FREE BOUNDARY REGULARITY FOR A MULTIPHASE SHAPE OPTIMIZATION PROBLEM

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ABSTRACT. In this paper we prove a $C^{1,\alpha}$ regularity result in dimension two for almost-minimizers of the constrained one-phase Alt-Caffarelli and the two-phase Alt-Caffarelli-Friedman functionals for an energy with variable coefficients. As a consequence, we deduce the complete regularity of solutions of a multiphase shape optimization problem for the first eigenvalue of the Dirichlet-Laplacian up to the fixed boundary. One of the main ingredient is a new application of the epiperimetric-inequality of [\[18\]](#page-21-0) up to the boundary. While the framework that leads to this application is valid in every dimension, the epiperimetric inequality is known only in dimension two, thus the restriction on the dimension.

1. INTRODUCTION

In [\[12\]](#page-20-0) David and Toro studied properties of the free boundaries for almost-minimizers of the one-phase Alt-Caffarelli and the two-phase Alt-Caffarelli-Friedman functionals considered in [\[1\]](#page-20-1) and [\[2\]](#page-20-2) respectively. They proved that, in any dimension, almost-minimizers are non-degenerate and Lipschitz continuous (see also $\boxed{17}$). More recently, David-Engelstein-Toro proved in $\boxed{11}$ that, under suitable assumption, the free boundaries of almost-minimizers are uniformly rectifiable for both functionals, and almost everywhere given as the graph of a $C^{1,\alpha}$ function for the one-phase functional.

In this paper we extend these regularity results in dimension two, proving general $C^{1,\alpha}$ regularity results for the free boundary associated to:

(OP) Almost-minimizers of the one-phase Alt-Caffarelli functional for an operator with variable coefficients, which may also satisfy a further geometric inclusion constraint (Theorem $\overline{1.1}$ and Corollary $[1.3]$;

(TP) Almost-minimizers of the two-phase Alt-Caffarelli-Friedman functional for an operator with variable coefficients (Theorem $\overline{1.5}$).

As pointed out by David-Toro, the difficulty of dealing with almost-minimizers is that they do not satisfy an equation. To overcome this, the approach we follow in this paper is different from the one given in $\boxed{11}$ and relies on an epiperimetric inequality (see $\boxed{18}$) and on a Weiss' almost-monotonicity formula, both of which are variational techniques. We stress out that the only obstruction to a generalization of this proof to any dimension comes from the fact that epiperimetric inequality is only known in dimension two.

The second purpose of this paper is a regularity result for solutions of a multiphase shape optimization problem for the first eigenvalue of the Dirichlet-Laplacian. In **3** Bogosel and Velichkov proved properties concerning the optimal shapes, such that the lack of triple points, and the Lipschitz continuity of the corresponding eigenfunctions. In this paper we prove, using the above results for almost-minimizers of the one-phase and the two-phase functionals, that if $(\Omega_1, \ldots, \Omega_n)$ is an optimal shape, then the entire boundary $\partial\Omega_i$, $i=1,\ldots,n$, is $C^{1,\alpha}$ regular (Theorem [1.8\)](#page-2-1). This line of study has become increasingly important in recent years, where regularity results for solutions of free-boundary problems, and in particular almost-minimizers, have been applied to study the regularity of shape optimization problems involving eigenvalues of the Dirichlet-Laplacian (see for instance $[16, 14, 15, 9]$ $[16, 14, 15, 9]$ $[16, 14, 15, 9]$ $[16, 14, 15, 9]$).

1.1. Regularity for almost-minimizers. Throughout this paper we will use the following notations. We fix a matrix valued function $A = (a_{ij})_{ij} : B_2 \to Sym_2^+$, where Sym_k^+ denotes the family of the real positive symmetric $k \times k$ matrices, for which there are constants δ , C_A , $M_A > 0$ such that

$$
|a_{ij}(x) - a_{ij}(y)| \le C_A |x - y|^\delta, \quad \text{for every} \quad i, j \quad \text{and} \quad x, y \in B_2;
$$

$$
M_A^{-1} |\xi|^2 \le \xi \cdot A(x)\xi = \sum_{i,j=1}^2 \xi_i \xi_j a_{ij}(x) \le M_A |\xi|^2, \quad \text{for every} \quad x \in B_2 \quad \text{and} \quad \xi \in \mathbb{R}^2.
$$

We fix Q_{OP} , Q_{TP}^+ and Q_{TP}^- to be three different δ-Hölder continuous functions on B_2 , for which there is a constant $C_Q > 0$ such that $C_Q^{-1} \leq Q_{\text{OP}}, Q_{\text{TP}}^+$, $Q_{\text{TP}}^- \leq C_Q$ on B_2 . Finally, for every function $u : \mathbb{R}^2 \to \mathbb{R}$, we will use the following standard notations

 $u_{\pm}(x) := \max\{\pm u(x), 0\}$, $\Omega_u := \{u \neq 0\}$, $\Omega_u^+ := \{u > 0\}$ and $\Omega_u^- := \{u < 0\}$.

We are now in position to state our main free boundary regularity results. The one-phase free boundaries. For every $u \in H^1(B_2)$, $x_0 \in B_1$ and $r \in (0,1)$, we define the one-phase functional

$$
J_{\text{OP}}(u, x_0, r) = \int_{B_r(x_0)} \left(\sum_{i,j} a_{ij}(x) \frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_j} + Q_{\text{OP}}(x) \mathbb{1}_{\{u > 0\}} \right) dx.
$$

Here and after $B_r(x)$ denotes the ball with center $x \in \mathbb{R}^2$ and radius $r > 0$ and we will write $B_r := B_r(0)$. Let $\mathcal{A}^+(B_r)$ be the admissible set

$$
\mathcal{A}^+(B_r) = \{ u \in H^1(B_r) : u \ge 0 \text{ in } B_r, u = 0 \text{ on } B_r \setminus B_r^+ \},
$$

where H stands for the upper half-plane $H := \{(x, y) \in \mathbb{R}^2 : y > 0\}$ and $B_r^+ := B_r \cap H$. We say that the nonnegative function $u : B_2 \to \mathbb{R}$ is a *almost-minimizer of the one-phase functional* J_{op} in the upper half-disk B_2^+ , if $u \in A^+(B_2)$ and there are constants $r_1 > 0$, $C_1 > 0$ and $\delta_1 > 0$ such that, for every $x_0 \in B_1 \cap \partial \Omega_u$ and $r \in (0, r_1)$, we have

(1.1)
$$
J_{\text{OP}}(u, x_0, r) \leq (1 + C_1 r^{\delta_1}) J_{\text{OP}}(v, x_0, r) + C_1 r^{2+\delta_1}
$$

for every $v \in \mathcal{A}^+(B_2)$ such that $u = v$ on $B_2 \setminus B_r(x_0)$.

We have the following result for the almost-minimizers of the one-phase Alt-Caffarelli functional J_{OP} constrained in the upper half-disk B_2^+ .

Theorem 1.1 (Regularity of the constrained one-phase free boundaries). Let $B_2 \subset \mathbb{R}^2$ and $u : B_2 \to \mathbb{R}$ be a non-negative and Lipschitz continuous function. If u is a almost-minimizer of the functional J_{OP} in $\mathcal{A}^+(B_2)$, then the free boundary $B_1 \cap \partial \Omega_u$ is locally the graph of a $C^{1,\alpha}$ function.

Remark 1.2. The Hölder continuity of the (exterior) normal vector n_{Ω} is the best regularity result that one can expect. Indeed, recently Chang-Lara and Savin [\[10\]](#page-20-9) showed that even for minimizers the regularity of the constrained free boundaries cannot exceed $C^{1,1/2}$. Moreover, we notice that the result analogous to Theorem $\boxed{1.1}$ was proved in any dimension in $\boxed{10}$, by a viscosity approach, but only for minimizers of the functional J_{op} .

