term, for each investigated complex E_M^0 was obtained after a

geometry optimization performed on the isolated monomer M at the same level of theory (MP2mod, *vide infra*) employed in the calculation of ΔE^{inter} for dimer AB .

On each investigated AB dimer geometry, the computation of the interaction energies was performed at two different levels of theory.

- i) For selected arrangements (either obtained by assembling the two rigid monomers or through dimer optimization) ΔE^{inter} was computed according to eq. (2), at the CCSD(T)@cbs⁵⁶ level, following the procedure employed by us for quinhydrone,⁶⁴ and reported in detail in the Supporting Information.
- ii) All considered dimer pairs were investigated through the MP2mod method,⁶⁴ computing both ΔE^{inter} and ΔE^{bind} , according to equations (2) and (3), respectively. All calculations were performed with the 6-31G*($\alpha_C, \alpha_O, \alpha_N$), where α_C , α_O and α_N are the modified polarization exponents for Carbon, Oxygen, and Nitrogen atoms, respectively. The first two exponents, namely 0.25 and 0.44, were transferred from the ones optimized for quinhydrone,⁶⁴ whereas α_N was optimized in this work through the EXOPT^{64,66,67} procedure against the reference CCSD(T)@cbs data obtained for pyrrole, as detailed in the Supporting Information and in the next section.

Finally, the absorption vertical transition energies for the first 25 excited states were obtained for eumelanin related monomers and dimers at the TD-DFT level, using the B3LYP functional with the aug-cc-pvDZ basis set. Absorption spectra were subsequently obtained by broadening each transition with a Gaussian function with an half-width at half maximum of 0.33 eV. All calculations were performed with the GAUSSIAN09 package.⁶⁸

3 Results and Discussion

3.1 MP2mod set up

The eumelanin related monomers considered in this work and displayed in Figure 1 all consist in two condensed aromatic moieties, namely an *o*-hydroquinone derivative and one pyrrole ring. Therefore, a first validation of the MP2mod reliability can be obtained

by testing the method performances on these two smaller molecules. To obtain the α polarization exponents for the modified 6-31G* basis set required by the MP2mod method, two different strategies were followed.

As far as *o*-hydroquinone derivatives are concerned, the polarization exponents for Carbon (0.25) and Oxygen (0.44) atoms were transferred from previous work performed on a very similar dimer (quinhydrone).⁶⁴ To test this assumption, *o*-benzoquinone/*o*-hydroquinone pairs were prepared in three different geometries (see Figure A of the Supporting Information): face-to-face (FF), antiparallel face-to-face (AFF) and hydrogen bonded (HB). The intermolecular distance between the monomers centers of mass was varied stepwise in the 2.8 Å – 7.0 Å range, and the interaction energy was computed for each of the resulting dimer arrangements at MP2mod/6-31G*(0.25, 0.44) level. From the computed curves, reported in detail in Figure B of the Supporting Information, two main observations may be drawn. First, as could be expected, it is clear from the difference between the MP2mod energies and their HF contributions that the dispersion forces routed in the dynamic electron correlation dominate the stacked wells (FF and AFF), and play only a minor (but no negligible) role on H-bonded (HB) dimers. As a consequence, the FF and AFF arrangements are much more favorable than the HB one.

Starting from the most stable AFF geometry, a complete energy optimization was carried out at the MP2mod level. The final arrangement (dimer I) is shown in the left

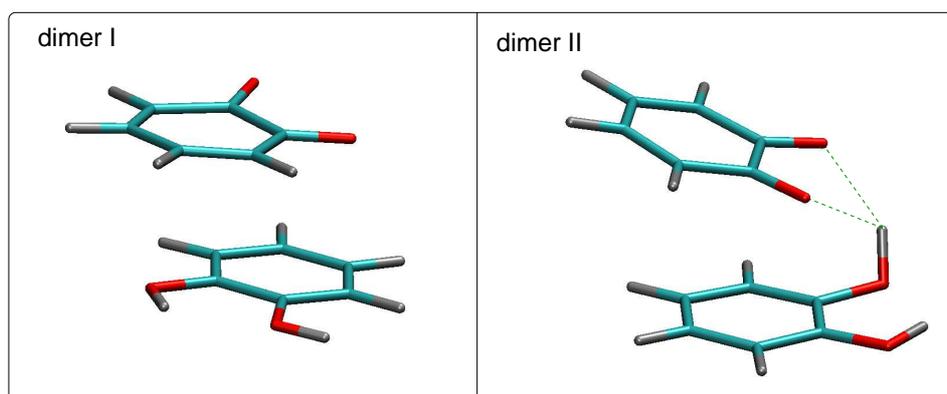


Figure 2: Optimized geometries of the *o*-benzoquinone/*o*-hydroquinone pair computed at the MP2mod level, starting from the AFF (dimer I, left) and HB (dimer II, right) arrangements, displayed in Figure A of the Supporting Information. Hydrogen bonds in dimer II are highlighted with green dashed lines.

panel of Figure 2. The interaction energy ΔE^{inter} of this geometry has been computed according to equation (2) at both MP2mod and CCSD(T) levels, as reported in Table 1. For a more extended validation, a second optimization was also carried out at MP2mod level, starting with the HB minimum conformation. As could be expected from the comparison of the AFF and HB curves (see Supporting information, Figure B), the HB dimer optimization ends up into a stacked conformation (dimer II), shown in the right panel of Figure 2. Differently from dimer I, in dimer II one hydroquinone’s Hydrogen is found out of the aromatic plane, allowing for the formation of two weak hydrogen bonds, which concur to further stabilize the stacked geometry. MP2mod and CCSD(T) interaction energies for dimer II were then computed, and are compared to each other in Table 1. For both dimers, the agreement between MP2mod and CCSD(T) reference values

Dimer	Method	Basis set	ΔE^{inter} (kcal/mol)	ΔE^{bind} (kcal/mol)
dimer I	CCSD(T)	cbs	-5.34	-
dimer I	MP2mod	6-31G*(0.25,0.44)	-5.80	-5.35
dimer II	CCSD(T)	cbs	-9.54	-
dimer II	MP2mod	6-31G*(0.25,0.44)	-9.91	-7.04

Table 1: Interaction (ΔE^{inter}) and binding energies (ΔE^{bind}) computed with different methods on quinone dimer I and dimer II. The ΔE^{bind} value at CCSD(T)/c.b.s. level is not reported as the calculation of the relaxation/distortion energy at cbs is not strictly defined.

is rather good, being the ΔE^{inter} difference less than 0.5 kcal/mol, and the adoption of the α_C and α_O exponents to handle *o*-hydroquinone derivatives appears reliable.

