

1 **Transdimensional and Hamiltonian Monte Carlo inversions of Rayleigh wave dispersion**
2 **curves: A comparison on synthetic datasets**

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ABSTRACT

10 We compare two Monte Carlo inversions that aim to solve some of the main problems of dispersion
11 curve inversion: deriving reliable uncertainty appraisals, determining the optimal model
12 parameterization, and avoiding entrapment in local minima of the misfit function. The first method
13 is a transdimensional Markov Chain Monte Carlo that considers as unknowns the number of model
14 parameters, that is the locations of layer boundaries, together with the V_s and the V_p/V_s ratio of
15 each layer. A reversible jump Markov Chain Monte Carlo (rjMCMC) algorithm is used to sample
16 the variable-dimension model space, while the adoption of a parallel tempering strategy and of a
17 delayed rejection updating scheme improves the efficiency of the probabilistic sampling. The
18 second approach is a Hamiltonian Monte Carlo (HMC) inversion that considers the V_s , the V_p/V_s
19 ratio and the thickness of each layer as unknowns, whereas the best model parameterization
20 (number of layer) is determined by applying standard statistical tools to the outcomes of different
21 inversions running with different model dimensionalities. This work has a mainly didactic
22 perspective and for this reason, we focus the attention to synthetic examples in which only the
23 fundamental mode is inverted. We perform what we call semi-analytical and seismic inversion tests
24 on 1D subsurface models. In the first case, the dispersion curves are directly computed from the
25 considered model making use of the Haskell-Thompson method, while in the second case they are
26 extracted from synthetic shot gathers. To validate the HMC and rjMCMC outcomes we analyse the

27 estimated posterior models and we also perform a sensitivity analysis in which we compute the
28 model resolution matrices, posterior covariance matrices, and correlation matrices from the
29 ensembles of sampled models. Our tests demonstrate that major benefit of the rjMCMC inversion is
30 its capability of providing a parsimonious solution that automatically adjusts the model
31 dimensionality. The downside of this approach is that many models must be sampled to guarantee
32 accurate posterior uncertainty. Differently, less sampled models are required by HMC algorithm,
33 but its limits are the computational effort related to the Jacobian computation, and the multiple
34 inversion runs needed to determine the optimal model parameterization.

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36

INTRODUCTION

37 Rayleigh wave measurements are highly sensitive to the S-wave velocity (V_s) and for this reason
38 they are extensively used for geotechnical characterization or seismic site response studies (Socco
39 and Strobbia 2004). Over the last years, the full-waveform inversion of surface waves is getting
40 growing attention thanks to the increased computational power of modern parallel architectures
41 (Gross et al. 2017; Xing and Mazzotti 2019). However, well-established methods still rely on
42 dispersion curve inversion under the assumption of a 1D subsurface structure (Socco and Boiero
43 2008; Maraschini and Foti 2010; Cercato, 2011; Foti et al. 2018; Di Giulio et al. 2019).

44 The dispersion curve inversion is a highly non-linear and ill-conditioned problem. For this reason, it
45 is crucial adopting inversion approaches that efficiently converge toward the global minimum, that
46 allow for a straightforward introduction of a-priori model constrains, and that also provide reliable
47 estimations of the uncertainties affecting the recovered solution (i.e. the estimation of the posterior
48 probability density “PPD”). In this context, gradient-based inversion methods (i.e. Gauss-Newton,
49 Steepest descent) exhibit fast convergence rates but limited capability to explore the parameter
50 space, resulting in a final solution highly dependent on the initial model (Luke et al. 2003; Wathelet
51 2005; Kritikakis et al. 2014). Metaheuristic search algorithms (i.e. genetic algorithms, particle
52 swarm optimization) exhaustively explore the model space but they usually require many forward

53 model evaluations and considerable computational effort to converge (Dal Moro et al. 2007; Picozzi
54 and Albarello 2007; Sen and Stoffa 2013; Sajeve et al. 2014; Sajeve et al. 2017). In addition, both
55 gradient-based and metaheuristic approaches do not satisfy the importance sampling principle and
56 for this reason they hamper a reliable uncertainty appraisal (Tarantola, 2005). Markov Chain Monte
57 Carlo (MCMC) algorithms exhibit global convergence capabilities and honour the importance
58 sampling principle (Sen and Stoffa, 1996) but they usually require a higher computational effort
59 with respect to linearized, gradient-based inversion algorithms.

60 Markov Chain Monte Carlo methods convert the inversion into a sampling problem in which the
61 sampling density is proportional to the posterior (Sambridge and Mosegaard, 2002). The first stage
62 of the MCMC sampling is the burn-in period in which the chain moves from a random starting
63 model to a high-probability region. The second stage is often called the sampling stage in which the
64 small fluctuations of the misfit value indicate that the MCMC algorithm reaches the stationary
65 regime. Usually the samples accepted during the burn-in period do not accurately represent the
66 target density and for this reason they are disregarded in the computation of the PPD.

67 Another issue of any geophysical inversion is determining the optimal model parameterization (i.e.
68 number of layers) that guarantees a good compromise between model resolution and posterior
69 uncertainty (Fernández-Martínez, 2015; Aleardi 2015; Menke 2018). A transdimensional MCMC
70 inversion treats the number of model parameters as an unknown to be inferred from the data (Green
71 1995; Malinverno 2000; Sambridge et al. 2006). In the past years, transdimensional MCMC
72 algorithms have been successfully applied to solve inverse problems especially at the earthquake
73 seismology scale (Bodin and Sambridge, 2009; Bodin et al. 2012), whereas over the last few years
74 their applications have also been extended to exploration and applied geophysics (Dadi et al. 2016;
75 Ray et al. 2017; Mandolesi et al. 2018; Xiang et al. 2018; Galletti and Curtis 2018; Zhu and Gibson,
76 2018; Cho et al. 2018). In the context of surface waves inversion, it is known (e.g. Socco and
77 Strobbia, 2004) that the chosen model parameterization (i.e. number of layers) heavily affects the
78 outcomes of dispersion curve inversion. For this reason, it is desirable to use an inversion approach

79 that automatically adapts the underlying model parametrization and that produces solutions with an
80 appropriate level of complexity to fit the data to statistically meaningful levels. However, specific
81 MCMC recipes targeted to the problem at hand must be implemented to maintain the computational
82 effort of the Markov chain sampling affordable. More specifically, MCMC methods are primarily
83 affected by low acceptance rates and show high correlations between successively sampled models.
84 Hamiltonian Monte Carlo (HMC; Duane et al. 1987) algorithm was designed to circumvent these
85 critical issues of MCMC algorithms. The main benefits of HMC with respect to standard MCMC
86 algorithms are its rapid convergence toward the stationary regime, so that the burn-in phase is
87 drastically reduced, and its ability to make long jumps in the model space, so that the independence
88 of the sampled models is guaranteed. HMC exploits the derivative information of the target PPD to
89 focus on the most promising regions of the model space, that is regions containing more plausible
90 models. HMC was developed for problems in which the derivative of the target probability density
91 with respect to the model parameters can be computed quickly (MacKay, 2003; Bishop 2006; Neal
92 2011; Betancourt 2017). More in detail, HMC treats a model as the mechanical analogue of a
93 particle that moves from its current position (current model) to a new position (proposed model)
94 along a given trajectory (Muir and Tkalcic, 2015; Sen and Biswas 2017; Fichtner and Simuté, 2018;
95 Fichtner et al. 2019). The geometry of the trajectory is controlled by the kinetic energy (K), by the
96 mass of the particle, and by the potential energy (U) that is interpreted as the misfit function. Note
97 that the kinetic energy and the mass matrix are artificially introduced as auxiliary quantities and
98 allow for the inclusion of the derivative information of the misfit function into the sampling
99 framework. In case of optimal tuning of the hyperparameters, the acceptance rate in standard
100 MCMC sampling (such as for the random walk Metropolis) lies between 0.2-0.4 (Sambridge and
101 Mosegard, 2002), while in case of optimal tuning the acceptance rate of HMC is very close to 0.6
102 (Neal, 2011).

103 In this work we compare the transdimensional MCMC and HMC algorithms for inverting Rayleigh
104 wave dispersion curves and we discuss some of their main benefits and drawbacks. The

105 transdimensional MCMC has been already applied for inverting surface waves and some examples
106 can be found in Bodin et al. (2012), Dettmer et al. (2012), and Galletti et al. (2016). Differently, this
107 is the first application of HMC to dispersion curve inversion and to the best of our knowledge this is
108 the first paper in which these two methods are extensively compared in solving this inverse
109 problem. For this reason, we develop a mainly didactic perspective without including any
110 methodological novelties to the rjMCMC and HMC approaches and by limiting the attention to
111 synthetic data inversions in which only the fundamental mode is considered.

112 We apply a birth-death reversible jump MCMC algorithm (rjMCMC; Geyer and Møller, 1994), in
113 which the amplitude of the variation of the number of dimensions is one (i.e. only one layer can be
114 added or deleted to the current model). Both algorithms invert for the shear wave velocity (V_s) and
115 the compressional wave velocity over the shear wave velocity (V_p/V_s) ratio. The rjMCM also
116 considers the number of layers and the depth location of the layer interfaces as additional unknowns
117 to be inverted for, whereas the HMC uses a fixed model parameterization but considers the layer
118 thickness as an unknown parameter. To avoid biased parameter estimations, we treat the V_p/V_s ratio
119 as an unknown, although this parameter exerts a small influence on the observed dispersion curves.
120 The density has an even smaller influence on the observed dispersion curve than the V_p/V_s ratio,
121 and for this reason this parameter is not inverted but fixed to a previously determined value.

122 We increase the efficiency of the implemented rjMCMC approach by applying a parallel tempering
123 strategy (Sambridge, 2014) and a delayed rejection updating scheme (Bodin and Sambridge 2009).
124 As previously mentioned, HMC requires the model parameterization (number of layers) as an input
125 to the inversion. For this reason, we here estimate the most probable model dimensionality by
126 making use of common statistical tools, such as the chi-squared (χ^2) probability and the Bayesian
127 information criterion “BIC” (Schwartz 1978; Ando 2010; Sambridge et al. 2006), which are
128 applied to models sampled in different inversion runs in which different model parameterizations
129 are employed.

130 For both the rjMCMC and HMC we perform what we call semi-analytical and seismic inversion
131 tests on 1D reference models. In the first case the observed data (dispersion curves) are semi-
132 analytically derived from the true model, while in the seismic tests the dispersion curves are
133 extracted from seismic shot gathers computed on the reference models by means of the reflectivity
134 algorithm (Mallick and Fraser 1987). In all the following experiments we limit to consider the
135 fundamental mode as the observed data, although it is known that higher modes are essentially to
136 better constraints the solution in case of shear velocity inversions and/or high stiffness contrasts
137 within the soil column (Feng et al. 2005; Luo et al. 2009; Cercato, 2011; Farrugia et al. 2016;
138 Sajeve and Menanno 2017; Qiu et al. 2019). We return to this aspect in more detail in the discussion
139 section. The extraction of the dispersion curves is always a very delicate step especially when
140 higher modes are more energetic than the fundamental one (Xia et al. 2003; Luo et al. 2009; Boiero
141 et al. 2011). In this work concerning synthetic data, we simply extract the observed data by picking
142 the maxima of the fundamental mode in the frequency-phase velocity spectra. In both the semi-
143 analytical and seismic inversions, the forward modelling is based on Haskell-Thompson (Haskell
144 1953) method that considers a stack of horizontal, homogeneous, layers. This method is also used in
145 the semi-analytical tests to compute the observed data from the reference models.

146 For two inversion examples we quantitatively validate the rjMCMC and HMC outcomes by
147 comparing the estimated marginal PPDs and by analysing some of the most popular sensitivity
148 analysis tools such as the model resolution matrices, posterior model covariance matrices, and
149 posterior correlation matrices (Menke, 2018), which are numerically computed from the ensemble
150 of sampled models. In addition, we relate these matrices and the estimated PPDs to the
151 geometrical properties of the misfit function (L2 norm difference between observed and predicted
152 dispersion curves).

153 In the following we start with a theoretical description of the rjMCMC and HMC approaches. Only
154 the main mathematical aspects of the methods are described, while the interested reader is referred
155 to pertinent references for further details. Then, we discuss the outcomes of the semi-analytical and

156 seismic inversion tests. Additional theoretical and practical insights into Monte Carlo inversion of
 157 geophysical data can be found in Mosegaard and Tarantola (1995), Sambridge and Mosegaard
 158 (2002), Mosegaard and Sambridge (2002), and Scalzo et al. (2019).

159 METHODS

160 Birth-death reversible jump Markov Chain Monte Carlo

161 The Bayes theorem can be written as:

$$162 \quad p(\mathbf{m}|\mathbf{d}) = \frac{p(\mathbf{d}|\mathbf{m})p(\mathbf{m})}{p(\mathbf{d})}, \quad (1)$$

163 or

$$164 \quad p(\mathbf{m}, n|\mathbf{d}) = \frac{p(\mathbf{d}|\mathbf{m}, n)p(\mathbf{m}|n)p(n)}{p(\mathbf{d})}, \quad (2)$$

165 where equation 1 represents the Bayesian equation for inverse problems with a fixed number of
 166 unknowns, while equation 2 is the Bayesian formulation in case of parameterizations with different
 167 number of unknowns n . In both equations 1 and 2, \mathbf{d} is the N -dimensional observed data vector, and
 168 \mathbf{m} is the Q -dimensional model parameter vector, whereas the left-hand side terms represent the
 169 target PPD that could be numerically estimated from the ensemble of models sampled by the
 170 MCMC algorithm after the burn-in phase. In a transdimensional MCMC, the Metropolis-Hasting
 171 rule determines the acceptance probability α , which is the probability to move from a model \mathbf{m} with
 172 dimension n to a model \mathbf{m}' with dimension n' at a given step of the chain:

$$173 \quad \alpha = p(\mathbf{m}', n'|\mathbf{m}, n) = \min \left[1, \frac{p(\mathbf{m}', n')}{p(\mathbf{m}, n)} \times \frac{p(\mathbf{d}|\mathbf{m}', n')}{p(\mathbf{d}|\mathbf{m}, n)} \times \frac{q(\mathbf{m}, n|\mathbf{m}', n')}{q(\mathbf{m}', n'|\mathbf{m}, n)} \times |\mathbf{J}| \right], \quad (3)$$

174 where \mathbf{J} is the Jacobian of the transformation from \mathbf{m} to \mathbf{m}' and is needed to account for the scale
 175 changes when the transformation involves a jump between models with different dimensions; $q()$ is
 176 the proposal distribution that defines the new model \mathbf{m}' as a random deviate from a probability

177 distribution $q(\mathbf{m}'|\mathbf{m})$ conditioned only on the current model \mathbf{m} . Note that the proposal ratio term in
 178 equation 3 vanishes if symmetric proposals (for example Gaussian proposals) are employed in a
 179 fixed-dimensional model space (i.e. \mathbf{m} and \mathbf{m}' have the same dimensions). In addition, for the
 180 birth-death rjMCMC algorithm the determinant of the Jacobian is equal to 1 and thus it can be
 181 conveniently ignored in computing the acceptance ratio (Bodin and Sambridge 2009).