Analogously, we say that the nonnegative function $u : B_2 \to \mathbb{R}$ is a *almost-minimizer of the* one-phase functional J_{OP} in B_2 , if $u \in H^1(B_2)$ and there are constants $r_1 > 0$, $C_1 > 0$ and $\delta_1 > 0$ such that, for every $x_0 \in B_1 \cap \partial \Omega_u$ and $r \in (0, r_1)$, we have

(1.2)
$$
J_{\text{op}}(u, x_0, r) \le (1 + C_1 r^{\delta_1}) J_{\text{op}}(v, x_0, r) + C_1 r^{2 + \delta_1} \text{for every } v \in H^1(B_2) \text{ such that } u = v \text{ on } B_2 \setminus B_r(x_0).
$$

The regularity of the unconstrained one-phase free boundary $\partial \Omega_u$ follows directly by Theorem [1.1.](#page-1-0) For the sake of completeness, we give the precise statement in Corollary [1.3](#page-1-1) below.

Corollary 1.3 (Regularity of the unconstrained one-phase free boundaries). Let $B_2 \subset \mathbb{R}^2$ and $u : B_2 \to \mathbb{R}$ be a non-negative and Lipschitz continuous function. If u is a almost-minimizer of the functional J_{op} in B_2 , then the free boundary $B_1 \cap \partial \Omega_u$ is locally the graph of a $C^{1,\alpha}$ function.

Remark 1.4. We notice that the regularity of the free boundaries of the one-phase (unconstrained) almost-minimizers is already proved by David, Engelstein and Toro in [\[11\]](#page-20-3) in every dimension and by a different approach.

The two-phase free boundaries. For every $u \in H^1(B_2)$, $x_0 \in B_1$ and $r \in (0,1)$, we define the two-phase functional

$$
J_{\text{TP}}(u, x_0, r) = \int_{B_r(x_0)} \Big(\sum_{i,j} a_{ij}(x) \frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_j} + Q_{\text{TP}}^+(x) 1\!\!1_{\{u > 0\}} + Q_{\text{TP}}^-(x) 1\!\!1_{\{u < 0\}} \Big) \, dx.
$$

We say that the function $u \in H^1(B_2)$ is a almost-minimizer of the two-phase functional J_{TP} in B₂, if there are constants $r_2 > 0$, $C_2 > 0$ and $\delta_2 > 0$ such that, for every $x_0 \in B_1 \cap \partial \Omega_u$ and $r \in (0, r_2)$, we have

(1.3)
$$
J_{\text{TP}}(u, x_0, r) \le (1 + C_2 r^{\delta_2}) J_{\text{TP}}(v, x_0, r) + C_2 r^{2 + \delta_2}
$$

for every $v \in H^1(B_2)$ such that $u = v$ on $B_2 \setminus B_r(x_0)$.

Then, we have the following result:

Theorem 1.5 (Regularity of the two-phase free boundaries). Let $B_2 \subset \mathbb{R}^2$ and let $u : B_2 \to \mathbb{R}$ be Lipschitz continuous such that the functions u_{\pm} are solutions of the PDEs

$$
-\text{div}\,(A\nabla u_{\pm}) = f_{\pm} \quad in \quad \Omega_u^{\pm} \cap B_2,
$$

where $f_{\pm} : \overline{\Omega}_{u}^{\pm} \to \mathbb{R}$ are Hölder continuous functions and A is fixed as above. If u is a almostminimizer of the two-phase functional J_{TP} in B_2 , then the free boundaries $B_1 \cap \partial \Omega_u^+$ and $B_1 \cap \partial \Omega_u^$ are locally graphs of $C^{1,\alpha}$ functions.

Remark 1.6. In particular, we improve from C^1 to $C^{1,\alpha}$ the regularity of the free boundaries proved by Alt, Caffarelli and Friedman [\[2\]](#page-20-2) for minimizers in the case $A = Id$.

Remark 1.7 (Remark on the Lipschitz continuity). In Theorem [1.3](#page-1-1) and Theorem [1.5](#page-2-0) we assume that the function u is Lipschitz continuous. In the case of the Laplacian, David and Toro $[12]$ proved that the Lipschitz continuity is a consequence of the the almost-minimality condition. It is of course natural to expect that the same will hold when the operator involved has variable coefficients. We will not address this question in the present paper since our main motivation comes from the application to shape optimization problems as [\(1.4\)](#page-2-2), for which the Lipschitz continuity is often already known. Actually, in the case of [\(1.4\)](#page-2-2), the Lipschitz continuity of the eigenfunctions is used to deduce the almost-minimality (see Section [7\)](#page-15-0).

1.2. Multiphase shape optimization problem for the first eigenvalue. As a consequence of Theorem [1.1](#page-1-0) and Theorem [1.5,](#page-2-0) we get the complete regularity of the following multiphase shape optimization problem

(1.4)
$$
\min \left\{ \sum_{i=1}^{n} \left(\lambda_1(\Omega_i) + q_i |\Omega_i| \right) : \Omega_1, \dots, \Omega_n \text{ are disjoint open subsets of } \mathcal{D} \right\},
$$

where, we will use the following notations:

- $1 \leq n \in \mathbb{N}$, and $0 < q_i \in \mathbb{R}$, for every $i = 1, \ldots, n$;
- $\mathcal{D} \subset \mathbb{R}^2$ is a bounded open planar set with C^2 regular boundary;
- $|\Omega|$ denotes the Lebesgue measure of Ω ;
- $\lambda_1(\Omega)$ is the first eigenvalue of the Dirichlet Laplacian on Ω .

Theorem 1.8. Let $(\Omega_1, \ldots, \Omega_n)$ be a solution of [\(1.4\)](#page-2-2). Then, all the sets Ω_i , $i = 1, \ldots, n$, are bounded open sets with $C^{1,\alpha}$ regular boundary.

We notice that, in the above theorem, the *entire* boundary $\partial\Omega_i$, $i=1,\ldots,n$, is regular. In particular, this holds at the contact points with the other phases Ω_j , $j \neq i$, and also with the boundary of the box $\partial \mathcal{D}$.

Remark 1.9 (On the regularity of the box). We state the regularity result of Theorem [1.8](#page-2-1) (and of the corollary below) with a box D of class C^2 . However, it is possible to weaken this assumption, with exactly the same proof, by assuming that D is a bounded open set of class $C^{1,\alpha}$ such that the solution w to the PDE

$$
-\Delta w = 1 \text{ in } D, \qquad w \in H_0^1(D),
$$

is a Lipschitz continuous function in \mathbb{R}^d . It is for instance the case if D is $C^{1,1}$ regular.

Moreover, note that, in the special case $n = 1$, [\(1.4\)](#page-2-2) reduces to the classical shape optimization problem

(1.5)
$$
\min \{ \lambda_1(\Omega) + \Lambda |\Omega| : \Omega \text{ open, } \Omega \subset \mathcal{D} \}.
$$

The existence in the class of open sets and the regularity of the free boundary (precisely, of the part contained in the open set \mathcal{D}) was proved by Briançon and Lamboley in [\[4\]](#page-20-10). As a direct corollary of our Theorem [1.8,](#page-2-1) we obtain that the entire boundary is $C^{1,\alpha}$ regular.

Corollary 1.10 (Regularity of the optimal sets for the first eigenvalue). Let $\mathcal{D} \subset \mathbb{R}^2$ be a bounded open set of class C^2 and let $\Lambda > 0$. Then, there is $\alpha \in (0,1)$ such that every solution $\Omega \subset \mathcal{D}$ of (1.5) is $C^{1,\alpha}$ regular.