Since no α exponent was available for the Nitrogen atom in pyrrole, α_N was optimized following the procedure adopted for quinhydrone in Ref. [64], and discussed in some detail in the Supporting Information. The optimized polarization exponent for the Nitrogen atom resulting from the EXOPT minimization is 0.37, a value not far from that found for Oxygen atom in hydroquinone (0.44).⁶⁴ The pyrrole IPES was scanned according to selected cross sections, described in Figure C of the Supporting Information. By looking at Figure D in the Supporting Information, the MP2mod results well match the scenario that emerged from previous calculations performed on pyrrole dimer and reported in literature (see Ref.[69] and reference therein). Indeed, the stacked face-to-face (FF) arrangements are found not to be stable, whereas the antiparallel displaced

one shows a significant well (~ 3.5 kcal/mol) at rather short ring-ring distances (~ 3 Å). More important, the most favorable conformations (> 5 kcal/mol) among the four investigated arrangements are those where an hydrogen bond between the -NH group of one monomer and the π cloud of the other is established. To further check the reliability of the MP2mod protocol, two additional dimer optimizations were performed, starting from the two most favorable geometries (TS and APD, see Figure C and D in Supporting Information). The final structures, labeled dimer I and II, are reported in Figure 3. In the first case, the two monomers are found in a T-shaped like arrangement, where

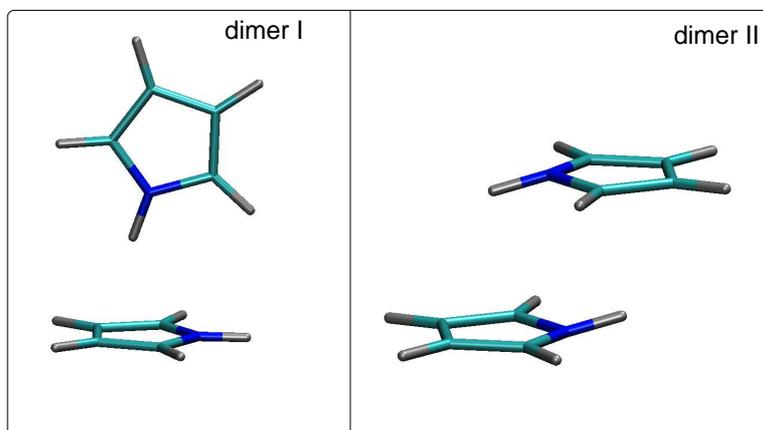


Figure 3: Optimized geometries of the pyrrole dimer computed at the MP2mod level, starting from the TS (dimer I, left) and APD (dimer II, right) arrangements, displayed in Figure C and D of the Supporting Information.

the Hydrogen of on -NH group strongly interacts with the π cloud of the neighboring monomer, being the distance with the aromatic plane around 2 Å. However, because the HF contribution is only slightly attractive (-0.5 kcal/mol), dispersion forces still play the major role to dimer stabilization. As far as dimer II is concerned, the planar rings are found in a stacked conformation, though displaced by ~ 2 Å and with the NH bonds pointing in opposite directions. The interaction energies were computed for both dimers at MP2mod level with the optimized 6-31G*(0.25,0.37) basis set, finding $\Delta E^{inter} = -5.41$ and -5.97 for dimer I and II, respectively. Despite these values are well in the interaction energy range (-4.0 – 6.3 kcal/mol) found by previous high level calculations,⁶⁹ the many local minima expected for the pyrrole IPES do not allow for a direct comparison. To

better assess MP2mod reliability, the ΔE^{inter} was computed for dimer I and II also at CCSD(T)@cbs level, and the results are compared with the MP2mod counterparts in Table 2. It appears that the agreement with the reference values is remarkable, being 0.1 kcal/mol and 0.5 kcal/mol the differences between reference and MP2mod values for dimer I and II geometries, respectively.

Dimer	Method	Basis set	ΔE^{inter} (kcal/mol)	ΔE^{bind} (kcal/mol)
dimer I	CCSD(T)	cbs	-5.51	-
dimer I	MP2mod	6-31G*(0.25,0.37)	-5.41	-5.35
dimer II	CCSD(T)	cbs	-6.52	-
dimer II	MP2mod	6-31G*(0.25,0.37)	-5.97	-4.77

Table 2: Interaction (ΔE^{inter}) and binding energies (ΔE^{bind}) computed with different methods on pyrrole dimer I and dimer II. The ΔE^{bind} value at CCSD(T)/c.b.s. level is not reported as the calculation of the relaxation/distortion energy at cbs is not strictly defined.

Exploiting the exponent reliability found for both o-quinones and pyrrole dimer cases, the 6-31G*(0.25,0.44,0.37) basis set was implemented, and the MP2mod procedure was applied in the calculations of the interaction energy between pairs of the eumelanin related monomers displayed in Figure 1.

3.2 DHI/IQ IPES

The most probable protomolecules that may concur to form eumelanin structure have been recently reviewed in references [36], [38] and [7]. Despite several models were proposed (ranging from monomers to octamers), all candidates share the same kind of building blocks,² *i.e.* the DHI derivatives shown in Figure 1. Among all possible combinations, we started the IPES investigation with the DHI/IQ pair. Indeed, beside being one of the most probable combinations that are generated when two of the aforementioned protomolecule’s models are considered in a stacked arrangement, it is expected that the most stable DHI/IQ non-covalent dimers will be originated by dispersion (stacking) interactions and hydrogen bonds, whose delicate interplay is a challenging benchmark for the MP2mod protocol. Finally, many features of the DHI/IQ dimer closely resemble those of quinhydrone, recently⁶⁴ studied by us with the same computational technique.