182 We consider a Dirichlet prior distribution for interface locations and bounded and independent
 183 uniform priors for V_s and V_p/V_s ratio. We also assume a Gaussian distributed, uncorrelated noise
 184 with data covariance matrix equal to \mathbf{C}_d , whereas the likelihood function is based on a L2 norm
 185 misfit between observed and predicted data.

186 In particular, in the rjMCMC inversion the model vector \mathbf{m} includes the V_s and V_p/V_s ratio (in the
 187 following equations often indicated with γ for notational convenience) of the n layers and the vector
 188 $\mathbf{z} = [z_1, z_2, \dots, z_{n-1}]$ expressing the location of the of the $n - 1$ interfaces; $\mathbf{m} = [\mathbf{Vs}, \boldsymbol{\gamma}, \mathbf{z}] = [\mathbf{e}, \mathbf{z}]$.

189 The prior for model dimension is independent of the prior of layer properties and layer partitioning
 190 (distribution of interfaces depths over the depth range of interest), so that the prior model can be
 191 written as:

$$192 \quad p(\mathbf{m}, n) = p(\mathbf{e}|n)p(\mathbf{z}|n)p(n), \quad (4)$$

193 where $p(\mathbf{z}|n)$ is the prior on layer partitioning, $p(n)$ is the prior on the number of layers, whereas
 194 $p(\mathbf{e}|n)$ represents the prior on V_s and γ . By assuming uncorrelated layer properties, equation 4
 195 becomes:

$$196 \quad p(\mathbf{m}, n) = p(\mathbf{Vs}|n)p(\boldsymbol{\gamma}|n)p(\mathbf{z}|n)p(n), \quad (5)$$

197 The prior $p(n)$ is defined as a uniform distribution bounded in the interval $[n_{min}, n_{max}]$:

$$p(n) = \begin{cases} \frac{1}{\Delta n}, & \text{if } n_{min} \leq n \leq n_{max}, \\ 0, & \text{otherwise} \end{cases}, \quad (6)$$

198

199 where $\Delta n = n_{max} - n_{min}$. A uniform bounded prior for the interface depth within the partition is

200 given by the Dirichlet distribution (Steininger et al. 2013):

$$p(\mathbf{z}|n) = n! \mathbf{z}_{max}^{-n}, \quad (7)$$

201

202 where \mathbf{z}_{max} is the maximum allowed depth position of an interface.

203 We define the prior for each layer property (V_s and γ) to be a bounded uniform distribution

204 between $[e_{min}, e_{max}]$:

$$p(\mathbf{e}|n) = \begin{cases} \prod_{i=1}^n p(e_i|n) = \prod_{i=1}^n \frac{1}{\Delta e} = \frac{1}{(\Delta e)^n}, & \text{if } e_{min} \leq e \leq e_{max} \\ 0, & \text{otherwise} \end{cases}, \quad (8)$$

205

206 where \mathbf{e} generically represents both V_s and γ with $\Delta e = e_{max} - e_{min}$.

207 The assumption of uncorrelated V_s and V_p/V_s ratio is an oversimplification that results in an

208 augmented null-space of solutions and in an decreased convergence of the sampling. However, in

209 the following applications the limited number of model parameters (always less than 12) guarantees

210 the rapid convergence of the algorithm even though an oversimplified prior model is used.

211 If we assume Gaussian distributed, uncorrelated noise with a covariance matrix \mathbf{C}_d , the likelihood

212 function takes the following form:

$$p(\mathbf{d}|\mathbf{m}, n) = \frac{1}{\sqrt{(2\pi)^L |\mathbf{C}_d|}} \exp\left(-\frac{\varphi(\mathbf{m}, n)}{2}\right) \propto \exp\left(-\frac{\varphi(\mathbf{m}, n)}{2}\right), \quad (9)$$

213

214 where L is the number of data points, whereas $\varphi(\mathbf{m}, n) = \|\mathbf{C}_d^{-1/2}(\mathbf{d} - \mathbf{d}_{\text{pre}})\|_2^2$ is the weighted L2
 215 norm difference between observed and predicted data.

216 To increase the computational efficiency of the algorithm, we employ a parallel tempering strategy
 217 (Dosso et al. 2012; Sambridge, 2014) in which multiple and interactive chains are simultaneously
 218 run at different temperature levels $T=[T_1, T_2, \dots, T_{max}]$. High-temperature chains ensure a wide
 219 exploration of the model space, whereas low-temperature chains exploit the high-probability
 220 regions. According to stochastic criteria, swaps of models are allowed between chains at different
 221 temperatures, and in this context the high temperature chains ensure that low-temperature chains
 222 access all the high probability regions while maintaining an efficient exploitation capability. Only
 223 the models collected at $T=1$ are considered in the computation of the PPD because the models
 224 collected at $T>1$ sample a biased posterior probability density function. In this context the
 225 likelihood ratio of equation 3 becomes:

$$\frac{p(\mathbf{d}|\mathbf{m}', n', T)}{p(\mathbf{d}|\mathbf{m}, n, T)} = \exp\left(-\frac{\varphi(\mathbf{m}', n') - \varphi(\mathbf{m}, n)}{2T}\right). \quad (10)$$

226

227 In the approach to parallel tempering applied here, the two chains for a proposed swap are chosen at
 228 random from all chains at each iteration, and these chains are allowed to exchange their current
 229 models (or equivalently their current temperature levels). To derive the interchange probability let
 230 us consider two independent Markov chains with temperature T_i and T_j with current models \mathbf{m}_i and
 231 \mathbf{m}_j , respectively. At a given Monte Carlo step the two chains exchange their models with a
 232 probability equal to:

$$p(i, j) = \min \left[1, \exp \left[\left(\frac{1}{T_i} - \frac{1}{T_j} \right) (\varphi(\mathbf{m}_i) - \varphi(\mathbf{m}_j)) \right] \right], \quad (11)$$

233

234 and if the swap is accepted:

$$(\mathbf{m}_i, T_i), (\mathbf{m}_j, T_j) \rightarrow (\mathbf{m}_j, T_i), (\mathbf{m}_i, T_j). \quad (12)$$

235

236 In the dispersion curve inversion, the spread of the PPD is influenced by the model parameter
 237 illumination that, for example, decreases as the depth of investigation increases. In other terms, the
 238 optimal variance of the proposal for a given model parameter (e.g. V_s and V_p/V_s) is expected to
 239 vary with depth. For this reason, we increase the efficiency of the implemented rjMCMC sampling
 240 by using a delayed rejection scheme: after a rejected perturbation, a second attempt is made with a
 241 different proposal that could be theoretically dependent or independent from the previously rejected
 242 model. If the second move is independent from \mathbf{m}' we can simply write (Bodin and Sambridge
 243 2009):

$$p(\mathbf{m}''|\mathbf{m}) = \min \left[1, \frac{p(\mathbf{m}''|\mathbf{d})}{p(\mathbf{m}|\mathbf{d})} \frac{q_1(\mathbf{m}'|\mathbf{m}'')}{q_1(\mathbf{m}'|\mathbf{m}')} \frac{(1 - p(\mathbf{m}'|\mathbf{m}''))}{(1 - p(\mathbf{m}'|\mathbf{m}))} \right]. \quad (13)$$

244

245 In practice we derive \mathbf{m}'' by perturbing \mathbf{m} according to a second proposal distribution q_2
 246 characterized by a reduced variance with respect to the first proposal q_1 . To make the computation
 247 of the $p(\mathbf{m}''|\mathbf{m})$ straightforward, the delayed rejection scheme is here applied only to the fixed-
 248 dimension moves (see the following discussion). This strategy automatically adapts the
 249 characteristics of the proposal distribution to the spread of the PPD along different directions in the
 250 parameter space.

251 For even iteration numbers, we apply the following perturbation scheme:

252 • Property move: Randomly pick one layer and perturb its V_s and V_p/V_s values according to a
253 Gaussian proposal distribution with a null mean value and a previously selected variance.

254 For odd iteration numbers we select one of the following updating strategies with equal probability:

255 • Interface move: Randomly pick one interface and slightly perturb its position using a
256 Gaussian proposal with a zero mean and a previously selected variance. This move
257 guarantees a small variation in the likelihood between the current and the candidate models.

258 However, to prevent the creation of a too thin layer, if an interface is moved too close to
259 another interface, the move is rejected.

260 • Birth move: Add a new interface to the model at a given vertical location. Then, pick the V_s
261 and V_p/V_s values for the newly created layer from the corresponding prior models.

262 • Death move: choose at random one interface and remove it from the model. Then, randomly
263 select the layer above or below the deleted interface and assign its properties to the whole
264 vertical interval pertaining to the deleted layer.

265 It can be demonstrated that the previously described MCMC recipe results in an acceptance ratio for
266 all the moves equal to (see Dosso et al. 2014 and Xiang et al. 2018 for a complete mathematical
267 derivation):

$$p(\mathbf{m}', n' | \mathbf{m}, n) = \min \left[1, \frac{p(\mathbf{d} | \mathbf{m}', n')}{p(\mathbf{d} | \mathbf{m}, n)} \right]. \quad (14)$$

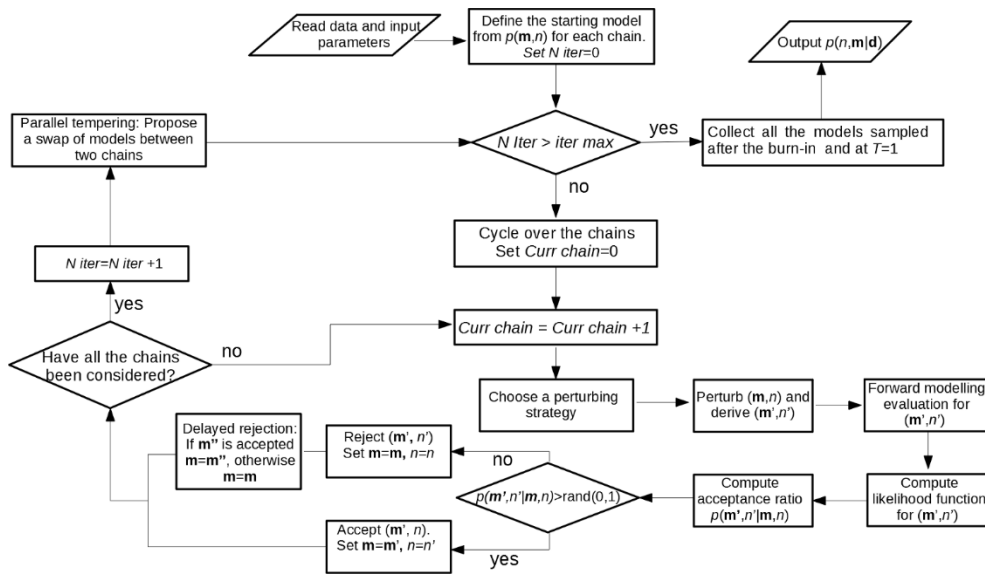
268

269 In the following examples we consider a uniform prior for V_s defined over [50 m/s, 400 m/s], a
270 uniform prior for V_p/V_s ratio defined over [1.8, 5.8], whereas the number of layers is bounded
271 between [2, 5], for a maximum allowed depth position for an interface equal to 20 m. The variances

272 of the proposals are properly set to have an acceptance probability between 0.2-0.4 for the chains
273 running at $T = 1$. For each chain the starting models are randomly generated from the prior. In all

274 the following experiments, we employ 20 tempered chains, with 10 chains at $T=1$ and the remainder

275 with logarithmically distributed temperature values between 1 and 100. The highest temperature
 276 values are set to obtain an acceptance probability around 0.6-0.7 for the corresponding chain. The
 277 number of iterations is set to 3000 with a burn-in period of 300. Figure 1 represents the flow-chart
 278 of the rjMCMC algorithm.
 279



280
 281 *Figure 1: A schematic representation of the rjMCMC inversion. $N iter$, $iter max$, and $Curr chain$*
 282 *identify the current iteration number, the maximum number of iterations, and the considered chain,*
 283 *respectively.*

284

285 Hamiltonian Monte Carlo

286 The Hamiltonian Monte Carlo is another method to numerically solve the Bayes theorem (equation
 287 1). HMC improves upon random walk Metropolis by proposing states that are distant from the
 288 current one, but nevertheless have a high probability of acceptance (Duane et al. 1987; Neal, 1996;
 289 Neal 2011; Betancourt 2017). These distant proposals are found by numerically simulating the
 290 Hamiltonian dynamics, in which a model is viewed as a moving particle with a physical state
 291 uniquely determined by the position and momentum vectors: These vectors define the phase space.
 292 HMC samples an auxiliary distribution defined over the $2Q$ -dimensional phase space, from which

293 samples of the posterior can be obtained by ignoring the momentum space component. The particle
 294 trajectory in the phase space is determined by the potential energy (U), the kinetic energy (K) and
 295 the mass matrix (\mathbf{M}). The potential energy is the negative natural logarithm of the posterior (see
 296 equation 1) or in other terms is the misfit function associated to the inverse problem. In this context
 297 more plausible models with large values of the posterior are associated to low potential energies.
 298 Generally speaking, the potential energy is given by:

$$U(\mathbf{m}) = -\ln(p(\mathbf{m}|\mathbf{d})). \quad (15)$$

299
 300 HMC determines the kinetic energy by introducing an auxiliary variable (momentum variable) \mathbf{p}
 301 defined over a Q -dimensional space. It is usually assumed that the auxiliary momentum variable has
 302 a multivariate normal distribution with zero mean and a covariance matrix equal to the so-called
 303 mass matrix:

$$K(\mathbf{p}) = N(\mathbf{p}; \mathbf{0}, \mathbf{M}), \quad (16)$$

304
 305 where N represents the Gaussian distribution, and \mathbf{M} is the $Q \times Q$ mass matrix that must be
 306 accurately set to ensure the convergence of the HMC algorithm (see Fichtner et al. 2019). The
 307 vectors \mathbf{p} and \mathbf{m} define the $2Q$ -dimensional phase space, whereas the Hamiltonian $H(\mathbf{p}, \mathbf{m})$ is the
 308 total energy of the particle:

$$H(\mathbf{p}, \mathbf{m}) = U(\mathbf{m}) + K(\mathbf{p}). \quad (17)$$

310 After defining the kinetic and potential energies, the Hamiltonian dynamics can be simulated. In
 311 this context, the model \mathbf{m} moves through the $2Q$ -phase space according to Hamilton's equations:

$$\frac{d\mathbf{m}_i}{dt} = \frac{\partial K}{\partial \mathbf{p}_i}, \quad \frac{d\mathbf{p}_i}{dt} = -\frac{\partial U}{\partial \mathbf{m}_i}, \quad \text{with } i = 1, 2, \dots, Q, \quad (18)$$

312
 313 where t indicates the artificially introduced time variable. Note that the rightmost term of equation
 314 18 contains the partial derivative of the potential energy (i.e. the misfit function) with respect to the

315 considered model \mathbf{m} . This makes it possible introducing information about the gradient of the misfit
316 function into the Monte Carlo sampling framework.