1.3. Organization of the paper. In Section [2](#page-3-1) we recall the definitions of the Weiss' boundary adjusted energies and the statements of the epiperimetric inequalities. Moreover, we show how to use the algebraic properties of these quantities to deduce the rate of convergence of the blow-up sequences and the uniqueness of the blow-up limits. In Section [3](#page-6-0) we prove a technical lemma that reduces the one-phase and two-phase problems to the case of the Laplacian, which allows us to apply the results of Section [2.](#page-3-1) Section [4](#page-8-0) is dedicated to the classification of the blow-up limits for the one-phase and the two-phase problems. In Section [5](#page-11-0) and Section [6](#page-13-0) we prove Theorem [1.1](#page-1-0) and Theorem [1.5,](#page-2-0) respectively. In Section [7](#page-15-0) we prove that the (eigenfunctions associated to the) solutions of the multiphase problem [\(1.4\)](#page-2-2) are locally almost-minimizers of the one-phase or the two-phase problems, and we prove Theorem [1.8.](#page-2-1)

2. Boundary adjusted energy and epiperimetric inequality

All the arguments in this section hold in every dimension $d > 2$, except the epiperimetric inequalities Theorem [2.2](#page-4-0) and Theorem [2.3,](#page-5-0) which are known only in dimension two.

2.1. One-homogeneous rescaling and excess. Let $d \geq 2$ and $u \in H_{loc}^1(\mathbb{R}^d)$. For $r > 0$ and $x_0 \in \mathbb{R}^d$, we define the one-homogeneous rescaling of u as

(2.1)
$$
u_{x_0,r}(x) := \frac{u(x_0 + rx)}{r} \quad \text{for every} \quad x \in \mathbb{R}^d.
$$

Then, $u_{x_0,r} \in H^1_{loc}(\mathbb{R}^d)$ and for almost every $r > 0$, $E(u_{x_0,r})$ is well defined, where we set

(2.2)
$$
E(v) := \int_{\partial B_1} |x \cdot \nabla v - v|^2 d\mathcal{H}^{d-1},
$$

where $x \in \partial B_1$ is the exterior normal derivative to ∂B_1 at the point $x \in \mathbb{R}^d$ and \mathcal{H}^{d-1} stands for the $(d-1)$ -dimensional Hausdorff measure. The excess function $e(r) = E(u_{x_0,r})$ controls the asymptotic behavior, as $r \to 0^+$, of the one parameter family $u_{x_0,r} \in L^2(\partial B_1)$. Precisely, we have the following elementary estimate.

Lemma 2.1. Let $u \in H_{loc}^1(\mathbb{R}^d)$ and $x_0 \in \mathbb{R}^d$. Suppose that there are constants $r_0 > 0$, $\gamma \in (0,1)$ and $I > 0$ such that

(2.3)
$$
\int_0^{r_0} \frac{E(u_{x_0,r})}{r^{1+\gamma}} dr \le I.
$$

Then, there is a unique function $u_{x_0} \in L^2(\partial B_1)$ such that

$$
||u_{r,x_0} - u_{x_0}||_{L^2(\partial B_1)}^2 \le \gamma^{-1} Ir^{\gamma} \quad \text{for every} \quad r \in (0,r_0).
$$

Proof. We set for simplicity, $x_0 = 0$ and $u_r := u_{x_0,r}$. Let $0 < r < R \le r_0$. Notice that, for any $x \in \partial B_1$, we have

$$
\frac{u(Rx)}{R} - \frac{u(rx)}{r} = \int_r^R \left(\frac{x \cdot (\nabla u)(sx)}{s} - \frac{u(sx)}{s^2} \right) ds = \frac{1}{s} \int_r^R \left(x \cdot \nabla u_s(x) - u_s(x) \right) ds.
$$

Thus, by the Cauchy-Schwarz inequality, we obtain

$$
\int_{\partial B_1} |u_R - u_r|^2 d\mathcal{H}^{d-1} \leq \int_{\partial B_1} \left(\int_r^R \frac{1}{s} |x \cdot \nabla u_s - u_s| ds \right)^2 d\mathcal{H}^{d-1}
$$

\n
$$
\leq \int_{\partial B_1} \left(\int_r^R s^{\gamma - 1} ds \right) \left(\int_r^R \frac{1}{s^{1 + \gamma}} |x \cdot \nabla u_s - u_s|^2 ds \right) d\mathcal{H}^{d-1}
$$

\n
$$
\leq \frac{R^{\gamma} - r^{\gamma}}{\gamma} \int_r^R \frac{E(u_s)}{s^{1 + \gamma}} ds \leq \frac{R^{\gamma}}{\gamma} \int_0^{r_0} \frac{E(u_s)}{s^{1 + \gamma}} ds,
$$

which implies the claim by a standard argument. \Box

2.2. The one-phase boundary adjusted energy. Let $d \ge 2$ and $u \in H^1(B_1)$. For any $\Lambda > 0$, we define the one-phase Weiss' boundary adjusted energy as

(2.4)
$$
W_{\text{OP}}(u) := \int_{B_1} |\nabla u|^2 dx - \int_{\partial B_1} u^2 d\mathcal{H}^{d-1} + \Lambda |\{u > 0\} \cap B_1|.
$$

Let $r > 0$, $x_0 \in \mathbb{R}^2$ and $u \in H_{loc}^1(\mathbb{R}^d)$. The relation between W_{OP} and the excess E is given by the following formula, which holds for any function u and can be obtained by a direct computation (see [\[19\]](#page-21-2) and [\[16\]](#page-20-5)).

(2.5)
$$
\frac{\partial}{\partial r}W_{\text{OP}}(u_{x_0,r}) = \frac{d}{r}(W_{\text{OP}}(z_{x_0,r}) - W_{\text{OP}}(u_{x_0,r})) + \frac{1}{r}E(u_{x_0,r}),
$$

where $z_{x_0,r}$ denotes the one-homogeneous extension of the trace of $u_{x_0,r}$ in B_1 , that is,

(2.6)
$$
z_{x_0,r}(x) := |x| u_{x_0,r} (x/|x|) = \frac{|x|}{r} u (r x/|x|) , \text{ for every } x \in B_1.
$$

Theorem 2.2 (Epiperimetric inequality for W_{OP}). Let $d = 2$. Let $C_0 > 0$ be a given constant. There exists a constant $\varepsilon > 0$ such that: for every non-negative $c \in H^1(\partial B_1)$ satisfying Z $\int_{\partial B_1} c d\mathcal{H}^1 \geq C_0$, there exists a non-negative function $h \in H^1(B_1)$ such that $h = c$ on ∂B_1 and ∂B_2

(2.7)
$$
W_{\text{OP}}(h) - \Lambda \frac{\pi}{2} \le (1 - \varepsilon) \left(W_{\text{OP}}(z) - \Lambda \frac{\pi}{2} \right),
$$

where W_{OP} is given by [\(2.4\)](#page-4-1) and $z \in H^1(B_1)$ denotes the one-homogeneous extension of c into B_1 . Moreover, the competitor h has the following properties:

- (a) There is a universal numerical constant $C > 0$ such that $||h||_{H^1(B_1)} \leq C||c||_{H^1(\partial B_1)}$.
- (b) If $H_{x_0,\nu} := \{x \in \mathbb{R}^2 : (x-x_0) \cdot \nu \ge 0\}$, for some $x_0 \in \mathbb{R}^2$ and $\nu \in \partial B_1$, is a half-plane such that

(2.8)
$$
0 \in H_{x_0,\nu} \quad \text{and} \quad z = 0 \quad \text{on} \quad \mathbb{R}^2 \setminus H_{x_0,\nu},
$$

then we can choose h such that $h = 0$ on $\mathbb{R}^2 \setminus H_{x_0,\nu}$.