DHI/IQ pairs were initially placed in eight different arrangements, devised to consider all types of intermolecular interactions (stacking, H- π , hydrogen bond, *etc.*) that could arise between such non-covalent dimers, *i.e.* stacked face-to-face (FF), anti-parallel face-to-face (AFF), side-to-face (SF), hydrogen bonded (HB), parallel displaced (PD), anti-parallel displaced (APD), stacked rotated (ST) and not stacked rotated (NST). For each arrangement, a number of dimer geometries was built by moving one monomer with respect to another along one IPES coordinate (either a distance or an angle), without changing the monomer internal geometries. All considered configurations and the chosen displacement coordinates are displayed in the insets of Figure 4. Finally, for each arrangement, the IPES cross sections, also displayed in Figure 4, were obtained at the MP2mod/6-31G*(0.25,0.44,0.37)a level by computing the interaction energy of the proper set of configurations.

The first general observation that arises by looking at the resulting curves, is that the DHI and IQ monomers show rather strong non-covalent interactions, which remarkably depend on their relative orientation, and span ΔE^{inter} range between -10 and -3 kcal/mol. Notwithstanding the stacking interactions are dominant, the most stable conformers are found when the molecular planes are twisted (ST, -9.51 kcal/mol) or displaced along the plane direction (APD, -9.28 kcal/mol). In all cases, an antiparallel arrangement of the CO dipoles is preferred to a parallel one. It is also worth noticing that in all the stacked configurations considered, the inter-ring distance corresponding to the curve minimum is found between 3.4 to 3.6 Å, a value in good agreement with range (3-4 Å) proposed^{7,36,40,41} for eumelanin inter-layer distance. Interestingly, while the lack of stacking interactions causes the HB configurations to be the least stable, the side-to-face like pairs (SF, -7.95 kcal/mol and NST, -8.10 kcal/mol) also show a non-negligible interaction, although the distance between the monomer centers of mass is around 5 Å. As a consequence, even in the (most likely) hypothesis that eumelanin protomolecules consist on DHI related oligomers, the latter side-to-face arrangements between the building blocks of two neighboring oligomers cannot be ruled out, as they could be realized by a rotation around the C-C bond connecting two protomolecules moieties, if the gain in non-covalent energy is more than the one needed for the dihedral torsion. However, a full investigation

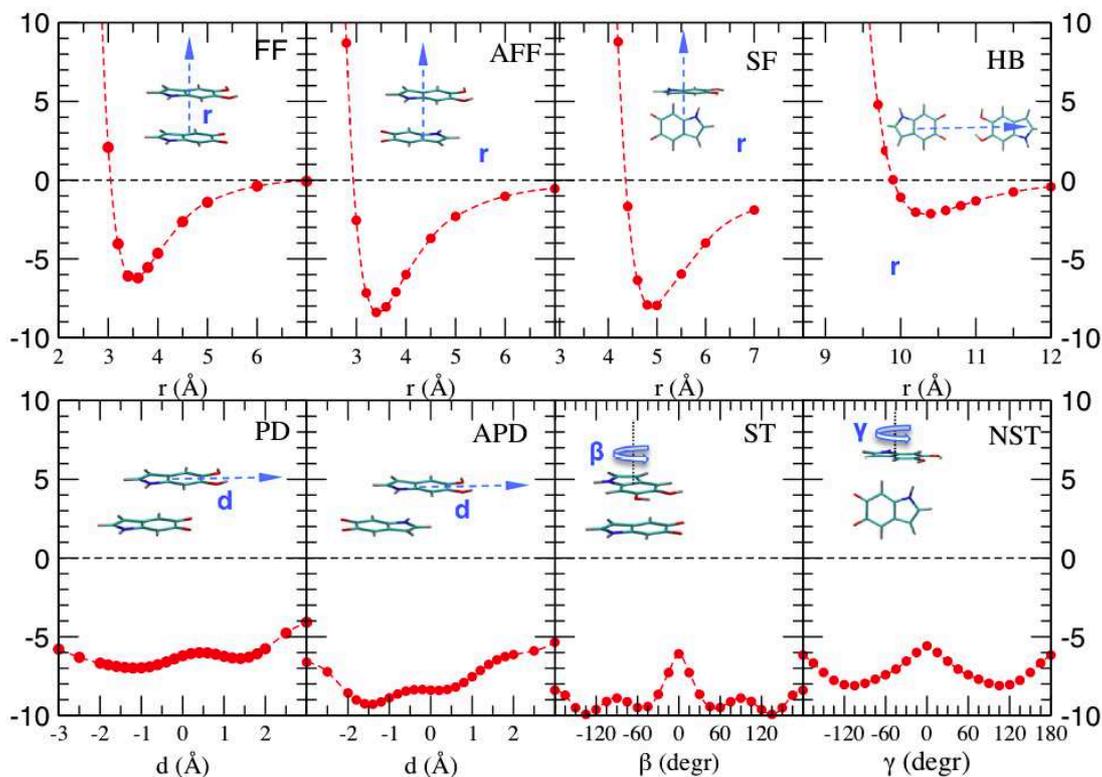


Figure 4: MP2mod interaction energy curves, computed for the DHI/IQ pair in different arrangements. From left to right: face-to-face (FF, top) and parallel displaced (PD, bottom), antiparallel face-to-face (AFF, top) and antiparallel displaced (APD, bottom), side-to-face (SF, top) and stacked (ST, bottom), hydrogen bonded (HB, top) and not stacked (NST). The scanned IPES coordinates in abscissa is highlighted in blue in the inset of each panel.

of the subtle balance of the interaction energy between oligomers building blocks and the torsional barriers driving the protomolecules internal structure is beyond the aims of the present work, but will be the main topic of a future investigation currently in progress in our group.

3.3 MP2mod validation on DHI/IQ dimers

Before attempting to study other DHI related pairs possibly involved in eumelanin, as for instance those involving the MQ tautomer, a final test was performed on the optimized basis set, to validate the assumption of polarization exponents transferability. To this purpose, four dimer optimizations were performed at MP2mod level starting from selected

DHI/IQ arrangements, namely ST, APD, NST and HB. The final geometries, labeled A-D, are shown in Figure 5 and their interaction energies, computed both at MP2mod and CCSD(T)@cbs level, are reported in Table 3.