317 For each current model \mathbf{m} , and for each iteration, HMC executes the following steps:

- 318 1. Randomly draw the Q momenta \mathbf{p} from the normal distribution $N(\mathbf{p}; \mathbf{0}, \mathbf{M})$;
- 319 2. Derive the proposed model $\mathbf{m}(t)$ and the new momenta $\mathbf{p}(t)$ by solving Hamilton's
320 equations (18) for a given propagation time t . In this work we use the leap-frog method as
321 the numerical integration method (Betancourt, 2017);
- 322 3. Accept the proposed model with probability α :

$$323 \quad \alpha = \min \left[1, \frac{\exp [-H(\mathbf{p}(t), \mathbf{m}(t))]}{\exp [-H(\mathbf{p}, \mathbf{m})]} \right]. \quad (19)$$

324 If accepted, the proposed $\mathbf{m}(t)$ point constitutes the starting model for the next trajectory
325 ($\mathbf{m} = \mathbf{m}(t)$). Otherwise, the current model \mathbf{m} is again used as the starting point in the
326 following iteration;

- 327 4. Return to step 1.

328

329 Note that standard MCMC algorithms explore model space only slowly compared to HMC, because
330 in HMC all model parameters are updated at each iteration, so that long distances in phase space
331 can be traversed with a single move. This promotes a high level of acceptance and independence of
332 the sampled models. Differently, in standard MCMC only subsets of parameters are updated in each
333 iteration. This maintains a relatively high level of acceptance but at the expense of a high degree of
334 correlation between successively sampled models. For a more detailed comparison of MCMC and
335 HMC see, for example, Fichtner et al. (2019).

336 The model vector in the implemented HMC approach includes the V_s , the V_p/V_s ratio and the layer
337 thicknesses h of each layer: $\mathbf{m} = [\mathbf{V}_s, \boldsymbol{\gamma}, \mathbf{h}]$. We use equation 16 to define the momentum

338 distribution. Since it is known that standard HMC exhibits poor performance in sampling
 339 multimodal target densities (see the discussion section), we simply employ a Gaussian prior model.
 340 In this context the prior model can be compactly written as follows:

$$p(\mathbf{m}) = N(\mathbf{m}; \mathbf{m}_{prior}, \mathbf{C}_m), \quad (20)$$

341
 342 where \mathbf{m}_{prior} and \mathbf{C}_m are the a-priori mean vector and model covariance matrix that can be derived,
 343 for example, from previous knowledge about the investigated area. In the following inversion
 344 examples, we assume an uncorrelated and depth independent prior model (i.e. \mathbf{C}_m is diagonal).

345 In this work the potential energy is defined as:

$$U(\mathbf{m}) = -\ln(p(\mathbf{m}|\mathbf{d}))$$

$$= \frac{1}{2} \left[(\mathbf{d} - G(\mathbf{m}))^T \mathbf{C}_d^{-1} (\mathbf{d} - G(\mathbf{m})) + (\mathbf{m} - \mathbf{m}_{prior})^T \mathbf{C}_m^{-1} (\mathbf{m} - \mathbf{m}_{prior}) \right], \quad (21)$$

346
 347 The partial derivative of equation 21 with respect to the model parameter vector \mathbf{m} is equal to:

$$\frac{\partial U(\mathbf{m})}{\partial \mathbf{m}} = \mathbf{J} \mathbf{C}_d^{-1} (\mathbf{d} - G(\mathbf{m})) + \mathbf{C}_m^{-1} (\mathbf{m} - \mathbf{m}_{prior}), \quad (22)$$

348
 349 In equations 21 and 22, G is the non-linear forward modelling that computes the dispersion curves
 350 from the current model (i.e. the Haskell-Thompson method), \mathbf{C}_d is the data covariance matrix that
 351 expresses the error affecting the observed dispersion curve (under the assumption of a Gaussian
 352 distribution). \mathbf{J} is the Jacobian matrix (also called Fréchet derivative matrix) that we compute with a
 353 central finite difference (FD) approach (Aster et al. 2018). We use the FD approach for its
 354 straightforward implementation but if needed other more sophisticated and more efficient methods
 355 can be used (Sen and Roy, 2003).

356 A proper setting of the mass matrix (\mathbf{M}) is of crucial importance in any HMC sampling (see for
357 example Fichtner et al. 2019) because it serves as a tuning parameter to adjust the speed with which
358 the phase space is traversed, or in other words it is used to decorrelate the target distribution
359 (Betancourt, 2017). In practice, a proper setting of this matrix maximizes the independency of the
360 sampled models, while preventing the exploration of similar model space regions. The optimal
361 setting of the mass matrix is strongly case dependent, but typical applications use a diagonal \mathbf{M}
362 defined as a scalar multiple of the identity matrix. However, this strategy often hampers an efficient
363 sampling of the parameter space. In this work, the mass matrix is computed as a local
364 approximation (around the considered model) of the inverse of the posterior covariance matrix (see
365 Fichtner et al. 2019):

$$\mathbf{M} = \mathbf{J}^T \mathbf{C}_d^{-1} \mathbf{J} + \mathbf{C}_m^{-1}, \quad (23)$$

366
367 In the leapfrog method, we update the momentum and position variables sequentially. First, we
368 simulate the momentum dynamics (changing momentum) by $\delta/2$ time units, then simulating the
369 position dynamics (moving in model space) for full δ time units, then again completing the
370 momentum simulation for the remaining half-time units, $\delta/2$, so that the momentum and model
371 perturbations can be at the same time of δ time units. This process is repeated for a total of L times
372 after which the algorithm reaches a new state. The leapfrog method integration is described by the
373 following equations:

$$p_i(t + \delta/2) = p_i(t) - \frac{\delta}{2} \times \frac{\partial U}{\partial m_i} \Big|_t, \quad (24)$$

374

$$m_i(t + \delta) = m_i(t) + \delta \times \frac{\partial K}{\partial p_i} \Big|_{t+\frac{\delta}{2}}, \quad (25)$$

375

$$p_i(t + \delta) = p_i\left(t + \frac{\delta}{2}\right) - \frac{\delta}{2} \times \frac{\partial U}{\partial m_i} \Big|_{t+\delta}, \quad (26)$$

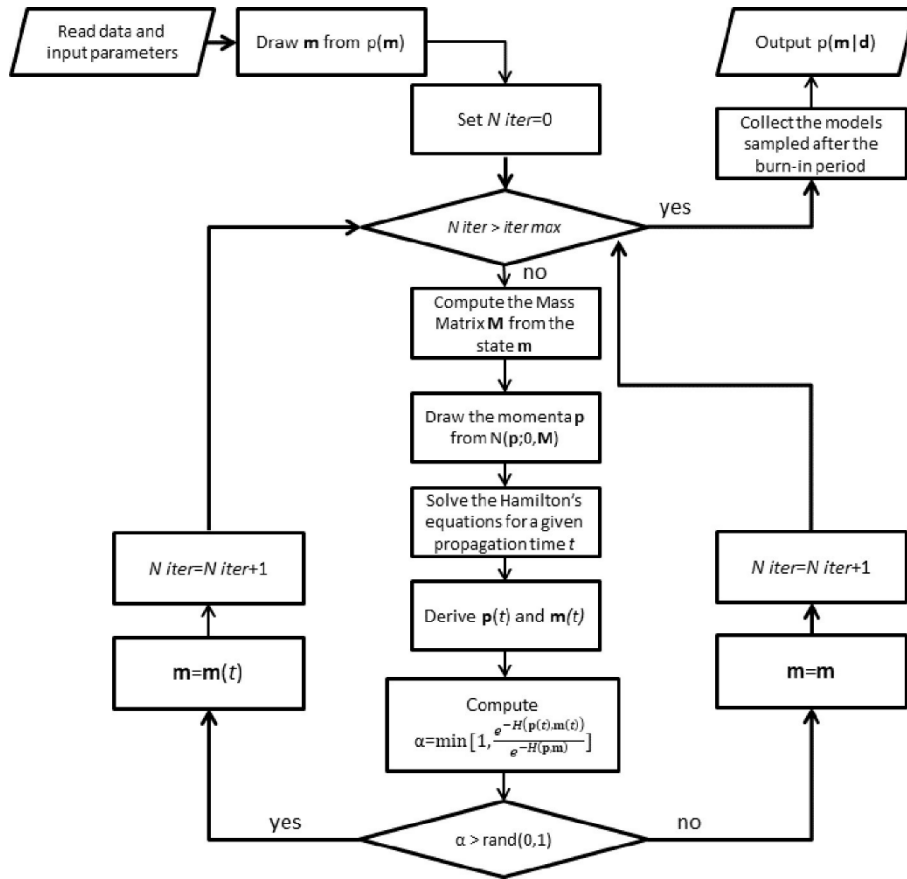
376

377 where t is the time variable. Although, many other integrator methods exist (see for example Blanes
378 et al. 2014), here we use the leapfrog method for its easy implementation.

379 In addition to the mass matrix, also the choice of L and δ plays an important role in the efficiency
380 of the sampling and for this reason these parameters need to be tuned properly to get the desired
381 acceptance rate. In particular, they determine the sampling trajectory in the augmented model space.
382 A too short trajectory generates proposal models very close to the current position, thus slowing
383 down the exploration of the model space. In contrast, if the trajectory is too long the algorithms
384 revisits points in model space that have been already sampled, thereby wasting computing time. To
385 avoid locking in periodic trajectories we follow the strategy discussed in Mackenze (1989) and in
386 each HMC iteration we draw the L parameter randomly from a previously defined uniform
387 distribution (see the discussion section for more details). More theoretical details about HMC
388 inversion can be found in Neal (2011) and Betancourt (2017).

389 Differently from the rjMCMC, the model dimensionality (layer number) is fixed in an HMC
390 inversion run. For this reason, we propose to determine the most probable number of layers by
391 performing different HMC inversions with different parameterizations. Then, standard statistical
392 tools, such as the χ^2 probability and the BIC, can be used to determine the most probable model
393 dimensionality. In the following examples we use a single HMC chain running for 3000 iterations,
394 for a burn-in period of 30 samples. Similarly to the rjMCMC, the starting points for the HMC
395 inversions are randomly generated from the prior distribution. Figure 2 shows the schematic flow-
396 chart of the implemented HMC inversion.

397

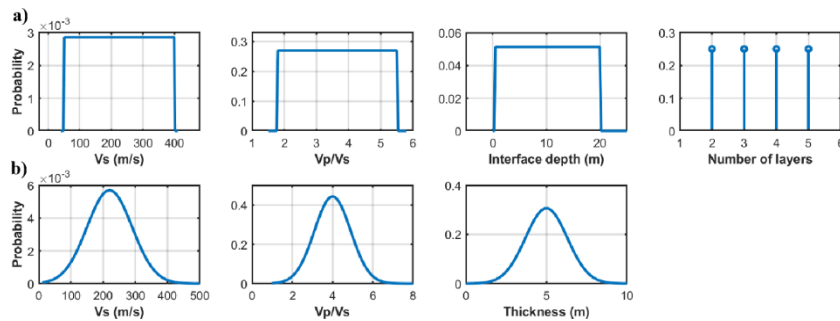


398

399 *Figure 2: The schematic workflow of the HMC inversion. N_{iter} and $iter_{max}$ identify the current*
 400 *iteration number and the maximum number of iterations, respectively.*

401 SEMI-ANALYTICAL INVERSIONS

402 We represent in Figure 3 the prior models for the rjMCMC and HMC algorithms. Note that these
 403 priors are depth independent and different for the two inversion approaches: uniform for rjMCMC
 404 and Gaussian for HMC.



405

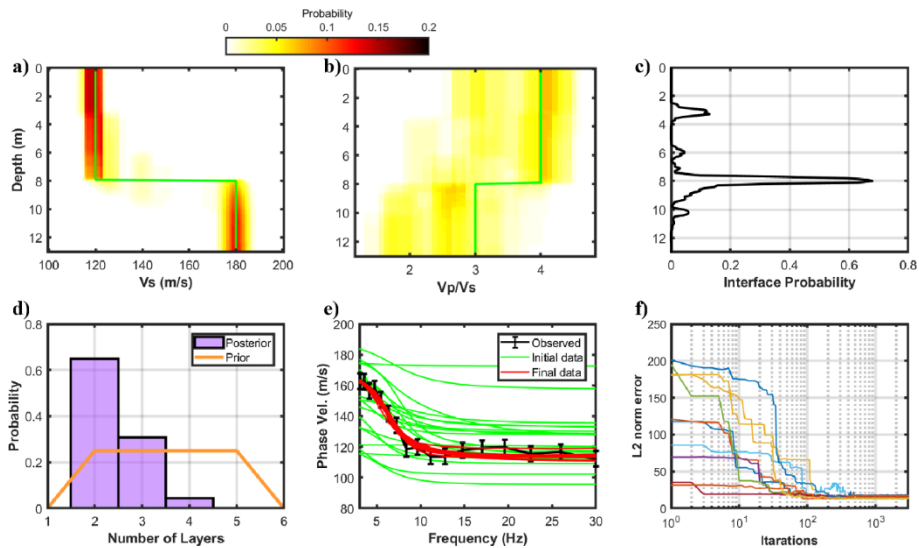
406 *Figure 3: a) Prior distributions of V_s , V_p/V_s ratio, interface depth, and number of layers for the*
 407 *rjMCMC algorithm. b) Prior distributions for the HMC algorithm projected onto the V_s , V_p/V_s , and*
 408 *layer thickness axes.*

409 The use of uniform and Gaussian priors highly simplifies the derivation of the acceptance ratio in
 410 the rjMCMC and of the posterior model in the HMC inversions, respectively. From a theoretical
 411 point of view the use of different priors complicates a direct comparison of the posterior models
 412 provided by the two inversion algorithms. Indeed, the posterior tends to the prior as the likelihood
 413 function becomes less informative. Despite this difference, we will see that the two methods
 414 provide PPD estimations congruent with the likelihood function and with the considered prior
 415 models. In particular, they yield similar posterior estimations for the model parameters better
 416 constrained by the data. For detailed investigation of the influence played by the prior model in a
 417 Bayesian inversion we refer the reader to Malinverno and Briggs (2004), and Theune et al. (2010).