2.3. The two-phase boundary adjusted energy. For every $\Lambda_1, \Lambda_2 > 0$ and $v \in H^1(B_1)$, we define the two-phase Weiss' boundary adjusted energy as

$$
(2.9) \t WTP(v) = \int_{B_1} |\nabla v|^2 dx - \int_{\partial B_1} v^2 d\mathcal{H}^{d-1} + \Lambda_1 |\{v > 0\} \cap B_1| + \Lambda_2 |\{v < 0\} \cap B_1|.
$$

As in the one phase case, we have

(2.10)
$$
\frac{\partial}{\partial r} W_{\text{TP}}(u_{x_0,r}) = \frac{d}{r} (W_{\text{TP}}(z_{x_0,r}) - W_{\text{TP}}(u_{x_0,r})) + \frac{1}{r} E(u_{x_0,r}),
$$

where $z_{x_0,r}$ is given by [\(2.6\)](#page-4-2).

Theorem 2.3 (Epiperimetric inequality for W_{TP}). Let $d = 2$. For every $C_0 > 0$ there is $\varepsilon > 0$ such that: for every $c \in H^1(\partial B_1)$ such that ∂B_1 $c^+ d\mathcal{H}^1 \geq C_0$ and $\int_{\partial B_1}$ $c^- dH^1 \geq C_0$, there exists a function $h \in H^1(B_1)$ with $h = c$ on ∂B_1 such that

(2.11)
$$
W_{\text{TP}}(h) - (\Lambda_1 + \Lambda_2) \frac{\pi}{2} \leq (1 - \varepsilon) \left(W_{\text{TP}}(z) - (\Lambda_1 + \Lambda_2) \frac{\pi}{2} \right),
$$

where $z \in H^1(B_1)$ is the one-homogeneous extension of the trace of c to B_1 . Moreover, there is a universal numerical constant $C > 0$ such that $||h||_{H^1(B_1)} \leq C||c||_{H^1(\partial B_1)}$.

2.4. Almost-monotonicity and almost-minimality. Let $u \in H_{loc}^1(\mathbb{R}^d)$ and $x_0 \in \mathbb{R}^d$. For any $r > 0$, the function $u_{x_0,r}$ and $z_{x_0,r}$ are defined as in [\(2.1\)](#page-3-2) and [\(2.6\)](#page-4-2), respectively. In the next lemma we will show that a almost-minimality of u , with respect to radial perturbations, implies that the function $r \mapsto W_{\Box}(u_{x,r})$ is monotone up to a small error term (\Box stands for OP or TP).

Lemma 2.4 (Monotonicity of W_{\Box}). Let $u \in H_{loc}^1(\mathbb{R}^d)$ and $x_0 \in \mathbb{R}^d$. Suppose that there are constants $r_0 > 0$, $C > 0$ and $\delta > 0$ such that

(2.12)
$$
W_{\square}(u_{x_0,r}) \leq W_{\square}(z_{x_0,r}) + Cr^{\delta} \quad \text{for every} \quad r \in (0,r_0),
$$

where \Box stands for OP or TP. Then, the function

(2.13)
$$
r \mapsto W_{\Box}(u_{x_0,r}) + \frac{Cd}{\delta}r^{\delta},
$$

is non-decreasing on the interval $(0, r_0)$.

Proof. Using [\(2.5\)](#page-4-3) for \square =Op (resp. [\(2.10\)](#page-4-4) for \square =Tp), and the condition [\(2.12\)](#page-5-1) we get

$$
\frac{\partial}{\partial r}W_{\Box}(u_{x_0,r}) \geq \frac{d}{r}\left(W_{\Box}(z_{x_0,r}) - W_{\Box}(u_{x_0,r})\right) \geq C d r^{\delta-1},
$$

which gives (2.13) .

2.5. Epiperimetric inequality and energy decay. In this section we show how to use the epiperimetric inequality to obtain at once the decay for the energy $W_{\Box}(u_{x_0,r})$ and the convergence of $u_{x_0,r}$ in $L^2(\partial B_1)$. The argument is very general and we treat the cases \square = op and \square = TP simultaneously.

Lemma 2.5. Let $u \in H_{loc}^1(\mathbb{R}^d)$, $x_0 \in \mathbb{R}^d$ and W_{\Box} be as in [\(2.4\)](#page-4-1), if \Box = op, and [\(2.9\)](#page-4-5), if \Box = TP. Suppose that there are constants $r_0 \in (0,1)$, $C > 0$, $\delta > 0$ and $\varepsilon \in (0, \frac{\delta}{2d})$ $\frac{\delta}{2d+\delta}$) such that: (a) [\(2.12\)](#page-5-1) holds and the limit $\Theta_{\Box} := \lim_{r \to 0} W_{\Box}(u_{x_0,r})$ (which exists due to Lemma [2.4\)](#page-5-3) is finite; (b) for every $r \in (0, r_0)$ there is a function $h_{x_0,r} \in H^1(B_1)$ such that

(2.14)
$$
W_{\Box}(u_{x_0,r}) \leq W_{\Box}(h_{x_0,r}) + Cr^{\delta},
$$

and we have the epiperimetric inequality

$$
(2.15) \t W_{\Box}(h_{x_0,r}) - \Theta_{\Box} \le (1-\varepsilon)\big(W_{\Box}(z_{x_0,r}) - \Theta_{\Box}\big).
$$

Then, there is a unique function $u_{x_0} \in L^2(\partial B_1)$ such that

$$
||u_{r,x_0} - u_{x_0}||_{L^2(\partial B_1)}^2 \le \gamma^{-1} Ir^{\gamma} \quad \text{for every} \quad r \in (0, r_0),
$$

where
$$
\gamma = \frac{d\varepsilon}{1-\varepsilon}
$$
 and $I = r_0^{-\gamma} (W_{\square}(u_{x_0,r_0}) - \Theta_{\square}) + \frac{dC}{\delta - \gamma} r_0^{\delta - \gamma}$.

Proof. We use [\(2.5\)](#page-4-3) for \square =op (resp. [\(2.10\)](#page-4-4) for \square =TP), then the epiperimetric inequality [\(2.15\)](#page-5-4) and the almost-minimality condition [\(2.14\)](#page-5-5).

$$
\frac{\partial}{\partial r}\big(W_{\Box}(u_{x_0,r})-\Theta_{\Box}\big) \geq \frac{d}{r}\Big(\big(W_{\Box}(z_{x_0,r})-\Theta_{\Box}\big)-\big(W_{\Box}(u_{x_0,r})-\Theta_{\Box}\big)\Big) \geq \frac{d}{r}\Big(\frac{1}{1-\varepsilon}\big(W_{\Box}(h_{x_0,r})-\Theta_{\Box}\big)-\big(W_{\Box}(u_{x_0,r})-\Theta_{\Box}\big)\Big) \geq \frac{d}{r}\Big(\frac{\varepsilon}{1-\varepsilon}\big(W_{\Box}(u_{x_0,r})-\Theta_{\Box}\big)-Cr^{\delta}\Big),
$$

which implies that the function

$$
f(r) = \frac{W_{\Box}(u_{x_0,r}) - \Theta_{\Box}}{r^{\gamma}} + \frac{dC}{\delta - \gamma} r^{\delta - \gamma}
$$

is non-decreasing on $(0, r_0)$ for $\gamma = \frac{d\varepsilon}{1-\varepsilon}$, where we notice that $\gamma \leq \frac{\delta}{2}$ due to the choice $\varepsilon \leq \frac{\delta}{2d+\delta}$. In particular, using again (2.5) (resp. (2.10)), we get

$$
f'(r) \ge \frac{1}{r^{\gamma+1}} E(u_{x_0,r}),
$$

which integrated gives

$$
f(r_0) - f(s) \ge \int_s^{r_0} \frac{1}{r^{\gamma+1}} E(u_{x_0,r}) dr,
$$

for every $s \in (0, r_0)$. Now, notice that, up to choosing a bigger constant C in [\(2.14\)](#page-5-5), Lemma [2.4](#page-5-3) implies that $f(s) \geq 0$ for every $s > 0$. Thus, we get

$$
f(r_0) \ge \int_0^{r_0} \frac{1}{r^{\gamma+1}} E(u_{x_0,r}) dr,
$$

which is precisely [\(2.3\)](#page-3-3) with $I := f(r_0)$.