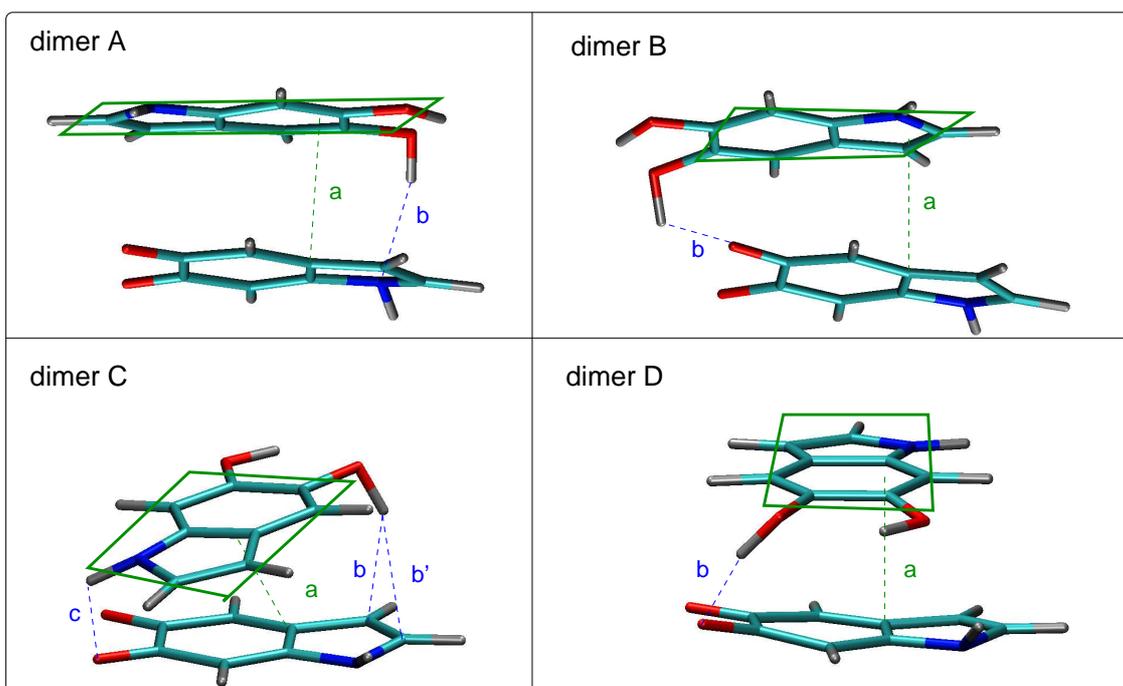


Figure 5: Optimized geometries of the DHI/IQ pair, starting from selected minima of the APD (dimer A), HB (dimer B), NST (dimer C) and ST (dimer D) curves (see Figure 4). The geometrical parameters reported in Table 3 are highlighted with blue and green dashed lines.

The comparison between the MP2mod values and those obtained by the CCSD(T)@cbs method is rather good, with a standard deviation of ~ 1 kcal/mol, *i.e.* less than 10% of the total ΔE^{inter} value. It is worth noticing that the accuracy obtained with the MP2mod protocol is even more appreciable if one considers the computational times, reported in the last row of Table 3, that underline the feasibility of MP2mod calculations, even on a large number of dimer arrangements.

Once the reliability of the MP2mod has been finally assessed, a closer look to the four minimum energy conformations found by MP2mod optimizations can help in understanding the type and strength of the non-covalent interactions established between DHI and IQ units. First, a general feature shared by the four pairs is the stacked arrangement

Dimer	a (Å)	b (b') (Å)	$\Delta E_{CCSD(T)}^{inter}$ (kcal/mol)	$\Delta E_{MP2mod}^{inter}$ (kcal/mol)	ΔE_{MP2mod}^{bind} (kcal/mol)
A	2.82	2.11	-11.4	-12.4	-10.0
B	2.80	1.92	-14.4	-15.6	-10.7
C	2.95	2.25 (2.54)	-11.0	-12.4	-10.1
D	2.95	1.92	-14.3	-14.8	-12.1
CPU time (hours)			~ 15000	~ 1	

Table 3: Geometrical parameters, interaction (ΔE^{inter} , computed at MP2mod/6-31G*(0.25,0.37,0.44) and CCSD(T)@cbs level) and binding (ΔE^{bind} , last column, only at MP2mod level) energies computed on DHI/IQ dimers A-D displayed in Figure 5. a and b (or b') are highlighted in the same Figure, and correspond to the distance between the two aromatic planes and to the minimum distance between a DHI Hydrogen and one IQ heavy atom, respectively. In the last row CCSD(T) and MP2mod computational times on a Xeon E5-2670 2.60 GHz processor are also reported for comparison.

of the two monomers, which is a clear indication of the fundamental role of dispersion forces. Indeed, the two molecular planes are in all cases very close, being the inter-plane distance slightly less than 3 Å. Another specific characteristic of the DHI/IQ pair is the formation of a hydrogen bond between the two units, in addition to the intra-molecular one found in the DHI monomer for all A-D conformations. The inter-molecular H-bond is formed through an out-of-plane rotation of one Hydrogen of the -OH group around the C-OH bond. For dimer B and D, this rotation is very effective, as the intermolecular H-bond distance is found below 2 Å. It is worth noticing that the same kind of torsion takes place also in the quinhydrone dimer,⁶⁴ and is involved in both dimer stability and electron and proton transfers between the redox pair in solution.⁶⁵ Finally, by comparing the $\Delta E_{MP2mod}^{inter}$ and ΔE_{MP2mod}^{bind} values reported in Table 3 for all investigated dimers, it appears that the distortion energy [$\Delta E_{DHI} + \Delta E_{IQ}$] reported in equation (4) ranges from 2 to 5 kcal/mol, depending on dimer conformation. However, the former contribution (ΔE_{DHI}) is more than 75% of the total distortion energy, indicating that the larger distortions from the monomer equilibrium geometries are located in the DHI species, and essentially consist in the aforementioned dihedral rotation.