418

419 **Transdimensional Markov Chain Monte Carlo inversions**

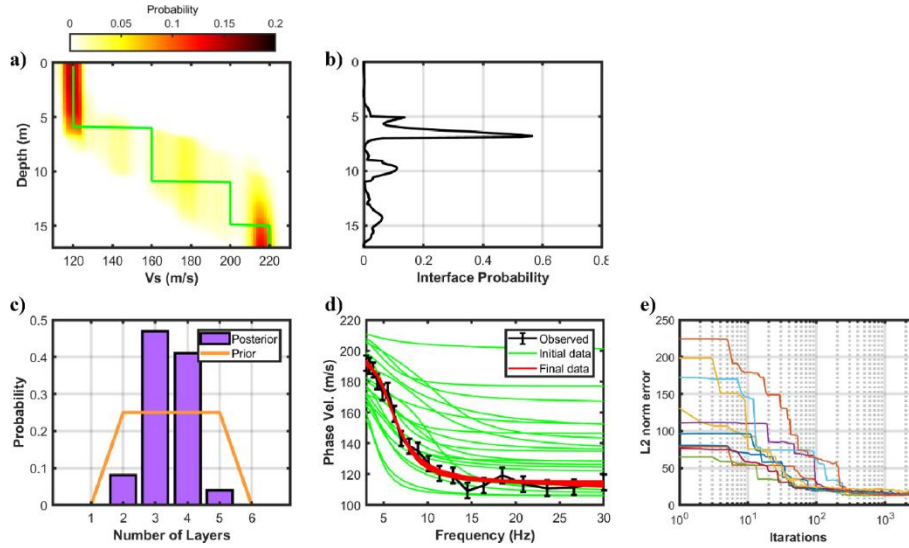
420 In this section we consider two schematic subsurface models with 2 and 4 layers, respectively. The
 421 frequency range of the data lies in the interval [4-30 Hz], whereas the noise standard deviation is
 422 assumed known and fixed to a value of 5 for all the frequencies.



423

424 *Figure 4: Semi-analytical inversion results provided by the rjMCMC algorithm for the 2-layer*
 425 *model. a) True Vs model (green line) and marginal PPD for Vs (colour scale). b) True Vp/Vs model*
 426 *(green line) and marginal PPD for Vp/Vs ratio (colour scale). c) Marginal PPD for interface*
 427 *location. d) Prior and posterior probability density functions for the number of layers. e)*

428 Comparison between the observed noisy data (black line), the data generated on the starting
 429 models (green lines) and the data generated on the last sampled models (red line). The black bars
 430 represent the noise standard deviation. f) Example of evolution of the L2 norm misfit for 10 chains
 431 out of 20.



432
 433 Figure 5: Semi-analytical inversion results provided by the rjMCMC algorithm for the 4-layer
 434 model. a) True V_s model (green line) and marginal PPD for V_s (colour scale). b) Marginal PPD for
 435 interface location. c) Prior and posterior probability density functions for the number of layers. d)
 436 Comparison between the observed noisy data (black line), the data generated on the starting
 437 models (green lines) and the data generated on the last sampled models (red line). The black bars
 438 represent the noise standard deviation. e) Example of evolution of the L2 norm misfit for 10 chains
 439 out of 20.

440 We start from a very schematic model constituted by two layers separated by an interface located at
 441 8 m depth. The marginal PPDs of V_s and interface location show that the inversion recovers the
 442 interface positions, the V_s of the first layer and that of the half-space (Figure 4). Conversely, as
 443 expected, the estimated V_p/V_s ratio is affected by much greater ambiguity. For this reason, in many
 444 of the following examples the estimated V_p/V_s ratio is not discussed. The algorithm has reliably
 445 sampled the PPD for the number of layers and yields a maximum-a-posteriori (MAP) solution that

446 matches the correct value of 2. The data comparison in Figure 5d demonstrates that the algorithm
447 successfully converges toward the so-called equivalence region of solutions. Finally, Figure 4f
448 proves that the rjMCMC reaches the stationary regime after 200 iterations approximately, when the
449 L2 errors of the different chains start to oscillate around stable values.

450 The second example considers a 4-layer model (Figure 5). As expected, we observe an increase of
451 uncertainties with respect to the previous test, related to the increased ill-conditioning of the
452 inversion procedure. The V_s of the shallowest layer and the position of the shallowest interface are
453 recovered with high accuracy, while the uncertainties on the estimated V_s and on the interface
454 positions increase for the deepest layers and deepest interfaces. The uncertainties affecting the
455 model parameterization (Figure 5c) increase with respect to the previous example. In this case both
456 the 3-layer and the 4-layer parameterizations exhibit similar posterior probability values, although
457 the MAP solution erroneously indicates that 3 is the most probable number layers, thus
458 underestimating the correct number of layers (4). This proves that in cases of more than 3 layers
459 and if we consider the fundamental mode only, the estimation of the number of layers becomes a
460 severely ill-conditioned problem. Differently, the inversion still provides satisfactory shear wave
461 velocity predictions. Figure 5d proves that the models sampled by the rjMCMC algorithm
462 successfully predict the observed data, whereas the evolution of the L2 norm data misfit (Figure 5e)
463 shows that the stationary regime is reached after 300 iterations, approximately.

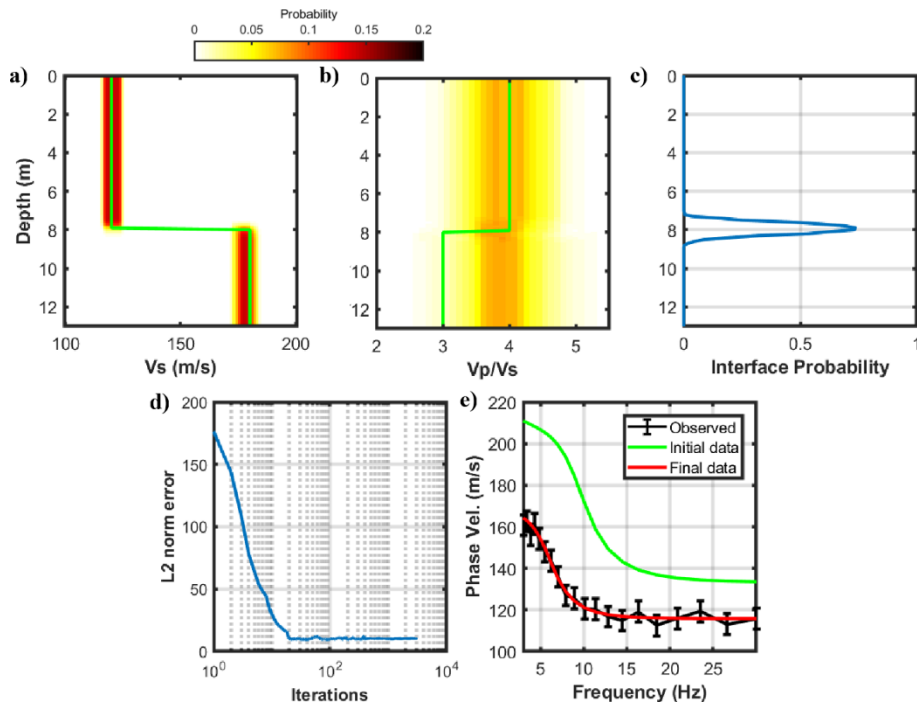
464

465 **Hamiltonian Monte Carlo inversion**

466 In this section we again consider the 2- and 4-layer models already used in the rjMCMC inversion.
467 In the 2-layer model example our aim is to compare the uncertainties affecting the final solution
468 when the dispersion curves lie in different frequency bands. In the first test the dispersion curve
469 extends over [3-30 Hz], whereas in the second case the observed data lie in the interval [6-30 Hz].
470 Both examples consider a correct number of layers equal to 2. In Figure 6 we represent the results

471 for the 3-30 Hz example. The marginal PPDs of V_s and interface location show that the inversion
 472 correctly and accurately estimates the true model. Differently the V_p/V_s ratio is not recovered, as
 473 demonstrated by a posterior distribution still very similar to the prior with a depth-independent
 474 MAP value equal to 4 (see Figure 6b). Note the fast convergence rate and that less than 20 iterations
 475 are needed to reach the stationary regime. The data comparison in Figure 6e demonstrates that the
 476 algorithm successfully predicts the observed data. Figure 7 shows the results for the 6-30 Hz
 477 example. We observe that the V_s of the first layer has been recovered with the same accuracy of the
 478 previous example. Differently, the position of the interface and particularly the velocity of the
 479 deepest layer are now estimated with higher uncertainties. This is mainly related to the lack of low
 480 frequencies in the observed data that are crucial to constraint the V_s of the deepest layer (Socco and
 481 Strobbia, 2004).

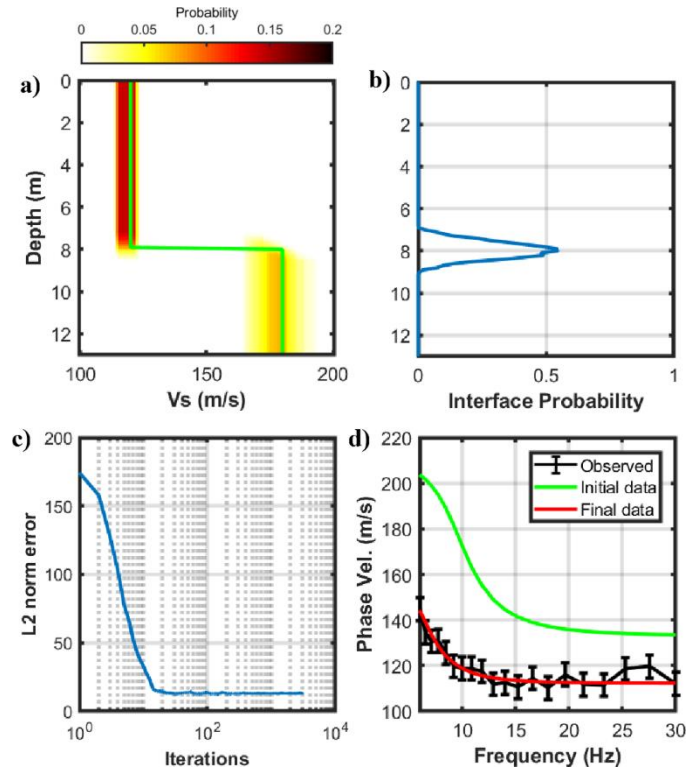
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483

484 *Figure 6: Semi-analytical inversion results for the 2-layer model provided by the HMC algorithm in*
 485 *the frequency range [3-30] Hz. a) True V_s model (green line) and marginal V_s PPD (colour scale).*
 486 *b) True V_p/V_s model (green line) and marginal V_p/V_s PPD (colour scale). c) Marginal PPD for*
 487 *interface location. d) Evolution of the L2 norm misfit. e) Comparison between the observed noisy*

488 data (black line), the data generated on the starting model (green line) and the data generated on
 489 the last sampled model (red line). The black bars represent the noise standard deviation.



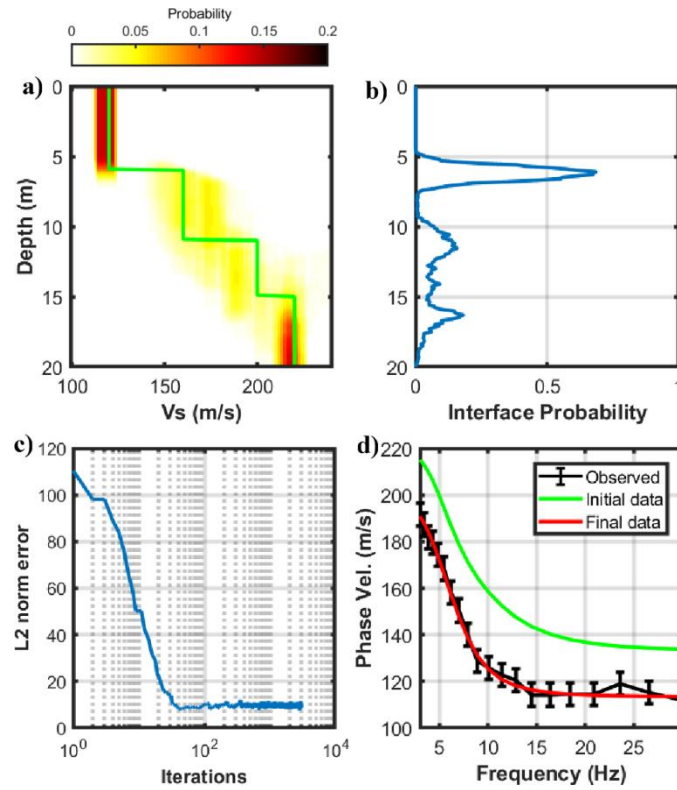
490
 491 *Figure 7: Semi-analytical inversion results for the 2-layer model provided by the HMC algorithm in*
 492 *the frequency range [6-30] Hz. a) True Vs model (green line) and marginal Vs PPD (colour scale).*
 493 *b) Marginal PPD for interface location. c) Evolution of the L2 norm misfit. d) Comparison between*
 494 *the observed noisy data (black line), the data generated on the starting model (green line) and the*
 495 *data generated on the last sampled model (red line). The black bars represent the noise standard*
 496 *deviation.*

497 We now perform different inversion runs with different model parameterizations (i.e. different
 498 number of layers) and considering a frequency band between 4 and 30 Hz. We compute the
 499 observed data on the 4-layer model already considered in the rjMCMC inversion (Figure 5), and we
 500 invert by considering a 4- and a 3-layer parameterization. Figure 8 shows that if the correct
 501 parameterization is employed, the inversion is still able to accurately predict the Vs of the two
 502 shallowest layers and to correctly locate the shallowest layer boundary. As expected, the shear wave

503 velocities of the intermediate layers and the position of the deepest interface are recovered with
 504 high uncertainties. If we consider the 3-layer parameterization, we are still able to accurately predict
 505 the velocity of the two shallowest layers and to correctly locate the positions of the two shallowest
 506 interfaces but the position of the deepest interface and the V_s of the third layer are not recovered.