3. Change of variables and freezing of the coefficients

The arguments of the previous section, the monotonicity formula and the decay of the blow-up sequences, can be applied only in the case when the operator in J_{OP} (resp. J_{TP}) is the identity. Thus, in order to prove the regularity results Theorem [1.1](#page-1-0) and Theorem [1.5](#page-2-0) we need to change the coordinates and reduce to the case $A = Id$. We prove the main estimate of this section in Lemma [3.2](#page-7-0) below, but before we will introduce several notations.

Let $A = (a_{ij})_{ij} : B_2 \to Sym_2^+$ and $Q_{\text{OP}}, Q_{\text{TP}}^+$, $Q_{\text{TP}}^- : B_2 \to \mathbb{R}^+$ be as in the Introduction and note that we have

$$
||A_x^{1/2}|| \le M_A^{1/2} \text{ and } ||A_x^{-1/2}|| \le M_A^{1/2} \text{ for every } x \in B_2,
$$

where $||A|| = \sup \{ |Au| : u \in \mathbb{R}^2, |u| = 1 \}.$

Remark 3.1. We recall that if $M \in Sym_d^+$, then there is an orthogonal matrix P such that $PMP^t = \text{diag}(\lambda_1, \ldots, \lambda_d)$, where P^t is the transpose of P and $\text{diag}(\lambda_1, \ldots, \lambda_d)$ is the diagonal matrix with eigenvalues $\lambda_1, \ldots, \lambda_d$. We set $D = \text{diag}(\sqrt{\lambda}_1, \ldots, \sqrt{\lambda}_d)$ and define $M^{1/2} := P^t D P$.

For $x_0 \in B_2$ and $r > 0$ we define the functionals

$$
J_{\rm op}^{x_0}(v,r) := \int_{B_r} \left(|\nabla v|^2 + Q_{\rm op}(x_0) \mathbb{1}_{\{v > 0\}} \right) dx ;
$$

$$
J_{\rm TP}^{x_0}(v,r) := \int_{B_r} \left(|\nabla v|^2 + Q_{\rm TP}^+(x_0) \mathbb{1}_{\{v > 0\}} + Q_{\rm TP}^-(x_0) \mathbb{1}_{\{v < 0\}} \right) dx.
$$

For every $x_0 \in B_1$, we define the function

(3.1)
$$
F_{x_0}(x) := x_0 + A_{x_0}^{1/2}(x)
$$

and the half-plane $\mathcal{H}_{x_0} := \{ x \in \mathbb{R}^2 : F_{x_0}(x) \cdot e_2 > 0 \}$, where $e_2 = (0, 1)$.

Lemma 3.2. Let $L > 0$. There are constants $C > 0$ and $r_0 \in (0,1)$ (depending only on C_A , C_Q , M_A , M_Q , δ_A , δ_Q , δ_1 , C_1 and L) and $\delta = \min{\{\delta_A, \delta_Q, \delta_1\}}$ such that: if $u \in H^1(B_1)$ is a nonnegative L-Lipschitz continuous function and a almost-minimizer of J_{OP} in B_2^+ , $x_0 \in B_{r_0} \cap \partial\Omega_u$ and $\bar{u} = u \circ F_{x_0}$ (F_{x_0} is defined in [\(3.1\)](#page-6-1) above), then for every $r \in (0, r_0)$,

(3.2)
$$
J_{\text{op}}^{x_0}(\bar{u},r) \leq (1+Cr^{\delta})J_{\text{op}}^{x_0}(\bar{v},r) + Cr^{2+\delta},
$$

for every $\bar{v} \in H^1(B_r)$ such that $\bar{u} - \bar{v} \in H_0^1(B_r)$ and $\bar{v} = 0$ on $\mathbb{R}^2 \setminus \mathcal{H}_{x_0}$.

Moreover, there is a numerical constant $C_0 > 0$, such that

$$
W_{\text{op}}(\bar{u}_r) \le \begin{cases} W_{\text{op}}(\bar{z}_r) + C_0(M_{\text{A}}L^2 + M_Q)Cr^{\delta}, \\ W_{\text{op}}(\bar{h}_r) + C_0(M_{\text{A}}L^2 + M_Q)Cr^{\delta}, \end{cases}
$$

for every $r \in (0,r_0)$, where C is the constant from (3.2) , $\bar{u}_r(x) := \frac{1}{r} \bar{u}(rx)$, \bar{z}_r is the one homogeneous extension of \bar{u}_r in B_1 , \bar{h}_r is the competitor given by Theorem [2.2](#page-4-0) and $\Lambda = Q_{\text{or}}(x_0)$ in $(2.4).$ $(2.4).$

Proof. Let $x_0 \in \partial \Omega_u \cap B_1$, $r > 0$ and $\rho = M_A^{1/2}r$ and notice that this implies $F_{x_0}(B_r) \subset B_\rho(x_0)$. Let $\bar{u} = u \circ F_{x_0}$ and $\bar{v} = v \circ F_{x_0}$. Then, the Hölder continuity of A and $Q := Q_{\text{op}}$ and the ellipticity of A give

$$
\tilde{J}_{\text{OP}}(u, x_0, \rho) := \int_{B_{\rho}(x_0)} \left(a_{ij}(x_0) \, \partial_i u \, \partial_j u + Q(x_0) \mathbb{1}_{\{u > 0\}} \right) dx \leq J_{\text{OP}}(u, x_0, \rho) \n+ C_{\text{A}} M_{\text{A}} \rho^{\delta_{\text{A}}} \int_{B_{\rho}(x_0)} a_{ij}(x) \, \partial_i u \, \partial_j u \, dx + C_{\text{Q}} M_{\text{Q}} \rho^{\delta_{\text{Q}}} \int_{B_{\rho}(x_0)} Q(x) \mathbb{1}_{\{u > 0\}} dx \n\leq (1 + Cr^{\delta}) J_{\text{OP}}(u, x_0, \rho),
$$

for some positive constant $C > 0$. Analogously, we get the following estimate from below:

(3.3)
$$
\tilde{J}_{\text{OP}}(v, x_0, \rho) \ge (1 - Cr^{\delta}) J_{\text{OP}}(v, x_0, \rho).
$$

Putting the two estimates together and using the almost-minimality of u, we get

$$
\tilde{J}_{\text{OP}}(u, x_0, \rho) \leq \frac{1 + Cr^{\delta}}{1 - Cr^{\delta}} (1 + C_1 \rho^{\delta_1}) \tilde{J}_{\text{OP}}(v, x_0, \rho) + C_1 (1 + Cr^{\delta}) \rho^{2 + \delta_1}.
$$

Now, notice that by the choice of the function F_{x_0} we have the identity

$$
|\nabla \bar{u}|^2(x) = a_{ij}(x_0) \, \partial_i u(F_{x_0}(x)) \, \partial_j u(F_{x_0}(x)), \qquad x \in B_{M_A^{-1/2}}.
$$

Therefore, a change of coordinates and the estimate [\(3.3\)](#page-7-2) give

$$
\int_{F_{x_0}^{-1}(B_{\rho}(x_0))} \left(|\nabla \bar{u}|^2 + Q(x_0) \mathbb{1}_{\{\bar{u} > 0\}} \right) dx = \det \left(A_{x_0}^{-1/2} \right) \tilde{J}_{\text{op}}(u, x_0, \rho)
$$
\n
$$
\leq (1 + Cr^{\delta}) \int_{F_{x_0}^{-1}(B_{\rho}(x_0))} \left(|\nabla \bar{v}|^2 + Q(x_0) \mathbb{1}_{\{\bar{v} > 0\}} \right) dx + Cr^{2+\delta},
$$

for some other positive constant $C > 0$. Finally, observing that $B_r \subset F_{x_0}^{-1}(B_\rho(x_0))$ we get

$$
J_{\rm op}^{x_0}(\bar{u},r) \le (1+Cr^{\delta})J_{\rm op}^{x_0}(\bar{v},r) + Cr^{2+\delta} + Cr^{\delta}J_{\rm op}^{x_0}(\bar{u},\mathbf{M}_{\rm A}^{-1/2}r),
$$

which gives [\(3.2\)](#page-7-1) since \bar{u} is Lipschitz continuous with Lipschitz constant $\|\nabla \bar{u}\|_{L^{\infty}} = M_A^{1/2}L$.