3.4 Additional pairs of eumelanin's building blocks

In the hypothesis that the structure of eumelanins can be described by groups of stacked layers, in turn composed by small size oligomers of covalently bound DHI related units,⁷

it can be expected that the intermolecular interactions between two neighboring building blocks belonging to different layers are very similar to those settled in the above discussed DHI/IQ arrangements. In this scenario, it is clear how the relatively high interaction energies found between DHI and IQ moieties can strongly influence both the relative orientation of two facing oligomers (through dispersion) and the distance between the stacked layers (through H-bonds) to which they belong. Nonetheless, these considerations rely on the results obtained for only one of the possible redox/tautomeric pairs that can be thought to take place in eumelanins; to draw more sound conclusions, other DHI related non-covalent dimers should be investigated. Among the many possible combinations between the monomers shown in Figure 1, to contain computational burden here will be only considered those pairs that can arise when two of the most reliable and investigated⁷ oligomer models proposed for eumelanin (*i.e.* the planar cyclic tetramer first proposed by Kaxiras *et al.*²⁶ from computational findings, and the cross-linked, non-cyclic tetramers hypothesized on experimental evidences by the D’Ischia group^{46,47}) are disposed in a stacked arrangements. As far as the not cyclic models are concerned, the most recurring monomeric building block is DHI, although IQ, MQ and SQ units can be also contained in the cross-linked tetramer.^{7,46,47} Hence, besides the DHI/IQ pair, considering the higher occurrence of DHI units within the oligomer, only the DHI/DHI, DHI/MQ and DHI/SQ combinations have been taken into account for additional MP2mod calculations. Moreover, since the cyclic tetrameric model is made up of IQ and MQ building blocks,^{26,36} the interactions between these two moieties (IQ/MQ) have been also investigated.

The starting conformations employed in the previous MP2mod optimizations, carried out on the DHI/IQ pair (that eventually led to A-D dimers displayed in Figure 5) were employed as starting points also for the other considered dimers, after proper addition/removal of Hydrogen atoms. The final dimer arrangements (labeled E to T) are displayed in Figure 6. By looking at these optimized conformers, and at the geometrical and energy data reported in Table 4 for all pairs, some general conclusions may be drawn.

The first observation is that, as already noticed for the DHI/IQ pair, most of the considered optimizations end up in a stacked arrangement, independently of the starting conformation. More precisely, after optimization, only P, R and T dimers are not found in

a stacked-like arrangement, but it should be mentioned that in all three cases the starting conformation was also not stacked. However, by looking at the ΔE^{bind} values reported in Table 4, it appears as they are among the less stable conformers. On the contrary, dimers B, F, J and N all start from a side-by-side, H-bonded conformer (see Figure 4, top left panel), but, during optimization, the molecular planes rotate around the intermolecular H-bond(s) and the two moieties are eventually found in a almost aligned, stacked geometry. By looking at the IQ/MQ pair, where only one or no H-bonds can be formed, it is clear that the stabilization of the stacked arrangements is ensured essentially by dispersion forces, which may contribute up to ~ 8 kcal/mol to the pair stability. The formation

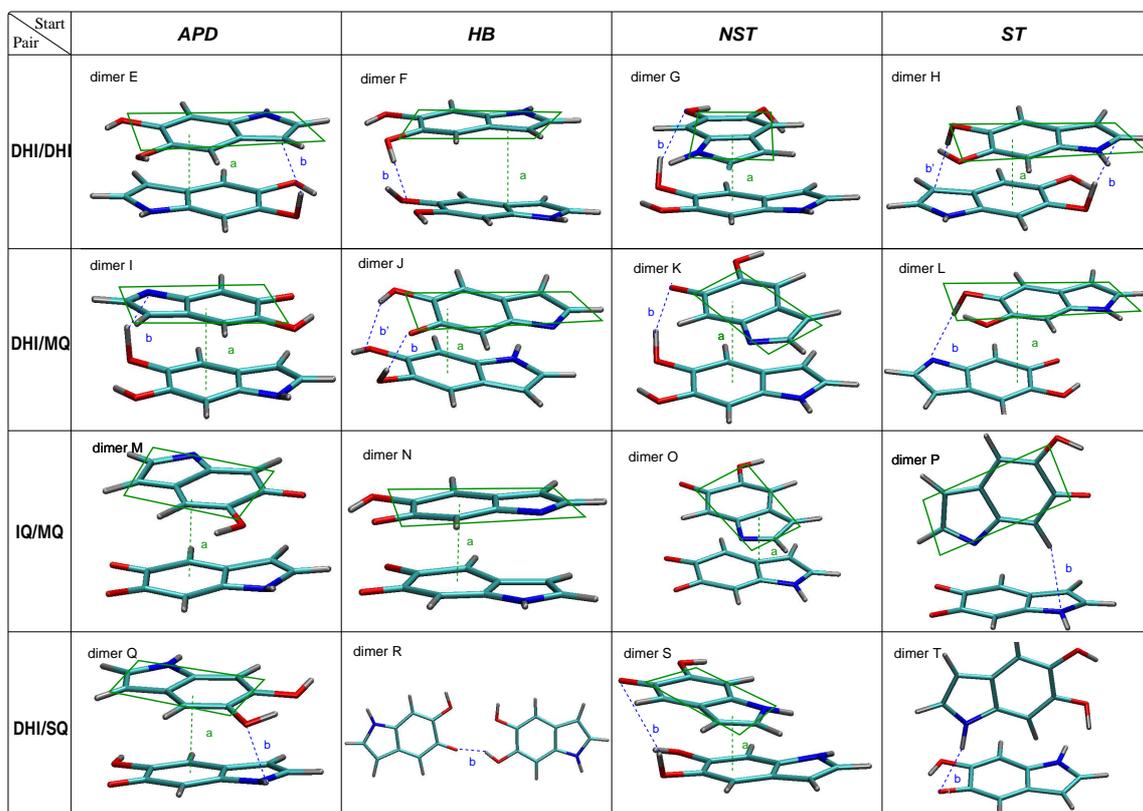


Figure 6: From top to bottom, Mp2mod optimized geometries of the DHI/DHI (E-H), DHI/MQ (I-L), IQ/MQ (M-P) and DHI/SQ (Q-T) pairs, starting from same initial conformations (*i.e.* antiparallel displaced (APD), hydrogen bonded (HB), not stacked (NST) and stacked (ST)) used in the optimization of the DHI/IQ shown in Figure 5 as indicated in the top panel. The geometrical parameters reported in Table 4 are highlighted with blue and green dashed lines.