	$p(\chi^2)$	BIC
2-layer parameterization	0.0003	38.93
3-layer parameterization	0.0386	21.93
4-layer parameterization	0.0199	32.30

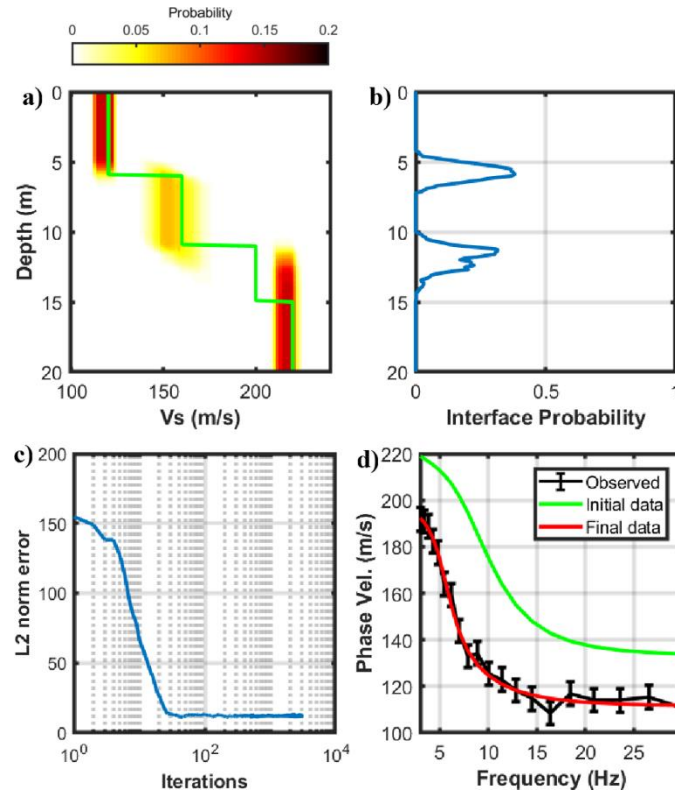
507 *Table 1: $p(\chi^2)$ and BIC values for the semi-analytical test resulting from different model*
 508 *parameterizations.*



509
 510 *Figure 8: Semi-analytical inversion results for the 4-layer model provided by the HMC algorithm*
 511 *for a 4-layer parameterization. a) True V_s model (green line) and marginal V_s PPD (colour scale).*
 512 *b) Marginal PPD for interface location. c) Evolution of the L2 norm misfit. d) Comparison between*
 513 *the observed noisy data (black line), the data generated on the starting model (green line) and the*

514 data generated on the last sampled model (red line). The black bars represent the noise standard
515 deviation.

516



517

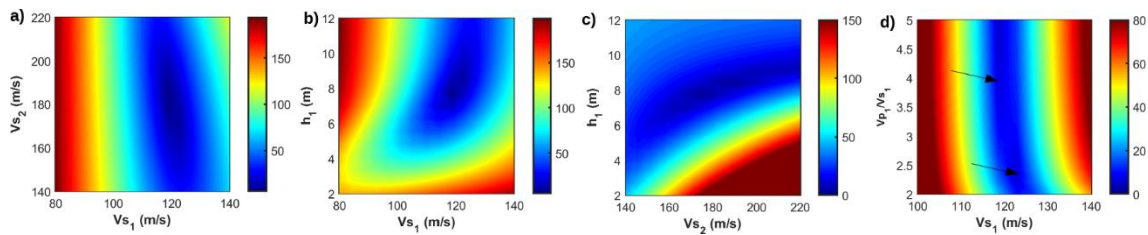
518 *Figure 9: As in Figure 8 but for a 3-layer parameterization.*

519 Table 1 shows the $p(\chi^2)$ and BIC values obtained in this semi-analytical test when different model
520 parameterizations are employed. The higher $p(\chi^2)$ and lower BIC values for the 3-layer
521 parameterization indicate that this is the most appropriate in this case. This result still proves that
522 estimating the correct number of layers for a 4-layer model from the fundamental model only is a
523 hopelessly ill-conditioned problem. Despite this limitation the Vs profile is still retrieved with
524 reasonable accuracy. Finally, in both Figures 8c and 9c note the fast convergence of the HMC
525 algorithm toward the stationary regime that is reached just after 20 iterations.

526

527 **Sensitivity analysis of the inversion results**

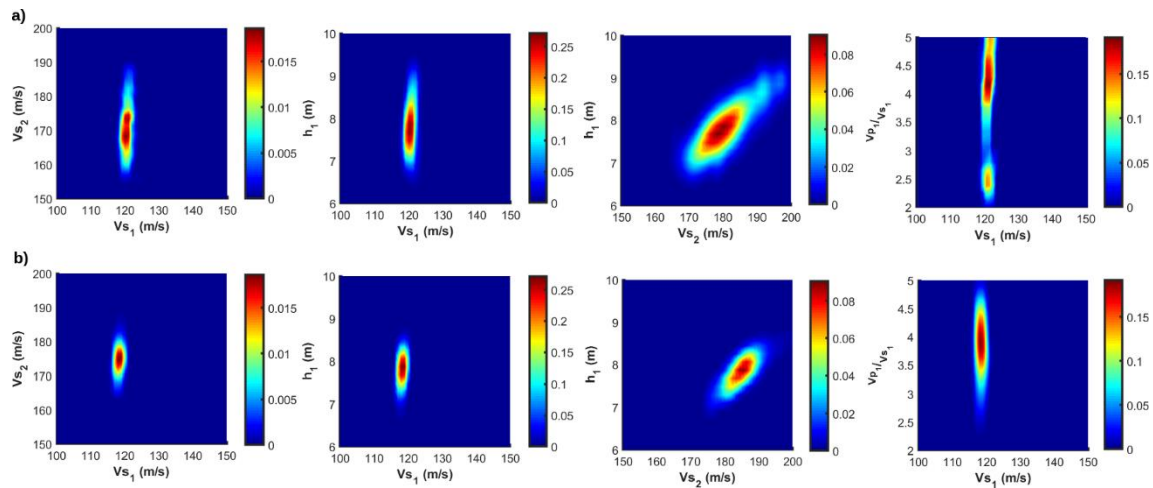
528 In this section we discuss in more detail the statistical properties of the ensemble of models sampled
529 by the HMC and rjMCMC algorithms. For brevity we limit to a single semi-analytical test running
530 on the previously considered 2-layer model and employing observed data lying in the same
531 frequency range (6-30 Hz). All the algorithm hyperparameters are the same previously used in the
532 2-layer examples. First, we compute the misfit function of the inversion procedure that is the L2
533 norm difference between observed and predicted dispersion curves. We remind that this misfit
534 function is directly related to the likelihood function of the Bayesian inversion through equation 9.
535 Figure 10 shows some examples of 2D projections of the misfit function, in which the elongated
536 valleys of minimum represent directions in the model space, and then model parameter
537 combinations, that weakly affect the observed data. In other terms the elongated valleys are aligned
538 with the eigenvectors in model space associated with the smallest singular values of the inversion
539 kernel (Fernandez Martinez et al. 2012).



540
541 *Figure 10: Some projections of the L2 norm misfit function onto different 2D sections. The*
542 *subscripts 1 and 2 refer to the first (shallowest) and second (deepest) layer, respectively, whereas h*
543 *indicates the layer thickness. The black arrows in d) indicate two local minima.*

544
545 As expected, the contours lines of the misfit function are very different from elliptic cylinders or
546 ellipses as it would be in a linear inverse problem (Fernandez Martinez et al. 2012). From the
547 inspection of Figure 10 we note that the dispersion curve is mostly influenced by the V_s of the
548 shallowest layer (V_{s1}), followed by the V_s of the deepest layer (V_{s2}), and by the thickness of the
549 first layer (h_1), while the V_p/V_s ratio plays a negligible role in determining the measured data. From
550 the analysis of the orientation of the misfit function we can claim that the shear wave velocities of

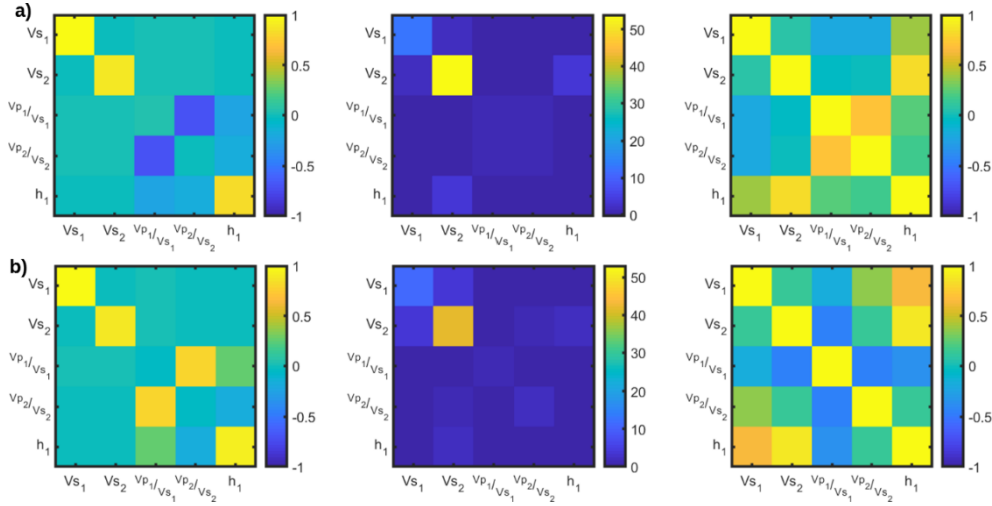
551 the two layers independently influence the data since the axes of the global minimum valley are
552 almost parallel to the model parameter axes (Figure 10a). In Figure 10b if we focus the attention to
553 the neighbourhood of the global minimum, we see that the V_s and the thickness of the first layer are
554 positively correlated. The orientation of minimum valley in Figure 10c indicates a strong positive
555 correlation between the V_s of the second layer and the thickness of the first one. In other terms, to
556 keep the data unchanged, an increase of the thickness of the first layer must be compensated by an
557 increase of the shear wave velocity of the underlying layer. In any case all these 3 parameters (V_{s1} ,
558 V_{s2} and h_1 in Figure 10) are informed by the data and none of them lies in the null space of
559 solutions. This means that the posterior model for these parameters will be mainly influenced by the
560 likelihood function instead by the a-priori assumptions. For this reason, in case of an accurate
561 sampling we expect that the geometrical characteristics of the misfit function influence the
562 statistical properties of the estimated PPDs. On the contrary, we observe in Figure 10d that the
563 V_p/V_s ratio is not illuminated by the data as the misfit function shows a very elongated valley with
564 two minima: one local minimum for a V_p/V_s ratio equal to 2.4, and the global minimum
565 corresponding to V_p/V_s ratio of the reference model equal to 4. In this case the two axes of the
566 minimum valley are aligned with the two considered model parameters and this means that the V_s
567 and the V_p/V_s of the first layer independently influence the observed dispersion curve. Similar
568 conclusions could have been drawn about the resolvability of the V_p/V_s ratio of the second layer
569 (not shown here for brevity). For this reason, we conclude that the V_p/V_s ratio lies in the null space
570 of solutions and for this reason we expect that the associated PPDs will be mainly determined by
571 the a-priori assumptions in case of an informative prior model (i.e. the Gaussian prior used by the
572 HMC algorithm), or by the likelihood function for an uninformative prior (i.e. the uniform prior
573 used by rjMCMC).



574

575 *Figure 11: 2D Marginal PPDs estimated by the rjMCMC (a) and by the HMC (b) algorithms.*

576 Figure 11 shows 2D marginal PPDs projected onto the same sections previously considered in
 577 Figure 10. Both the posterior models estimated by the rjMCMC and HMC algorithms confirm that
 578 the V_s of the first layer is the parameter mostly illuminated by the data, followed by the V_s of the
 579 second layer and the thickness of the first layer. As expected, the posterior estimations of these
 580 parameters are mainly driven by the likelihood function and for this reason both HMC and
 581 rjMCMC provide similar PPD predictions, although the informative, Gaussian prior adopted by
 582 HMC often gives narrower posterior estimations with respect to rjMCMC. Differently the estimated
 583 PPDs for the V_p/V_s ratio of the first layer are substantially different, because mainly influenced by
 584 the prior model. Indeed, the HMC provides a Gaussian posterior centred on the prior mean value of
 585 4, while the uninformative prior used by the rjMCMC results in a posterior model very similar to
 586 the likelihood. In this case, note that two local maxima of the posterior model (Figure 11a rightmost
 587 plot) correspond to the two local minima of the misfit function (Figure 10d). The model parameter
 588 correlation can also be correctly inferred from the sampled models. Indeed Figure 11 clearly shows
 589 the independence of the V_s of the two layers, the weak positive correlation between the V_s and the
 590 thickness of the first layer, the independence of the V_s and the V_p/V_s ratio of the first layer, and the
 591 strong positive correlation between the V_s of the second layer and the thickness of the underlying
 592 layer.



593

594 *Figure 12: From left to right we represent: non-dimensional model resolution matrix, posterior*
 595 *model covariance matrix, and posterior model correlation matrix. a) and b) represent the results*
 596 *for the rjMCMC and HMC algorithms, respectively.*

597 For a further quantitative quantification of the results provided by the two sampling methods we use
 598 the ensemble of sampled models to compute some of the most popular sensitivity analysis tools
 599 such as the model resolution matrix and the posterior model covariance matrix. Since the inversion
 600 considers different parameter types (seismic velocity and thickness) the off-diagonal elements of the
 601 model resolution matrix and of the posterior covariance matrix become dimensionally dependent,
 602 thus complicating their visual inspection. For this reason, we resort to scale these matrices with
 603 respect to the prior and posterior standard deviation of the model parameters, thus deriving the non-
 604 dimensional model resolution matrix and posterior correlation matrix (see Sambridge, 1999). For
 605 example, the values corresponding to the i -th row and j -th column of the non-dimensional model
 606 resolution matrix is derived as:

$$\hat{R}_{i,j} = R_{i,j} \frac{\sigma_i}{\sigma_j} , \quad (27)$$

607

608 where σ_i and σ_j are the a-priori standard deviations of the i -th and j -th model parameter
609 respectively, whereas R and \hat{R} are the standard and non-dimensional model resolution matrices,
610 respectively. The posterior correlation matrix C can be computed as follows:

$$C_{i,j} = \frac{C^M_{i,j}}{\sqrt{C^M_{i,i} C^M_{j,j}}}, \quad (28)$$

611
612 where C^M is the posterior model covariance matrix. Figure 12 demonstrates that the rjMCMC and
613 HMC provide similar non-dimensional model resolution matrices, posterior covariance matrices
614 and posterior model correlation matrices especially for the parameters better illuminated by the data
615 (V_s and layer thickness). We observe an almost perfect resolution on the V_s of the two layers and a
616 slightly lower resolution for the layer thickness, while the V_p/V_s ratios are not resolved. Due to the
617 different parameter dimensionality the analysis of the posterior covariance matrix only reveals that
618 the V_s of the first layer is more accurately estimated than that of the second layer. The posterior
619 model correlation matrices show the almost null correlation among the shear wave velocities of the
620 two layers, the weak positive correlation between the V_s and the thickness of the first layer, the null
621 correlation between the V_p/V_s ratio and the V_s of the first layer, and the strong positive correlation
622 between the V_s of the second layer and the thickness of the first one.