We next notice that we have the scaling

$$
J^{x_0}_{\rm OP}(\bar u_r, 1) = \frac{1}{r^2} J^{x_0}_{\rm OP}(\bar u, r).
$$

Thus, the almost-minimality inequality [\(3.2\)](#page-7-1) translates in

(3.4)
$$
J_{\text{op}}^{x_0}(\bar{u}_r, 1) \le (1 + Cr^{\delta})J_{\text{op}}^{x_0}(\bar{v}_r, 1) + Cr^{\delta}.
$$

Let $C_{\rm E} > 0$ be the constant from Theorem [2.2.](#page-4-0) Then, since \bar{u} is Lipschitz continuous, we have

$$
\int_{B_1} |\nabla \bar{h}_r|^2 dx \leq C_{\mathcal{E}} \int_{\partial B_1} \left(|\nabla \bar{u}_r|^2 + \bar{u}_r^2 \right) dx \leq C_0 M_{\mathcal{A}} L^2,
$$

where C_0 is a numerical constant and \bar{h}_r is the competitor from Theorem [2.2.](#page-4-0) Taking \bar{h}_r as a competitor in [\(3.4\)](#page-7-3), we obtain

$$
J_{\rm OP}^{x_0}(\bar{u}_r, 1) \le J_{\rm OP}^{x_0}(\bar{h}_r, 1) + Cr^{\delta} \Big(\int_{B_1} |\nabla \bar{h}_r|^2 dx + Q(x_0)|B_1| \Big) + Cr^{\delta}
$$

$$
\le J_{\rm OP}^{x_0}(\bar{h}_r, 1) + Cr^{\delta} \Big(C_0 M_{\rm A} L^2 + M_{\rm Q}|B_1| \Big) + Cr^{\delta},
$$

which concludes the proof, the case $\bar{v}_r = \bar{z}_r$ being analogous.

An analogous result, with essentially the same proof holds in the two-phase case.

Lemma 3.3. Let $L > 0$. There are constants $C > 0$ and $r_0 \in (0,1)$ (depending only on C_A , C_O , M_A , M_Q , δ_A , δ_Q , δ_1 , C_2 and L) and $\delta = \min\{\delta_A, \delta_Q, \delta_2\}$ such that: if $u \in H^1(B_1)$ is a L-Lipschitz continuous function and a almost-minimizer of J_{TP} in B_2 , $x_0 \in B_{r_0} \cap \partial \Omega_u$ and $\bar{u} = u \circ F_{x_0}$, then we have that for every $r \in (0, r_0)$,

(3.5)
$$
J_{\text{TP}}^{x_0}(\bar{u}, r) \le (1 + Cr^{\delta})J_{\text{TP}}^{x_0}(\bar{v}, r) + Cr^{2+\delta},
$$

for every $\bar{v} \in H^1(B_r)$ such that $\bar{u} - \bar{v} \in H_0^1(B_r)$.

Moreover, there is a numerical constant $C_0 > 0$ such that

$$
W_{\text{TP}}(\bar{u}_r) \leq \begin{cases} W_{\text{TP}}(\bar{z}_r) + C_0(M_{\text{A}}L^2 + M_Q)Cr^{\delta}, \\ W_{\text{TP}}(\bar{h}_r) + C_0(M_{\text{A}}L^2 + M_Q)Cr^{\delta}, \end{cases}
$$

for every $r \in (0,r_0)$, where C is the constant from (3.5) , $\bar{u}_r(x) := \frac{1}{r} \bar{u}(rx)$, \bar{z}_r is the one homogeneous extension of \bar{u}_r in B_1 , \bar{h}_r is the competitor given by Theorem [2.3](#page-5-0) and $\Lambda_1 = Q_{\text{TP}}^+(x_0)$, $\Lambda_2 = Q_{\text{TP}}^-(x_0)$ in [\(2.9\)](#page-4-5).

Remark 3.4 (On the non-degeneracy). In [\[12\]](#page-20-0) David and Toro proved that Lipschitz continuous almost-minimizers to the one-phase and the two-phase functionals for the Laplacian are non degenerate (see [\[12,](#page-20-0) Theorem 10.1]). Note that their definition of almost-minimizer is slightly different from ours. However, their proof still holds in our case with small changes which come from the additional term $Cr^{2+\delta}$ of our definition. It follows from Lemma [3.2](#page-7-0) and Lemma [3.3](#page-8-2) that if u is a almost-minimizer of the functional J_{op} (resp. u is a almost-minmizer of J_{TP}) then u (resp. u_{+}) is non-degenerate with respect to A in the sense of the following definition.

Definition 3.5 (Non-degeneracy). Let $d \geq 2$ and $A : \mathbb{R}^d \to Sym_d^+$ be a given function. We say that the non-negative function $u \in H^1(B_2)$ is non-degenerate (with respect to A), if there are constants $\eta > 0$, $\varepsilon \in (0,1)$ and $r_0 > 0$ such that, for every $x_0 \in B_1$ and $r \in (0,r_0)$, the following implication holds:

$$
\int_{\partial B_r} u \circ F_{x_0} d\mathcal{H}^{d-1} < \eta r^d \qquad \Rightarrow \qquad u \circ F_{x_0} \equiv 0 \quad in \quad B_{\varepsilon r}(x_0),
$$

where $F_{x_0}(x) := x_0 + A_{x_0}^{1/2}(x)$.

4. Blow-up sequences and blow-up limits

Let $u \in H^1(B_2)$ be a Lipschitz continuous function. Let $(x_n)_{n\geq 1}$ be a sequence of points in $B_1 \cap \partial \Omega_u$ converging to some $x_0 \in B_1 \cap \partial \Omega_u$, and $(r_n)_{n \geq 1}$ be an infinitesimal sequence in $(0, 1)$. Then, the sequence u_{x_n,r_n} is uniformly Lipschitz in every compact subset of \mathbb{R}^2 . Thus, up to extracting a subsequence, there is a Lipschitz function $u_0 : \mathbb{R}^2 \to \mathbb{R}$ such that

$$
\lim_{n \to \infty} u_{x_n, r_n} = u_0,
$$

where u_{r_n,x_n} is defined in [\(2.1\)](#page-3-2) and the convergence is uniform on every compact subset of \mathbb{R}^2 .

Definition 4.1. If [\(4.1\)](#page-8-3) holds, we will say that u_{x_n,r_n} is a blow up sequence (with fixed center, if $x_n = x_0$, for every $n \ge 1$). If the center is fixed, we will say that u_0 is a blow-up limit at x_0 .

We summarize the main properties of the blow-up sequences and the blow-up limits in the following two propositions. We notice that Proposition [4.2](#page-9-0) holds in every dimension $d \geq 2$, while Proposition [4.3](#page-9-1) is known to hold only for $2 \le d \le 4$.

Proposition 4.2 (Convergence of the blow-up sequences). Let $u \in H^1(B_2)$ be as in Theorem [1.1](#page-1-0) or Theorem [1.5](#page-2-0) and let $u_n := u_{r_n,x_n}$ be a blow-up sequence converging to some $u_0 \in H^1_{loc}(\mathbb{R}^2)$. Then:

- (1) the sequence u_n converges strongly to u_0 in $H_{loc}^1(\mathbb{R}^2)$;
- (2) the sequences of characteristic functions $\mathbb{1}_{\{u_n>0\}}$ and $\mathbb{1}_{\{u_n<0\}}$ converge in $L^1_{loc}(\mathbb{R}^2)$ to $\mathbb{1}_{\{u_0>0\}}$ and $\mathbb{1}_{\{u_0<0\}}$, respectively.