of one or more H-bonds, of course, favors the stability of the dimer with an additional gain in binding energy that can be quantified, depending on the case considered, between 2 and 4 kcal/mol. Furthermore, when H-bonds are formed between the two moieties, the two monomers tend to orient their long axes in a almost parallel (or antiparallel) disposition, whereas, in their absence, the molecular planes are fund mostly twisted with respect to each other. However, by looking at the binding energies reported in Tables 3 and 4, it seems that there is no systematic trend in the strength of the non-covalent interaction settled by different types of monomers pairs, although DHI/IQ, DHI/DHI and DHI/MQ dimers are in average more stable than those involving IQ/MQ and, in particular, DHI/SQ.

Dimer	a (Å)	b (b') (Å)	ΔE^{inter} (kcal/mol)	ΔE^{bind} (kcal/mol)
DHI/DHI				
E	3.00	2.21	-11.4	-7.8
F	3.10	2.04	-6.9	-4.5
G	2.95	1.95	-15.7	-10.7
H	2.92	2.16 (2.17)	-15.8	-11.3
DHI/MQ				
I	2.72	2.00	-21.9	-14.3
J	2.90	1.85 (2.90)	-10.5	-8.0
K	2.70	1.66	-11.4	-6.4
L	2.70	2.02	-18.5	-13.8
IQ/MQ				
M	3.20	-	-14.0	-6.3
N	2.80	2.51	-11.8	-7.1
O	2.85	-	-6.7	-4.8
P	-	2.56	-5.9	-5.7
DHI/SQ				
Q	3.14	2.56	-6.8	-4.9
R	-	1.91	-3.3	-3.1
S	2.87	2.92	-11.4	-7.2
T	-	2.30	-3.9	-3.5

Table 4: Geometrical parameters, interaction (ΔE^{inter}) and binding (ΔE^{bind}) energies, computed for the DHI related pairs shown in Figure 6. a and b (or b') are highlighted in the same Figure, and respectively correspond to the distance between the two aromatic planes and to the minimum distance between a Hydrogen of one monomer and one heavy atom of the other.

3.5 The effect of stacking on absorption spectrum

Another feature shared by all investigated pairs concerns with the inter-monomer separation between the stacked planes, which falls, in almost all cases, between 2.7 and 3.1 Å. At such small separations, the π clouds of the two monomers do show a non-negligible overlap, hence affecting frontier orbitals in their shape and energies. For instance, in the

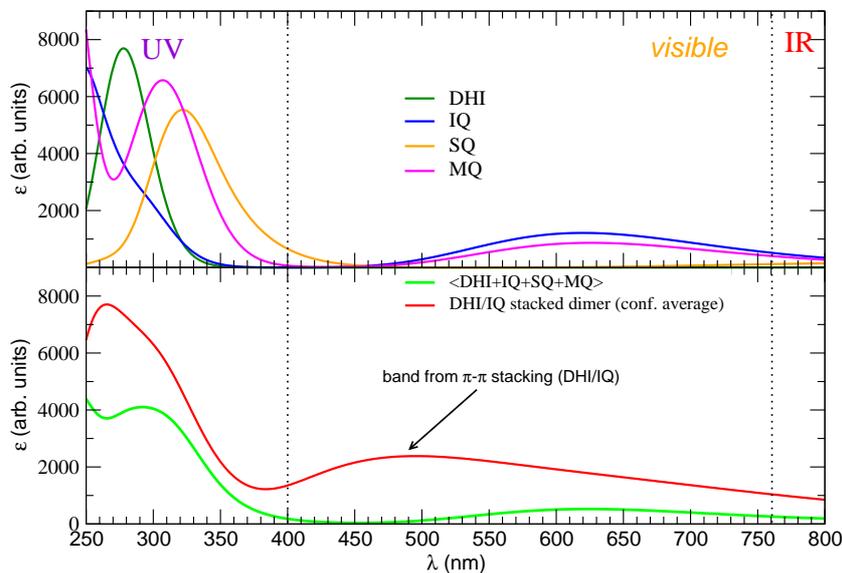


Figure 7: Top panel: TD-DFT absorption spectra of the monomer species displayed in Figure 1: DHI (blue line), IQ (magenta), MQ (cyan) and SQ (black). Bottom panel: Average of all monomer spectra (light green line) and averaged (red line) of the absorption spectra computed for DHI/IQ dimers A-D, displayed in Figure 5.

quinhydrone dimer, recent TD-DFT calculations⁶⁴ have shown that the stacking distance strongly influences the absorption spectrum, causing the rising of a broad band in the visible region as the separation between the two monomers forming quinhydrone is reduced. As a matter of fact, concerning eumelanin's broadband absorption, the important consequences of the excitonic coupling that takes place in stacked arrangements was first pointed out by Stark and coworkers,²³ and has been very recently invoked to complement the chemical disorder model with the latest geometrical order/disorder one.³⁸ However,

a possible limitation of these previous computational works, may stand in the method employed to originate the stacked conformers (empirical force-fields) as well as on the low level of theory employed for computing the spectrum (ZINDO or semi-empirical Hamiltonian). In this work, thanks to the reduced dimensions of the investigated targets, a more reliable approach was employed both to select the involved geometries (MP2mod) and to compute the optical properties (TD-DFT).