623 For the sake of clarity, we point out that for the HMC algorithm the matrices shown in Figure 12
624 have been derived from the ensemble of models sampled after the burn-in period. Differently, not
625 all the models sampled by rjMCMC have the same dimensionality and for this reason the matrices
626 shown in Figure 12 have been computed from the sampled models with two layers, that corresponds
627 to the estimated MAP solution for the layer number.

628 Finally, this analysis confirms that although the rjMCMC and HMC algorithms use different model
629 parameterizations and prior models, they not only achieve comparable and congruent estimates of

630 V_s , and interface positions but also provide the same insight into the model parameter correlation
631 and resolution, at least for the parameters better constrained by the data. Indeed, for these
632 parameters the posterior sampling is mainly driven by the likelihood function instead of by the a-
633 priori constraints.

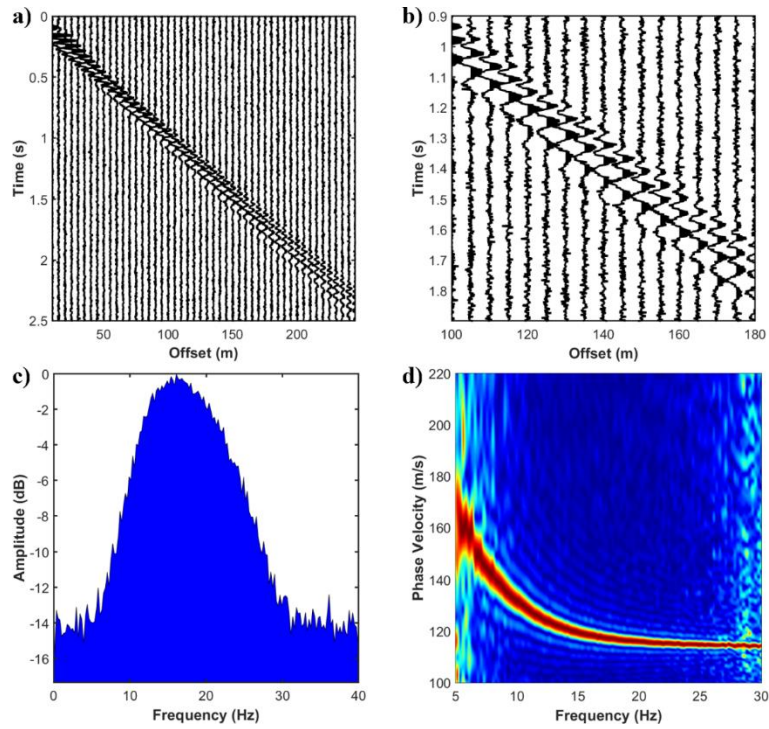
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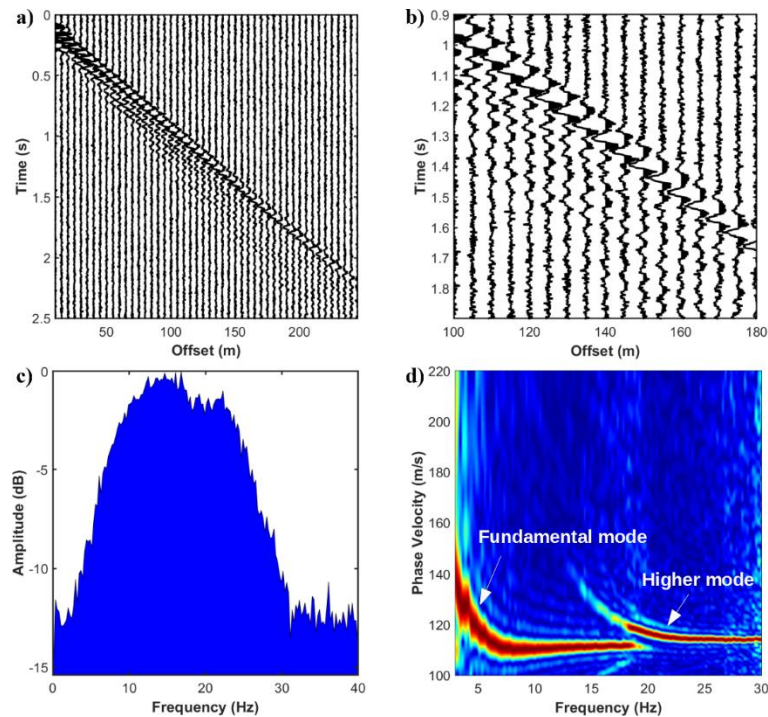
636 In the following experiments the observed dispersion curves have been picked on the frequency-
637 phase velocity spectra derived from synthetic shot gathers. For the sake of conciseness, we limit the
638 attention to two examples pertaining to two different 3-layer reference models. The first example
639 concerns a subsurface model with a steady increase of V_s with depth. The second example is more
640 challenging because it considers a model in which the V_s of the second layer is lower than that of
641 the first layer. This velocity inversion complicates the picking of the fundamental mode because
642 higher modes become more energetic than the fundamental mode at high frequencies. In both cases
643 we employ the reflectivity algorithm as the forward modelling, where the source signature is a 15-
644 Hz Ricker wavelet with a sampling interval of 2 ms. We simulate an off-end acquisition geometry
645 with a minimum offset of 10 m and 48 receivers equally spaced of 5 m, resulting in a maximum
646 offset of 245 m. Before the picking procedure we add to the synthetic seismic data, Gaussian
647 random noise with a standard deviation equal to the 100 % of the standard deviation of the noise-
648 free seismic dataset. As a result, differently from the previous inversion tests, the noise standard
649 deviation contaminating the observed dispersion pattern is not constant over the frequency range
650 but decreases as the frequency increases. For this reason, the diagonal entries of the data covariance
651 matrix are now related to the ambiguity affecting the picking of the dispersion curve. The prior
652 distributions are the same previously used in the semi-analytical tests.

653



654

655 *Figure 13: Data associated to the first 3-layer reference model used in the seismic tests. a) Noise*
 656 *contaminated shot gather. b) Close-up of a). c) Fourier amplitude spectra of a). d) Phase velocity*
 657 *spectra derived on a). In d) blue and red colors code low and high amplitude values, respectively.*



658

659 *Figure 14: Data associated to the second 3-layer reference model. a) Noise-contaminated shot*
 660 *gather. b) Close-up of a). c) Fourier amplitude spectra of a). d) Phase velocity spectra derived on*
 661 *a). In d) blue and red colors code low and high amplitude values, respectively.*

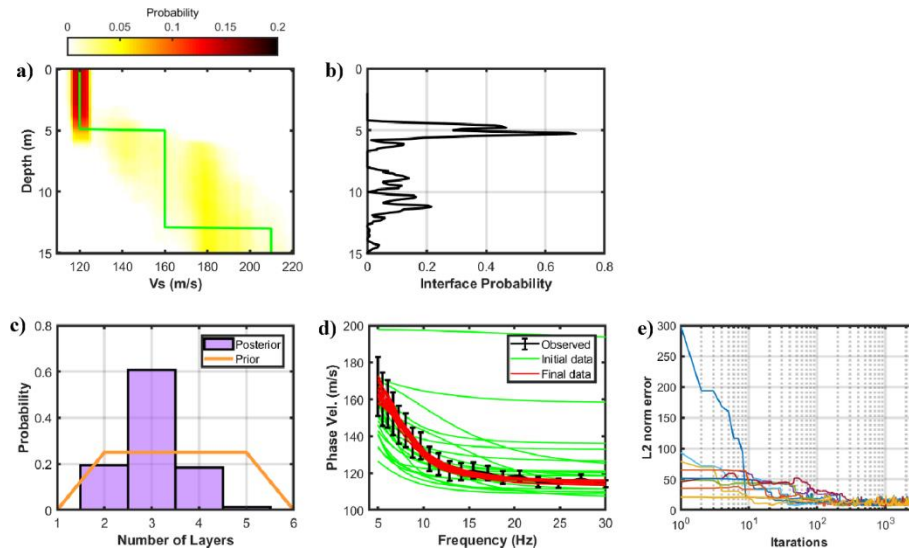
662 Figures 13 and 14 show the shot gathers, their Fourier amplitude spectra and the frequency-phase
 663 velocity spectra. In Figure 13 note that the steady V_s increase with depth makes it possible an
 664 accurate identification of the fundamental mode over the entire frequency range [5-30 Hz], but
 665 especially at high frequencies. Differently in Figure 14, the velocity inversion generates higher
 666 modes that become dominant for frequencies higher than 20 Hz. For this reason, in this case the
 667 rjMCMC and HMC inversions run in a narrower frequency band ranging from 5 to 20 Hz.

668

669 **Transdimensional Markov Chain Monte Carlo inversion**

670 Figure 15 shows the rjMCMC inversion results for the first 3-layer model. The V_s of the shallowest
 671 layer and the position of the first interface are well resolved. As expected, the uncertainties rapidly
 672 increase with depth. The overall loss of accuracy on the V_s and interface position with respect to the
 673 semi-analytical inversions is obviously produced by the loss of low frequencies (frequency lower

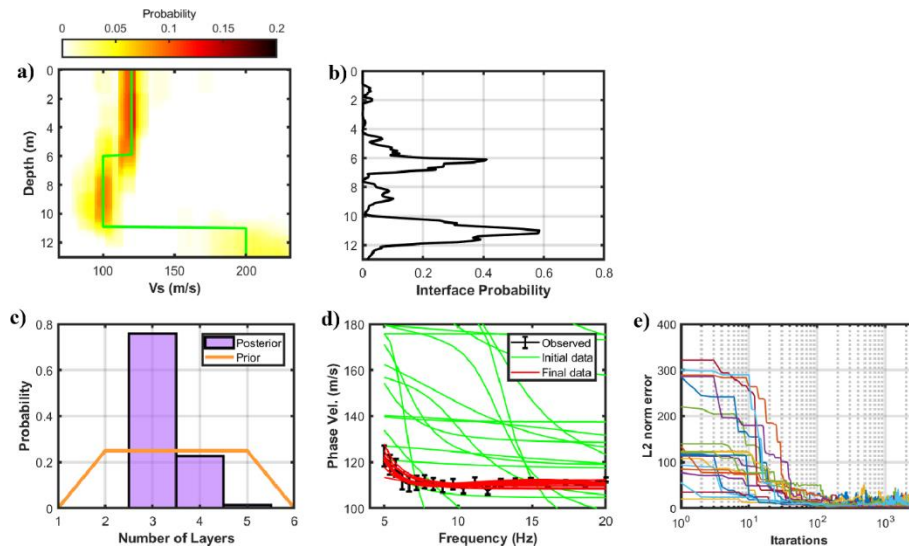
674 than 5 Hz) and by the higher uncertainties contaminating the picked dispersion curve at low
 675 frequencies. Figures 15c and 15d prove that the algorithm, even in this not favourable scenario,
 676 correctly identifies the correct number of layers and produces final dispersion curves that match the
 677 observed data. The evolution of the L2 norm error shows that the stationary regime is reached after
 678 200 iterations, approximately.



679
 680 *Figure 15: Seismic inversion results for the first 3-layer model. a) True Vs model (green line) and*
 681 *marginal PPD for Vs (colour scale). b) Marginal PPD for interface location. c) Prior and posterior*
 682 *probability density functions for the number of layers. d) Comparison between the picked*
 683 *dispersion curve (black line), the data generated on the starting models (green line), and the data*
 684 *generated on the last sampled models (red line). The black bars represent the noise standard*
 685 *deviation. e) Example of evolution of the L2 norm misfit for 10 chains out of 20.*

686
 687 Figure 16 displays the results for the second 3-layer model. We note that in spite of the more
 688 complex dispersion pattern generated by the shear-wave velocity inversion, the algorithm correctly
 689 estimates the velocity of the first two shallowest layers, and the velocity inversion occurring at 6 m
 690 depth, while the uncertainties significantly increase for the Vs of the deepest layer. The layer
 691 interfaces are still accurately located, and the number of layers is correctly determined. The

692 comparison of the picked and finally sampled dispersion curves confirms the convergence of the
693 algorithm to the stationary regime that is attained after 200 iterations, approximately.



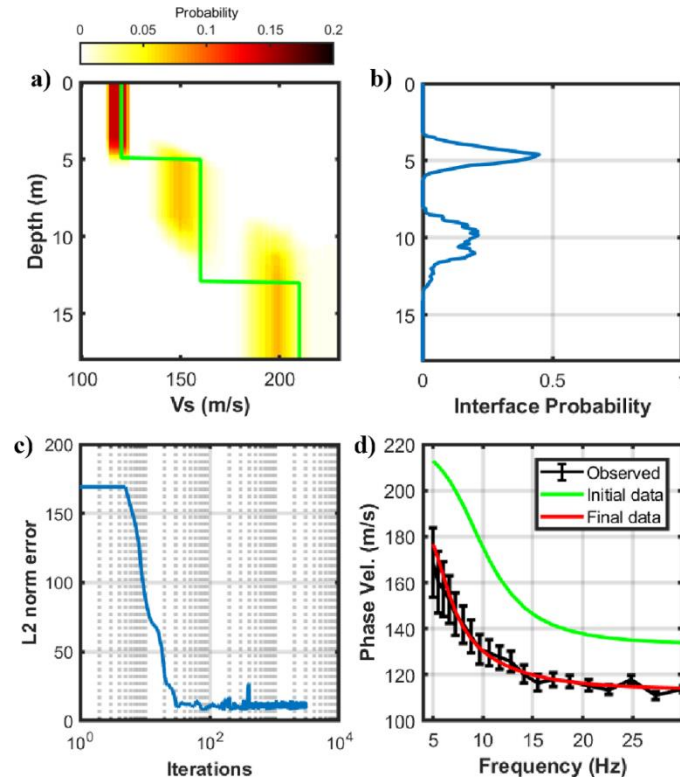
694

695 *Figure 16: As in Figure 15 but for the second 3-layer model.*

696

697 **Hamiltonian Monte Carlo inversion**

698 We now discuss the results obtained by the HMC algorithm on the data generated from the two
699 considered 3-layer models. Figure 17 represents the outcomes of the first example without the
700 velocity inversion. If we adopt a 3-layer parameterization we accurately predict the Vs of the
701 shallowest layer and the position of the first interface. The accuracy of the results rapidly decreases
702 as the depth increases, although the shear velocity of the second layer is still predicted with
703 acceptable accuracy, while the Vs of the third layer and the position of the second interface are
704 slightly underpredicted. This is obviously related to the lack of low frequencies in the data and to
705 the severe noise contamination at low frequencies that produces high uncertainties in the picking
706 phase (see Figure 14). However, the overall Vs trend and the positions of the layer interfaces are
707 still recovered with reasonable accuracy. The evolution of the data misfit shows the rapid
708 convergence of the HMC algorithm toward the stationary regime, that is attained in less than 30
709 iterations. Again, the final predicted data shows a good match with the picked dispersion curve.