Proposition 4.3 (Classification of the blow-up limits). Let $u \in H^1(B_2)$ be as in Theorem [1.1](#page-1-0) or Theorem [1.5.](#page-2-0) Let $x_0 \in \partial \Omega_u \cap B_1$ and $u_0 \in H_{loc}^1(\mathbb{R}^2)$ be a blow-up limit of u at x_0 .

(OP) If u is as in Theorem [1.1](#page-1-0) and $x_0 \in \partial\Omega_u \cap B_1^+$, then u_0 is of the form

(4.2)
$$
u_0(x) = Q_{\text{or}}^{1/2}(x_0) \max\left\{0, x \cdot A_{x_0}^{-1/2}[\nu]\right\}, \text{ where } \nu \in \partial B_1.
$$

(OP-c) If u is as in Theorem [1.1](#page-1-0) and $x_0 \in \partial \Omega_u \cap \partial H \cap B_1$, then u_0 is of the form

(4.3)
$$
u_0(x) = \mu \max\left\{0, x \cdot A_{x_0}^{-1/2}[\nu]\right\},\,
$$

where $\mu \geq Q_{\rm OP}^{1/2}(x_0)$ and $\nu \in \partial B_1$ is such that $A_{x_0}^{-1/2}[\nu]$ is normal to ∂H and pointing inwards.

(TP) If u is as in Theorem [1.5](#page-2-0) and $x_0 \in \partial \Omega_u^+ \cap \partial \Omega_u^- \cap B_1$, then u_0 is of the form

(4.4)
$$
u_0(x) = \mu_+ \max \left\{ 0, x \cdot A_{x_0}^{-1/2}[\nu] \right\} + \mu_- \min \left\{ 0, x \cdot A_{x_0}^{-1/2}[\nu] \right\},
$$

for some $\nu \in \partial B_1$ and some $\mu_+, \mu_- > 0$ such that $\mu_+^2 - \mu_-^2 = Q_{\text{TP}}^+(x_0) - Q_{\text{TP}}^-(x_0)$ and $\mu^2_+ \geq Q^+_{\text{TP}}(x_0), \, \mu^2_- \geq Q^-_{\text{TP}}(x_0).$

The proof of Proposition [4.2](#page-9-0) follows by a standard variational argument that only uses the almost-minimality of u; for more details, we refer to $\begin{bmatrix} 1 \end{bmatrix}$ (see also $\begin{bmatrix} 16 \end{bmatrix}$). Proposition [4.3](#page-9-1) follows by the optimality of the blow-up limits and the Weiss' monotonicity formula (Lemma [2.4\)](#page-5-3). We will need the following definition.

Definition 4.4 (Global solutions). Let $u : \mathbb{R}^2 \to \mathbb{R}$, $u \in H_{loc}^1(\mathbb{R}^2)$ be given.

(OP) We say that u is a global solution of the one-phase Bernoulli problem, if: $u \geq 0$ and, for every ball $B := B_R(x_0) \subset \mathbb{R}^2$, we have

(4.5)
$$
\int_{B} |\nabla u|^2 dx + \Lambda |\{u > 0\} \cap B| \le \int_{B} |\nabla v|^2 dx + \Lambda |\{v > 0\} \cap B|,
$$

for every $v \in H^1(B)$ such that $u - v \in H_0^1(B)$.

- (OP-c) We say that u is a global solution of the one-phase constrained Bernoulli problem in H, if $u \geq 0$ on H , $u = 0$ on $\mathbb{R}^2 \setminus H$ and (4.5) holds, for every ball $B := B_R(x_0) \subset \mathbb{R}^2$ and every $v \in H^1(B)$ such that $u - v \in H_0^1(B)$ and $\{v > 0\} \subset H$.
	- (TP) We say that u is a global solution of the two-phase Bernoulli problem if, for every ball $B := B_R(x_0) \subset \mathbb{R}^2$, we have

$$
(4.6) \qquad \int_{B} \left(|\nabla u|^{2} + \Lambda_{1} \mathbb{1}_{\{u>0\}} + \Lambda_{2} \mathbb{1}_{\{u<0\}} \right) dx \le \int_{B} \left(|\nabla v|^{2} + \Lambda_{1} \mathbb{1}_{\{v>0\}} + \Lambda_{2} \mathbb{1}_{\{v<0\}} \right) dx,
$$

for every $v \in H^1(B)$ such that $u - v \in H_0^1(B)$.

Lemma 4.5 (Optimality of the blow-up limits). Let $u \in H^1(B_2)$ be as in Theorem [1.1](#page-1-0) or Theorem [1.5](#page-2-0) and let $u_n := u_{r_n,x_0}$ be a blow-up sequence converging to the blow-up limit $u_0 \in H^1_{loc}(\mathbb{R}^2)$. Then, we have:

- (OP) If u is as in Theorem [1.1](#page-1-0) and $x_0 \in \partial \Omega_u \cap B_1^+$, then $u_0 \circ A_{x_0}^{1/2}$ is a global solution of the one-phase problem with $\Lambda = Q_{\text{op}}(x_0)$.
- (OP-c) If u is as in Theorem [1.1](#page-1-0) and $x_0 \in \partial \Omega_u \cap \partial H \cap B_1$, then, up to a rotation, $u_0 \circ A_{x_0}^{1/2}$ is a global solution of the constrained one-phase problem with $\Lambda = Q_{\text{or}}(x_0)$.
	- (TP) If u is as in Theorem [1.5](#page-2-0) and $x_0 \in \partial \Omega_u^+ \cap \partial \Omega_u^- \cap B_1$, then $u_0 \circ A_{x_0}^{1/2}$ is a global solution of the two-phase problem with $\Lambda_1 = Q_{\text{TP}}^+(x_0)$ and $\Lambda_2 = Q_{\text{TP}}^-(x_0)$.

Recall that the function $\bar{u} = u \circ F_{x_0}$, where F_{x_0} is as in [\(3.1\)](#page-6-1), is a almost-minimizer of the functional $J_{\rm op}^{x_0}$ (Lemma [3.2\)](#page-7-0). We then refer to Lemma 4.6 in [\[16\]](#page-20-5) applied to \bar{u} for the proof of Lemma [4.5.](#page-10-0) It is also worth mentioning that the strong convergence of the blow-up sequences and the optimality of the blow-up limits are equivalent.

Lemma 4.6 (Homogeneity of the blow-up limits). Let $u \in H^1(B_2)$ be as in Theorem [1.1](#page-1-0) or Theorem [1.5.](#page-2-0) Let $x_0 \in B_1 \cap \partial \Omega_u$ and let u_{x_0,r_n} be a blow-up sequence converging to a blow-up limit u_0 . Then, u_0 is one-homogeneous.

Proof. Assume that $x_0 = 0$ and set $\bar{u} = u \circ F_{x_0}$. Then

$$
u_{x_0,r} = \bar{u}_r \circ A_{x_0}^{-1/2}
$$
, where $\bar{u}_r(x) := \frac{\bar{u}(rx)}{r}$.