To investigate the effects of the afore-discussed dimers on the absorption properties of eumelanin, TD-DFT calculations were performed on both the isolated monomers and their optimized combinations. The accuracy of the chosen functional (B3LYP) was

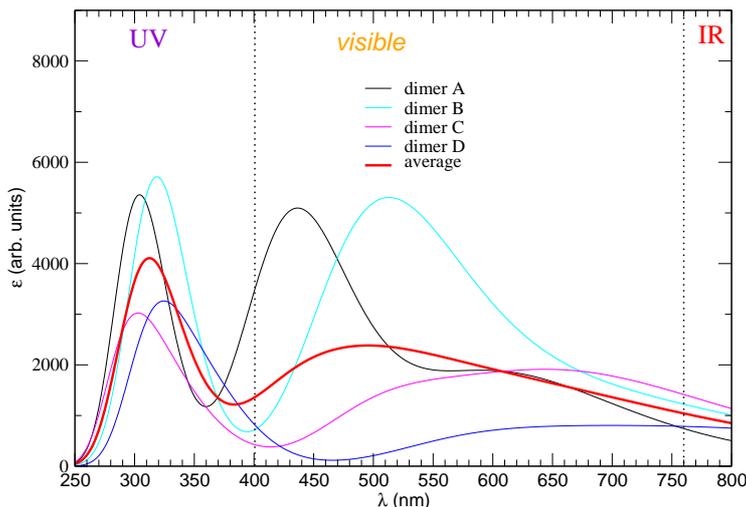


Figure 8: TD-DFT absorption spectra of the DHI/IQ dimers, in the optimized conformations displayed in Figure 5: A (black line), B (cyan), C (magenta) and D (blue). The spectra resulting from the superimposition of the former band shapes is reported with a solid red line.

validated (see Figure E in the Supporting Information) by comparing the computed absorption spectrum of the DHI monomer with the available experimental data.⁴² The absorption spectra of all monomers considered in this work have been computed and reported in the top panel of Figure 7. By looking at the resulting spectral lines, it appears as the differences in the redox form (DHI *vs* SQ, *vs* IQ) and in the geometrical

arrangement of their constituting atoms (MQ *vs* IQ) may cause significant shifts in the position of the strong absorption peak in the UV region and/or changes in intensity in the broad band in the visible range. In agreement with the chemical disorder model, when the separate monomer signals are averaged (see light green line in the bottom panel of Figure 7), a broad band is formed in the UV region, whose intensity is gradually increasing at decreasing wavelengths, as expected for eumelanin.^{2,3} On the contrary, not all the visible region is covered by the average absorption band, and small or null absorption intensity are found in the 350-550 nm range, confirming that the eumelanin broadband absorption, which covers monotonically (see top panel of Figure 10) the whole UV/visible spectrum, cannot be simply originated by the superposition of the absorption lines of the not-interacting monomers.

In Figures 8 and 9 the absorption spectra obtained for the considered non-covalent dimer arrangements A to T are displayed. By looking at Figure 8, where only the spectra

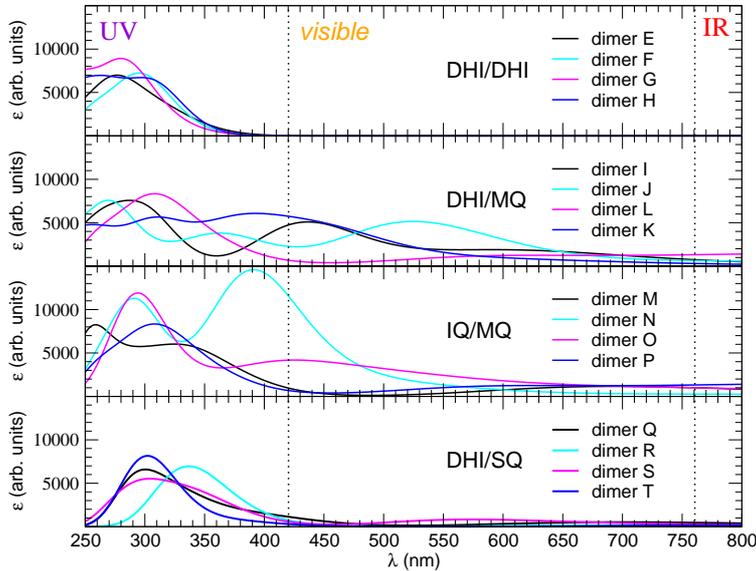


Figure 9: TD-DFT absorption spectra of the DHI related pairs in the optimized conformations displayed in Figure 6.

relative to the DHI/IQ pair are displayed, it is evident that the formation of the stacked

complex, and the consequent excitonic coupling, causes the appearance of a much more diffuse absorption band, which in some cases (dimer C) almost covers the whole visible region or presents maxima (dimer A) at wavelengths around 400-450 nm, where the separate monomers do not absorb at all. This last observation is also evidenced in Figure 7, where the absorption band of the DHI/IQ dimer, averaged over the four considered minima (A-D), is directly compared to the one arising from the monomeric species.

A more complete picture can be achieved by also taking into account the absorption due to the other dimer configurations E-T (see Figure 9), built by assembling different monomer pairs. Similarly to what previously found for the DHI/IQ dimer, all the DHI/MQ and IQ/MQ complexes (central panels of Figure 9) show rather diffuse bands, whose position and intensity strongly depend upon the dimer arrangement. In particular dimer M, formed by the IQ and MQ monomers, shows an intense band centered around 400 nm, a region not covered when considering all the isolated monomers. Conversely, dimers E - H and Q - T are characterized by a much reduced dependency upon the relative orientation of their constituting monomers, being all bands centered in the high energy region (~ 100 nm).

Finally, it is interesting to superimpose all the computed signals, either originated by monomers or non-covalent dimers, and compare the resulting curve with the experimental³ absorption band shape of real eumelanin. The outcome of both these operations are displayed in Figure 10. As shown in the upper panel, the computed curve well matches most of the characteristics of the experimental signal, as the monotonic increase towards lower wavelengths and the broad band covering great part of the investigated spectrum. Nonetheless, it is important to stress that the computed average displayed in Figure 10 neglects the contributions to the spectral signal due to chemical heterogeneity, which comes into play when covalent combinations of different building blocks (*e.g.* tetramers) and their relevant redox states are also taken into account. This however goes beyond the aims of the present study, and is the object of investigations currently in progress in our group.