710

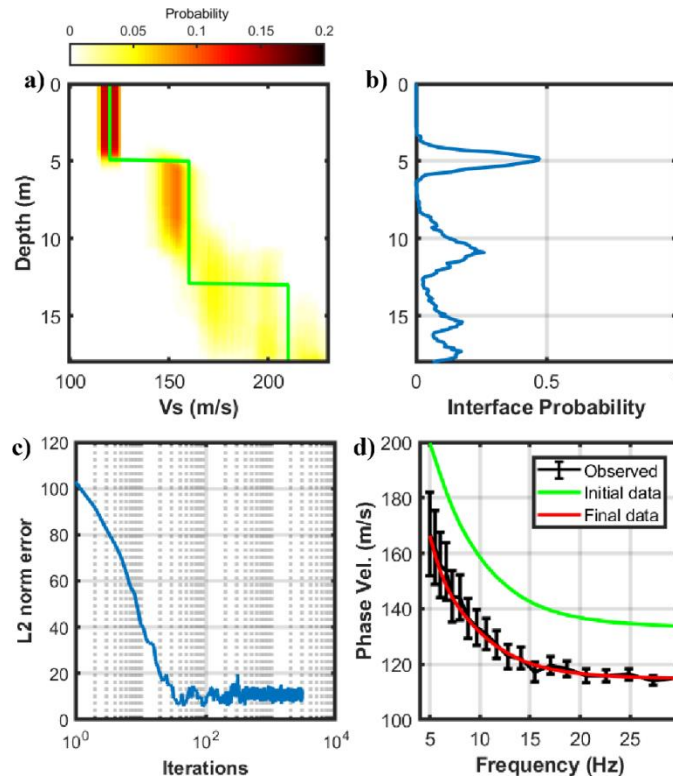
711 *Figure 17: Seismic inversion results for the first 3-layer model and considering a 3-layer*
 712 *parameterization. a) True Vs model (green line) and marginal Vs PPD (colour scale). b) Marginal*
 713 *PPD for interface location. c) Evolution of the L2 norm misfit. d) Comparison between the picked*
 714 *dispersion curve (black line), the data generated on the starting model (green line) and the data*
 715 *generated on the last sampled model (red line). The black bars represent the noise standard*
 716 *deviation.*

717 Similarly to the semi-analytical inversion tests, we now rerun the HMC inversion on the same data
 718 but employing a 4-layer parameterization. The results are shown in Figure 18. The Vs of the
 719 shallowest layer and the position of the first interface are accurately predicted and also the Vs of the
 720 second layer is well recovered, while high uncertainties affect the estimated positions of the deepest
 721 interfaces and the Vs of the deepest layer. Spurious interfaces are predicted by the algorithm below
 722 15 m depth. Similarly to the previous example, less than 30 iterations are enough to converge
 723 toward the stationary regime in which the predicted data closely match the dispersion curve
 724 extracted from the seismic shot gather.

725 Table 2 demonstrates that our inversion workflow correctly identifies the correct number of layers
 726 in this first seismic inversion and proves that the statistical tools we use ($p(\chi^2)$ and BIC) could
 727 constitute a valid help to select the optimal model parameterization.

	$p(\chi^2)$	BIC
2-layer parameterization	0.0092	38.93
3-layer parameterization	0.0342	25.20
4-layer parameterization	0.0105	31.30

728 *Table 2: $p(\chi^2)$ and BIC values for the first seismic inversion resulting from different model*
 729 *parameterizations.*



730
 731 *Figure 18: As in Figure 17 but for a 4-layer parameterization.*

732 If we compare Figure 17 with Figure 15, we note that HMC provides more accurate predictions
 733 (affected by lower uncertainties) with respect to the rjMCMC algorithm, especially for the V_s of the
 734 two deepest layers and the position of the deepest interface. Probably, conversely to the
 735 uninformative (i.e. uniform) priors used by the rjMCMC, the Gaussian prior employed by the HMC

736 method efficiently contributes to reduce the ensemble of possible solutions. In this case the a-priori
 737 constraints play a crucial role in determining the posterior model due to the ill-conditioning of the
 738 inverse problem that is mainly related to the narrow frequency band and to the noise contamination
 739 at low frequency.

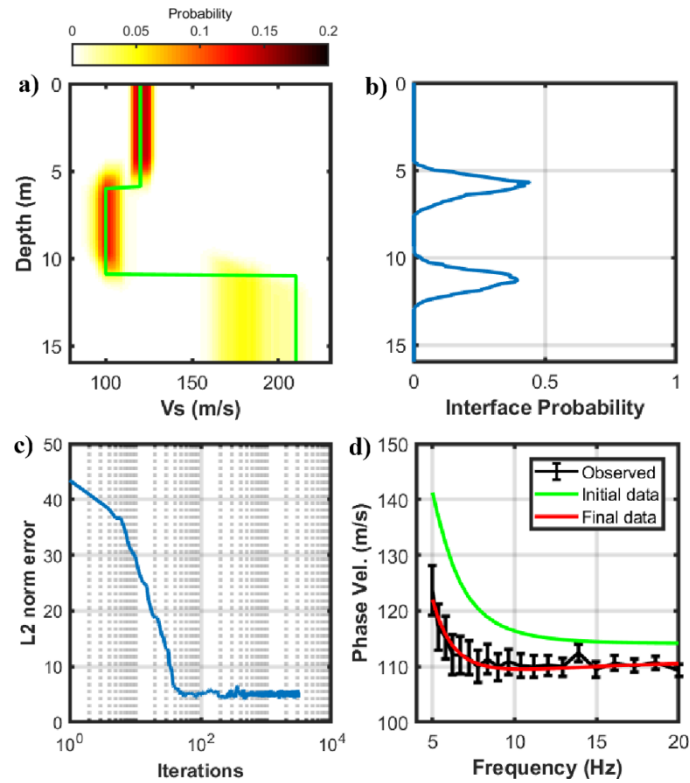
740 Figure 19 displays the results for the second model and for a 3-layer parameterization. The V_s of the
 741 two shallowest layers are well resolved, and despite the velocity inversion the algorithm accurately
 742 predicts a low velocity layer located between 6 m and 11 m depth. The uncertainties in the
 743 estimated V_s rapidly increase as the depth increases and the velocity of the deepest layer is not well
 744 recovered with a MAP solution that significantly underestimates the actual velocity value. Figure
 745 19b demonstrates that the HMC algorithm correctly identifies the position of the interfaces, whereas
 746 Figure 19c shows that in 30 iterations the algorithm converges toward the stationary regime, after
 747 which the observed data is fairly reproduced by the sampled models (Figure 19d).

	$p(\chi^2)$	BIC
2-layer parameterization	8×10^{-7}	48.92
3-layer parameterization	0.076	36.61
4-layer parameterization	0.021	41.42

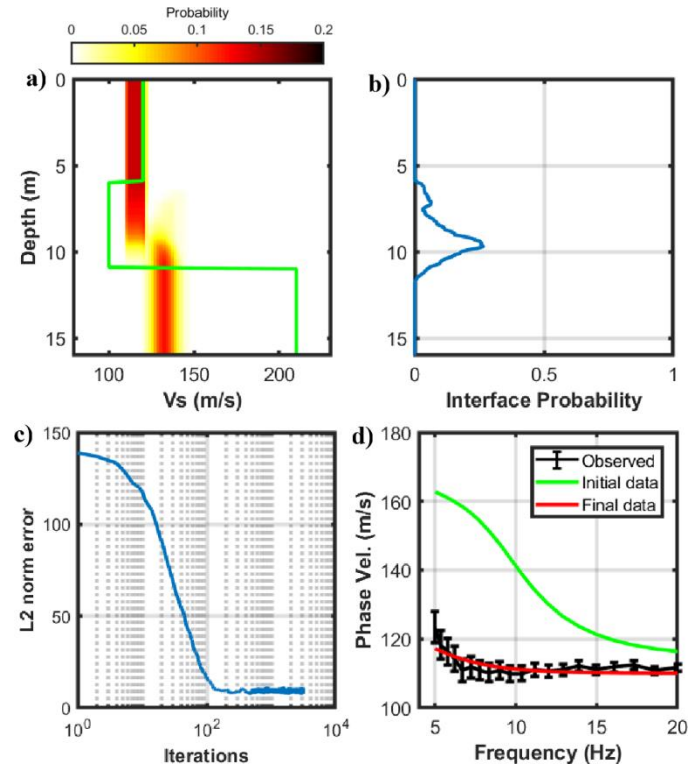
748 *Table 3: $p(\chi^2)$ and BIC values for the second synthetic test resulting from different model*
 749 *parameterizations.*

750 Figure 20 illustrates the results obtained on the same dataset but with a 2-layer parameterization.
 751 The velocity of the shallowest layer is still accurately retrieved, but the V_s of the deepest layers and
 752 the interface locations are not correctly estimated and show MAP solutions affected by significant
 753 biases. The evolution of the L2 norm error value demonstrates that in this case more iterations are
 754 needed to converge (for this reason we adopt only for this test a burn-in period of 150 samples),
 755 whereas the data generated on the final sampled model show some discrepancies with respect to the
 756 picked dispersion curve especially at low and high frequencies. Again, the $p(\chi^2)$ and BIC values

757 listed in Table 3 confirm that the proposed HMC workflow correctly estimates the actual number of
 758 layers.
 759



760
 761 *Figure 19: Seismic inversion results for the second 3-layer model and considering a 3-layer*
 762 *parameterization. a) True Vs model (green line) and marginal Vs PPD (colour scale). b) Marginal*
 763 *PPD for interface location. c) Evolution of the L2 norm misfit. d) Comparison between the picked*
 764 *dispersion curve (black line), the data generated on the starting model (green line) and the data*
 765 *generated on the last sampled model (red line). The black bars represent the noise standard*
 766 *deviation.*



767

768 *Figure 20: As in Figure 19 but for a 2-layer parameterization.*

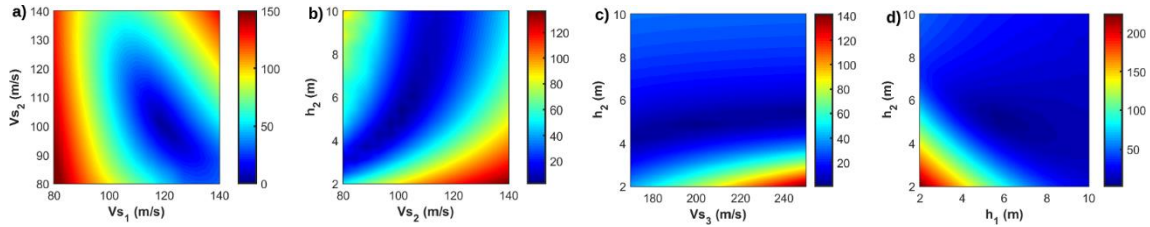
769 If we compare Figure 16 and Figure 19, we note that the rjMCMC and HMC algorithms provide
 770 congruent and comparable PPDs for the shear wave velocity and interface locations. This means
 771 that the posterior model for these parameters is mainly influenced by the observed data than by the
 772 a-priori constraints infused into the inversion framework. This aspect is deepened in the following
 773 section.

774

775 **Sensitivity analysis of the inversion results**

776 We now discuss the results of the sensitivity analysis for the second 3-layer model. Differently from
 777 the semi-analytical examples, we limit our discussion to the model parameters that can be inferred
 778 from the data (i.e. the V_p/V_s ratio is not discussed). Figure 21 shows some 2D projections of the L2
 779 norm misfit function. We observe that the shear wave velocities of the first two layers (V_{S1} and V_{S2})
 780 are the parameters that mostly influence the dispersion curve, followed by the layer thicknesses (h_1
 781 and h_2) and by the V_s of the deepest layers (V_{S3}). Between the considered parameters, the velocity
 782 of the third layer is the one that plays the minor role in determining the dispersion pattern.

783 From the orientation of the global minimum valleys we observe positive correlations between the
 784 V_s and the thickness of the second layer and between the V_s of the third layer and the thickness of
 785 the second one (Figures 21b and 21c), while the shear wave velocities and the thicknesses of the
 786 first two shallowest layers are negatively correlated (Figures 21a and 21d).

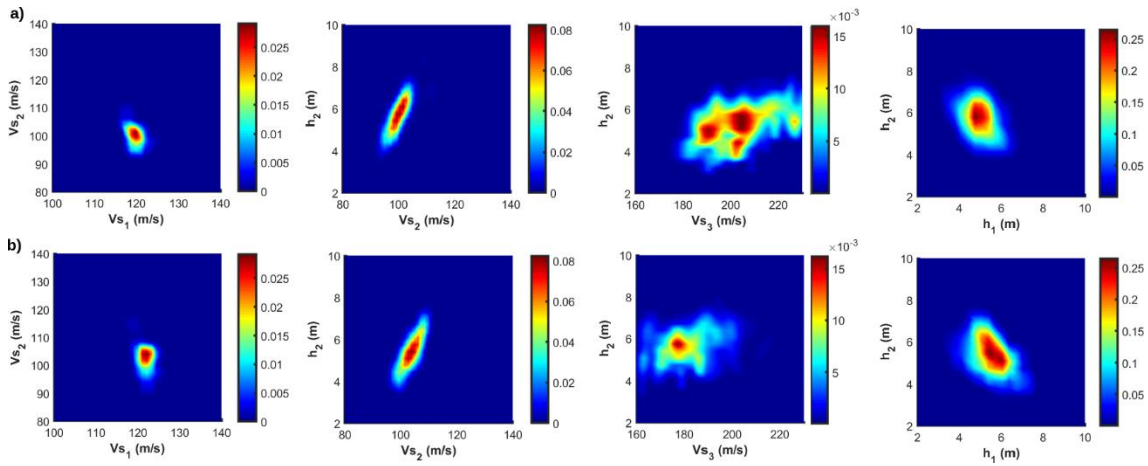


787
 788 *Figure 21: Some projections of the L2 norm misfit function onto different 2D sections. The*
 789 *subscripts 1 and 2 refer to the first (shallowest) and second (deepest) layer, respectively, whereas h*
 790 *indicates the layer thickness.*

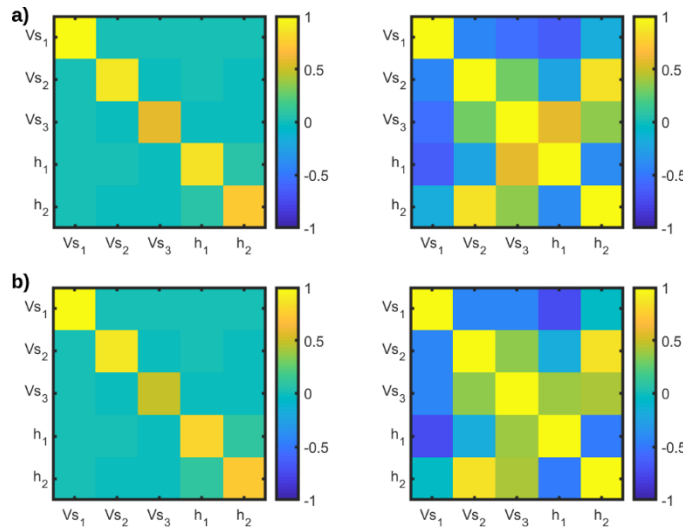
791
 792 Figure 22 shows 2D marginal PPDs estimated by the HMC and rjMCMC algorithms and projected
 793 onto the same sections previously considered in Figure 21. For both the rjMCMC and HMC
 794 inversions emerges that the velocities and the thicknesses of the first and second layer are well
 795 recovered, while higher uncertainties affect the estimated V_s of the deepest layer. We also observe a
 796 positive correlation between V_{S2} and h_2 , and between V_{S3} and h_2 , while negative correlations
 797 characterize the marginal PPDs projected onto the V_{S1} - V_{S2} and h_1 - h_2 planes.