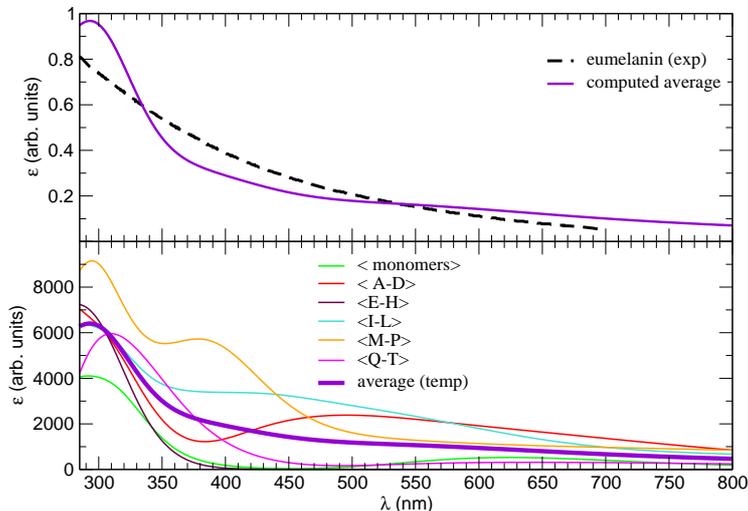


Figure 10: Top: eumelanin experimental³ absorption spectrum. Bottom: absorption spectrum computed in this work by averaging the TD spectra obtained for monomers and their combinations in the discussed optimized conformations.

4 Conclusions

A key point in the investigation of eumelanin properties certainly is a thorough comprehension of the intermolecular interactions that determine eumelanin’s peculiar structural properties. Because of the large dimensions of the protomolecules that are thought to compose eumelanin layers, a mandatory step in investigating their interactions is to set up an affordable yet reliable computational protocol. In this context, the Mp2mod method was here adopted, and the heavy atom polarization exponents of the small 6-31G* basis set were purposely tuned to reproduce the intermolecular interactions occurring between eumelanin’s smallest building blocks.

Exploiting the computational feasibility of the tuned basis set, the IPES of several dimer pairs was investigated. From the information retrieved, some general conclusions on the type and extent of the most likely non-covalent interactions characterizing eumelanin structure may be drawn. First of all, the stacking interactions are dominant, even if some hydrogen bonds can occur in the considered pairs. The latter indeed enforce dimer

stability but are not able to alter the preferred stacked disposition. More specifically, most of the investigated building block pairs are found with the molecular planes in a parallel arrangements, though often twisted or slightly displaced. The computed interaction energy are found to be remarkably dependent from monomer's relative orientation, spanning a rather wide range of values, from 5 to 12 kcal/mol. It might be thus speculated that such strong stacked interactions, sometimes complemented by non-negligible hydrogen bonds, should be able to promote the layered structure hypothesized for eumelanin by many authors. As a matter of fact, the extended IPES sampling carried out herein, revealed that for the most stable local minima the distance between two stacked rings of the considered building blocks is found between 3.0 and 3.5 Å, in agreement with the 3 – 4 Å range proposed^{7,36,40,41} for eumelanin inter-layer distance.

Once the presence of closely stacked aromatic planes has been confirmed by calculations, its consequences in the optical properties of the resulting non-covalent dimers were investigated. As a matter of fact, considering the rather small values found for the inter-ring separation, it can be hypothesized that the π clouds that characterize each of the two aromatic moieties will sensibly overlap, hence affecting the frontier orbitals involved in the electronic transition. Such supposition was confirmed by observing, for most of the considered pairs, the arising of a broad band upon dimer formation, whose position and intensity remarkably depend on dimer type and spatial arrangement. On the contrary, when the intermolecular interactions are completely neglected and the absorption spectrum is built as a superimposition of all the signals arising from the isolated monomers, a negligible absorption intensity is found in the 350-550 nm range, in contrast with the well known broad band features of real eumelanin. On the other hand, when the contribution of the stacked non-covalent dimers conformations is included in the calculation of the spectrum, the averaged computed signal shows two specific features that characterize eumelanin experimental absorption: a broad band that covers practically the whole investigated spectrum and a monotonic increase towards higher energies.

Despite these results seem to suggest that many of the considered configurations do effectively take place in eumelanin structure, the covalent bonds that may connect the investigated building blocks to form small protomolecules (as tetrameric structures) should

be taken into account. Indeed, work in this direction is currently in progress in our group. In this context, the validation of a reliable and effective method to compute the non-covalent interactions in eumelanin, taking place either between the simple building blocks investigated herein or between larger DHI based oligomers, paves the way to a further step in the comprehension of eumelanin supramolecular motifs.

It should be thus pointed out that this is the first part of a multi-step work, which will hopefully lead us toward an accurate large scale simulation of a model eumelanin material. Indeed, in the present work the interaction between the minimal building blocks was investigated and the proposed computational method reliably validated: a second step will consist in extending the investigation to the interaction between larger oligomers (protomolecules). This could be done, for instance, by resorting to the idea that the interaction energy between two large protomolecules (for instance two DHI tetramers) can be expressed as a sum of the interactions between all pairs of their building blocks.⁷⁰ In this hypothesis, the interaction energy between two protomolecules would strongly depend on their building blocks relative distance and orientation. However, unlike the case here investigated, the intramolecular flexibility of the involved oligomers should be also accounted for, since the constraints imposed by the structure of each protomolecule could sensibly alter the distribution of geometrical conformations of all possible building blocks pairs. Beside the protomolecules dimensions, a further obstacle to such an approach is also that the same level of accuracy is required to handle intramolecular flexibility and intermolecular interactions, if a correct balance has to be achieved. Therefore, the next step of this work will consist in applying to different kinds of protomolecules a number computational protocols, developed in our group⁷⁰⁻⁷³ to reliably investigate Van der Waals dimers of large dimensions.

Finally, an accurate sampling of the six dimensional interaction potential energy surface (IPES) of the building blocks pairs here considered could be also exploited to validate the performances of the FFs adopted in previous studies or employed as reference to parameterize a specifically tailored FF^{66,67,74,75} Once both the achieved intra- and intermolecular FF description will be validated, the last step of this project will consist in large-scale simulation on different eumelanin models, based on different combinations of

DHI building blocks.

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