798 The non-dimensional model resolution matrices (left hand side of Figure 23) confirm that the V_s of
 799 the two shallowest layers are the best resolved model parameters, while lower resolutions
 800 characterize the V_s of the third layer and the layer thicknesses. The mutual parameter correlations
 801 expressed by the 2D projections of the marginal PPDs can be also inferred from the inspection of
 802 the posterior correlation matrices (right hand side of Figure 23). The marginal PPDs displayed in
 803 Figure 22 and the non-dimensional model resolution matrices and posterior model correlation
 804 matrices of Figure 23 are in good agreement with the geometrical characteristics of the misfit

805 function (Figure 21). This proves that both the implemented algorithms accurately estimate the
 806 uncertainties affecting the model parameters and their mutual correlations.



807
 808 *Figure 22: 2D Marginal PPDs estimated by (a) the rjMCMC and by (b) the HMC algorithms.*



809
 810 *Figure 23: From left to right we represent: non-dimensional model resolution matrices, and*
 811 *posterior correlation matrices provided by the rjMCMC (a) and HMC(b) algorithms.*

812 DISCUSSION

814 The main advantage of the rjMCMC algorithm is the possibility to provide stable and reliable
 815 predictions with a minimum of a-priori constraints infused into the inversion kernel. Indeed, it does
 816 not require any statistical test to choose the adequate model parameterization and do not include any
 817 regularization operator to force the model to honour some external constraints. Instead, it
 818 automatically adjusts the underlying model parametrization to produce solutions with appropriate

819 level of complexity to fit the data to statistically meaningful levels. The downside of rjMCMC is
820 that a specific recipe for the problem at hand is usually needed to speed up to convergence of the
821 sampling toward the stationary regime. To mitigate this issue, in this work we employed a parallel
822 tempering strategy and a delayed rejection perturbation scheme: the former ensures optimal
823 exploitation and exploration capabilities, while the latter locally adapts the proposal to the shape of
824 the posterior. In particular, the convergence speed of the algorithm critically depends on the choice
825 of the statistical characteristics of such proposal distribution: a suboptimal proposal distribution
826 causes a persistent rejection of models, thus significantly reducing the exploitation of the algorithm.
827 Note that the updating of the different chains in the rjMCMC inversion can be done in parallel with
828 the different chains distributed across different processors. Such parallelizability and the 20
829 available CPUs make it possible the application of the parallel tempering and the delayed rejection
830 strategies with a limited extra computational cost. Indeed, in the parallel tempering the models
831 sampled at $T > 1$ must be disaggregated when computing the PPD and this fact implies that many
832 chains (each one involving a forward modelling run per iteration) are needed to converge to a stable
833 posterior. On the other hand, the delayed rejection implies the computation of an additional forward
834 solution in a single iteration when the first proposed model is rejected. For this reason, the
835 application of these two strategies must be carefully evaluated. In our inversion tests we found that
836 they are effective in speeding up the converge of the chains toward the stationary regime, in
837 promoting the exploration of the parameter space, and in increasing the acceptance rate of the
838 rjMCMC. Thanks to the limited number of unknown parameters (i.e. limited number of layers in the
839 inverted models) in our tests we found that an uncorrelated Gaussian proposal constitutes a good
840 compromise between acceptance rate and convergence of the sampling. If needed more
841 sophisticated strategies based on SVD decomposition and orthogonal projections could be used to
842 include a covariance model into the model perturbation scheme (Dosso et al. 2014). Another limit
843 of the rjMCMC is that the successively sampled models are usually highly correlated. For this
844 reason, to avoid biased PPD estimations (MacKay, 2003) not all the models sampled after the burn-

845 in phase have been used to numerically compute the PPD, but several iterations of the algorithm are
846 allowed to elapse in between successive samples. In all the previous examples only 1 model every 5
847 is considered in the computation of the posterior model.

848 The main advantage of HMC is that it guarantees an efficient sampling with a rapid convergence
849 toward the stationary regime. This ability rests on the exploitation of the derivative information that
850 other Monte Carlo methods, such the well-known random walk Metropolis (that is the basis of the
851 rjMCMC algorithm), do not consider. In particular, the HMC provides highly independent samples
852 and for this reason a lower number of iterations are needed for a reliable estimation of the posterior
853 model with respect to MCMC algorithms. Indeed, differently from the rjMCMC inversion, the high
854 independence of successively sampled models granted by the HMC framework, allows us to exploit
855 all the models collected after the burn-in phase in computing the PPD. In this work, the HMC
856 algorithm has been implemented for a Gaussian prior model, but another outstanding benefit of
857 HMC is the possibility to consider either parametric or non-parametric priors. However, making the
858 HMC sampling still efficient in highly multimodal distributions is an active research field nowadays
859 (Girolami and Calderhead, 2011; Nishimura and Dunson, 2016; Graham and Storkey, 2017).
860 Differently from rjMCMC, one limit of HMC is that the model parameterization is an input to the
861 inversion. In this study we propose to solve this issue by performing different inversions with
862 different parameterizations and then using the sampled models and standard statistical tools to
863 determine the most probable number of layers. Another limit is that the HMC algorithm can be
864 applied to problems where the derivative of the posterior model can be rapidly computed and where
865 this derivative is continuous. For example, such derivative becomes discontinuous if multiple
866 dispersion modes are considered. For this reason, in this work we have limited the attention the
867 fundamental mode only. Currently, we are investigating this issue in order to possibly extend the
868 HMC to the inversion of multiple modes. For the sake of consistency, in this work also the
869 rjMCMC inversion has been limited to the fundamental model. However, the rjMCMC algorithm
870 does not exploit any derivative information and for this reason the current implementation can be

871 directly used for inverting higher modes. Another critical element of HMC is the choice of a
872 suitable mass matrix. In the examples discussed here we found that setting this matrix equal to the
873 inverse of the locally approximated posterior covariance guaranteed stable posterior estimations.
874 Other choices of the mass matrix, for examples a diagonal matrix (e.g. a scalar multiple of the
875 identity matrix) resulted in very slow convergence and unreliable posterior assessments (see
876 Fichtner et al. 2019). Another critical aspect of HMC is the choice of the length of the trajectory in
877 the phase space. We overcome this issue by drawing the L parameter (which controls the number of
878 time integration steps and the trajectory length) from a previously defined uniform distribution in
879 each iteration. However, the trajectory length can be adaptively set, for example by adopting the so
880 called no U-turns sampling method (Gelman et al. 2013; Hoffman and Gelman et al. 2014) that
881 terminates the integration when the trajectory begins to return towards its starting point. In any case,
882 in the examples discussed here we found that 3–6 time steps per trajectory are usually optimal.
883 Using fewer time steps, leads to very high acceptance rates but at the expense of a limited model
884 space exploration and slow convergence. Otherwise, a larger number of time steps resulted in low
885 acceptance rates and in a decreased accuracy of the numerical integration.

886 To get the benefits of both the rjMCMC and HMC algorithms, we are also investigating the
887 possibility to hybridize these two methods and to implement a transdimensional Hamiltonian
888 algorithm for surface wave inversions, that is an inversion approach that infers the most appropriate
889 model parameterizations from the data and that is also characterized by an extremely efficient
890 sampling of the parameter space. An example of this approach applied to reflection seismic data can
891 be found in Sen and Biswas (2017).

892 We ran all our experiments on two deca-core intel E5-2630 @2.2 GHz (128 Gb RAM) and
893 employing parallel Matlab inversion codes. Table 4 lists some more quantitative details about the
894 HMC and rjMCMC seismic inversions on the second subsurface model (the V_s model with the
895 velocity inversion). The total computational time was 11 and 19 minutes for rjMCMC and HMC,
896 respectively. We point out that in our current implementations the rjMCMC algorithm is more

897 scalable than HMC, although the HMC requires a lower number of sampled models to provide a
898 stable posterior. In particular, note that the total number of forward modelling runs in the HMC is
899 higher than that associated to the rjMCMC. Indeed, for a single chain and for a single iteration, the
900 rjMCMC algorithm requires just one forward modelling run, while in the HMC several forward
901 evaluations are needed to compute the Jacobian matrix associated to each considered model. This
902 number of forward evaluations depends on the number of unknown parameters, on the length of the
903 trajectory in the phase space (value of the L parameter), and on the adopted FD scheme. For a Q -
904 dimensional model space, the number of forward evaluations per iteration is given by $(Q \times 2 + 1) \times L$ if
905 a central FD scheme is adopted. This number reduces to $(Q + 1) \times L$ if a forward FD is employed. For
906 this reason, a computationally efficient derivation of the Jacobian matrix would be crucial to
907 drastically reduce the computational time of each iteration. Although we employ a computationally
908 intensive numerical approach to derive the Jacobian, we reduce the computational effort of the
909 HMC by computing each column of this matrix in parallel. Indeed, the columns of the Jacobian are
910 independent to each other and their computation can be distributed across different processors. This
911 fact, together with the limited computational cost of the forward solver, did not hamper the
912 application of this method to the considered examples. Possible strategies to drastically reduce the
913 number of forward evaluations could be the adoption of a less computationally expensive FD
914 scheme (e.g. forward finite difference) and/or the adoption of a more parsimonious parameterization
915 of the model space. From the one hand, additional experiments we carried out (not shown here for
916 brevity) demonstrated that the application of the forward difference approach does not sensibly
917 affect the sampling efficiency of the method. On the other hand, since the number of columns on
918 the Jacobian is equal to the number of unknowns, we can reduce the number of forward modelling
919 runs by reducing the number of model parameters, for example by inverting for the V_s and layer
920 thickness and not for the V_p/V_s ratio that is known to be not constrained by the data. To put this in
921 perspective, if we consider a forward FD scheme and we remove the V_p/V_s ratio from the model
922 vector, the number of forward evaluations on the seismic inversion on the second subsurface model

923 is reduced from 229500 to 81000 (see Table 4). Note that this value is lower than the number of
 924 forward evaluations ran by the rjMCMC inversion on the same example (103200). Machine
 925 learning approaches could also be useful to reduce the computational cost of MCMC inversion
 926 (Hansen and Cordua 2017). In our specific case, machine learning algorithm can be trained to predict
 927 a local approximation of the Jacobian matrix around a considered model.
 928

	rjMCMC	HMC
Average computational time for a single forward modelling evaluation	≈ 0.063 s	
Number of iterations (per chain)	3000	3000
Number of chains	20	1
Total number of forward modelling evaluations	103200 (this number takes into account also the forward modellings related to the delayed rejection)	229500 (this number takes into account the forward modelling runs needed to compute the Jacobian with the central finite difference approach)
Total number of sampled models	60000	3000
Number of models used to compute the posterior	5600	2970
Average computational time of a single iteration	≈ 0.21 s (considering all the 20 chains	≈ 0.38 s (This number varies for

	running in parallel)	different iterations because is dependent on the drawn L value)
Average computational time for the Jacobian computation in a single iteration	<i>Not applicable</i>	≈ 0.32 s (This number varies for different iterations because is dependent on the drawn L value)

929 Table 4: Some details of the rjMCMC and HMC inversions for the last seismic inversion.

930

931

CONCLUSIONS

932 We implemented a Hamiltonian Monte Carlo (HMC) algorithm and a reversible jump Marko chain
933 Monte Carlo (rjMCMC) algorithm for Rayleigh wave dispersion curve inversion. This inverse
934 problem is highly non-linear and highly ill-conditioned. For this reason, the ultimate goal of this
935 study was to compare two inversion approaches that guarantee reliable uncertainty quantifications,
936 that estimate the most probable model parameterization (i.e. number of layers in the subsurface),
937 and that yield stable predictions not affected by the choice of the starting model. We limited the
938 attention to synthetic tests to maintain the discussion at a didactic level and to draw essential
939 conclusions about the main benefits and drawbacks of the two implemented approaches in the
940 context of dispersion curve inversion. Our experiments demonstrated that HMC and rjMCMC are
941 very promising approaches for dispersion curve inversion as they provide reliable assessment of the
942 posterior uncertainties also for this highly non-linear and severely ill-conditioned inverse problem.
943 In particular, the HMC and rjMCMC algorithms yielded uncertainty quantifications and model
944 predictions in accordance with the expected model parameter illuminations and model parameter
945 correlations. For example, the uncertainty increases passing from V_s to V_p/V_s ratio and increases as

946 the depth of investigation increases or when the noise in the observed data increases.
947 Notwithstanding the two algorithms use different model parameterizations and prior assumptions,
948 they estimate congruent and comparable posterior uncertainties and model parameter correlations
949 especially for the parameters better illuminated by the data, that is for those parameters for which
950 the PPD is mainly influenced by the likelihood function instead of by the a-priori constraints. The
951 transdimensional inversion approach and the combined used of HMC and statistical tools (such as
952 BIC and $p(\chi^2)$) demonstrated to be two valid strategies to determine the model parameterization (i.e.
953 number of layers), at least for subsurface models with a limited number of layers (less than 4 if only
954 the fundamental model is considered).
955 For both approaches an efficient parallel code is essential to decrease the computational time of the
956 inversion procedure related to the many forward modelling runs needed by the HMC to numerically
957 compute the Jacobian matrix, and by the rjMCMC to make evolve the different chains in a single
958 iteration.

959

960

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962 criticism and comments that helped to significantly improve the quality of the original manuscript.

963

964 **Conflict of interest**

965 The authors declare no conflict of interest

966

967 **Data Availability Statement**

968 Data available on request from the authors

969

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