We first notice that by Lemma [3.2,](#page-7-0) Lemma [2.4](#page-5-3) and the Lipschitz continuity of u , we get that the limit $\Theta_{\Box} := \lim_{n \to \infty} W_{\Box}(\bar{u}_r), \Box = \text{OP}, \text{TP},$ exists and is finite. Now the strong convergence of \bar{u}_{r_n} to $r\rightarrow 0$ $\bar{u}_0 := u_0 \circ A_{x_0}^{1/2}$ (Proposition [4.2\)](#page-9-0) implies that, for every $s > 0$, we have $\Theta_{\Box} := \lim_{r \to 0} W_{\Box}(\bar{u}_r) = \lim_{n \to \infty} W_{\Box}(\bar{u}_{r_n}) = \lim_{n \to \infty} W_{\Box}(\bar{u}_{r_ns}) = \lim_{n \to \infty} W_{\Box}((\bar{u}_{r_n})_s) = W_{\Box}((\bar{u}_0)_s).$

In particular, $s \mapsto W_{\Box}(\bar{u}_0, s)$ is constant. Now, since \bar{u}_0 is a global solution (Lemma [4.5\)](#page-10-0), [\(2.5\)](#page-4-3) and [\(2.10\)](#page-4-4) imply that $E((\bar{u}_0)_s) = 0$, for every $s > 0$. Thus we have $x \cdot \nabla \bar{u}_0 = \bar{u}_0$ in \mathbb{R}^2 , which implies that \bar{u}_0 (and thus, u_0) is one-homogeneous.

Proof of Proposition [4.3.](#page-9-1) We now notice that $\bar{u}_0 = u_0 \circ A_{x_0}^{1/2} : B_1 \to \mathbb{R}$ is one-homogeneous and harmonic on the cone $B_1 \cap {\bar{u}_0 \neq 0}$. Thus, the trace of \bar{u}_0 on the sphere satisfies the equation

$$
-\Delta_{\mathbb{S}}\bar{u}_0 = (d-1)\bar{u}_0 \text{ on } \mathbb{S}^{d-1} \cap {\bar{u}_0 \neq 0},
$$

where in dimension two the spherical Laplacian $\Delta_{\mathbb{S}}$ is simply the second derivative and $d-1=1$. Thus, \bar{u}_0 is of the form $\bar{u}_0(\theta) = \sin(\theta + \theta_0)$, $\theta \in \mathbb{S}^2$, for some constant θ_0 . This implies that ${\bar{u}}_0 \neq 0$ is a union of intervals of length π . In the one-phase case, since u is non-degenerate (see Remark [3.4\)](#page-8-4), this implies that \bar{u}_0 is of the form [\(4.2\)](#page-9-3), for some constant $\mu(x_0)$. Now, an internal variation argument (see [\[1\]](#page-20-1)) implies that $\mu(x_0) = Q_{\text{op}}^{1/2}(x_0)$, if $x_0 \in H \cap B_1^+$, and $\mu(x_0) \geq Q_{\text{op}}^{1/2}(x_0)$, if $x_0 \in \partial H \cap B_1$. The two-phase case follows again by an internal variation argument (see [\[2\]](#page-20-2)). □

Finally, we prove a uniqueness result for the one- and two-phase (Theorem [1.1](#page-1-0) and Theorem [1.5\)](#page-2-0) blow-up limits. This is the only result of this section that cannot be immediately extended to higher dimension. This is due to the fact that the epiperimetric inequality (Theorem [2.2](#page-4-0) and Theorem [2.3\)](#page-5-0) is known (for the moment) only in dimension two.

Proposition 4.7 (Uniqueness of the blow-up and rate of convergence of the blow-up sequences). Let $u : B_2 \to \mathbb{R}$ be as in Theorem [1.1](#page-1-0) or Theorem [1.5.](#page-2-0) There are constants $C > 0$, $\gamma > 0$ and $r_0 > 0$ such that the following claims do hold.

(OP) If u is as in Theorem [1.1,](#page-1-0) then for every $x_0 \in \partial \Omega_u \cap B_1$, there is a unique blow-up $u_{x_0}: \mathbb{R}^2 \to \mathbb{R}$ (of the form [\(4.2\)](#page-9-3) or [\(4.3\)](#page-9-4)) such that

(4.7)
$$
||u_{x_0,r} - u_{x_0}||_{L^{\infty}(B_1)} \leq C r^{\gamma} \text{ for every } r \in (0,r_0).
$$

(TP) If u is as in Theorem [1.5,](#page-2-0) then for every $x_0 \in \partial \Omega_u^+ \cap \partial \Omega_u^- \cap B_1$, there is a unique blow-up $u_{x_0}: \mathbb{R}^2 \to \mathbb{R}$ (of the form [\(4.4\)](#page-9-5)) such that

(4.8)
$$
||u_{x_0,r} - u_{x_0}||_{L^{\infty}(B_1)} \leq C r^{\gamma} \text{ for every } r \in (0,r_0).
$$

Proof. Let u be as in (OP) and $x_0 \in \partial\Omega_u \cap B_1$. We set $\bar{u} = u \circ F_{x_0}$ and $\bar{u}_r(x) := \frac{\bar{u}(rx)}{r}$, and we notice that $\bar{u}_r = u_{x_0,r} \circ A_{x_0}^{1/2}$. By Lemma [3.2](#page-7-0) and Lemma [2.4,](#page-5-3) $r \mapsto W_{\text{op}}(\bar{u}_r) + Cr^{\delta}$ is monotone. On the other hand, the homogeneity of the blow-up limits, imply that

$$
\Theta_{\text{op}} := \lim_{r \to 0} W_{\text{op}}(\bar{u}_r) = \frac{\pi}{2} Q_{\text{op}}(x_0).
$$

Thus, by the epiperimetric inequality (Theorem [2.2\)](#page-4-0), Lemma [3.2](#page-7-0) and Lemma [2.5,](#page-5-6) we have that there exists a one-homogeneous function \bar{u}_0 such that, for $r > 0$ small enough,

$$
\|\bar{u}_r - \bar{u}_0\|_{L^2(\partial B_1)} \le Cr^{\gamma_0/2}
$$

,

where γ_0 is the constant from Lemma [2.5.](#page-5-6) Integrating in r, we get that

$$
\|\bar{u}_r - \bar{u}_0\|_{L^2(B_1)} \le Cr^{\gamma_0/2}.
$$

Now, since $\bar{u}_r = u_{x_0,r} \circ A_{x_0}^{1/2}$ and $A_{x_0}^{1/2}$ is invertible, we get

$$
||u_{x_0,r} - u_{x_0}||_{L^2(B_1)} \le Cr^{\gamma_0/2},
$$

where $u_{x_0} = \bar{u}_0 \circ A_{x_0}^{1/2}$. Finally, we notice that the Lipschitz continuity of u implies that there is an universal bound on $\|\nabla u_{x_0,r}\|_{L^{\infty}(B_1)}$ and $\|\nabla u_{x_0}\|_{L^{\infty}(B_1)}$. Thus, we get [\(4.7\)](#page-10-1) with $\gamma = \gamma_0/4$. The proof of (TP) is analogous.

Remark 4.8. We notice that the above result does not hold at the one-phase points $x_0 \in \partial \Omega_u^+ \setminus \partial \Omega_u^$ of the solutions u of the two-phase problem (Theorem [1.5\)](#page-2-0). This is due to the fact that the positive part u_+ is not a solution of the one-phase problem in the balls $B_r(x_0)$ that have nonempty intersection with the negative phase Ω_u^- . In fact, the blow-up limit u_{x_0} (of u at x_0) is still unique, but the decay estimate [\(4.7\)](#page-10-1) holds only for $r < \frac{1}{2}$ dist (x_0, Ω_u^{-}) .

5. Regularity of the one-phase free boundaries. Proof of Theorem [1.1](#page-1-0)

Let $u \in H^1(B_2)$, $u \ge 0$, be as in Theorem [1.1.](#page-1-0) By Proposition [4.7](#page-10-2) we have that, for every $x_0 \in \partial \Omega_u \cap B_1$, there is a unique blow-up limit of u at x_0 . We denote it by

$$
u_{x_0}(x) = \mu(x_0) \max\{0, \nu_{x_0} \cdot x\},\,
$$

where ν_{x_0} is of the form $A_{x_0}^{1/2}[\nu]$, for some $\nu \in \partial B_1$; and $\mu(x_0)$ is such that $Q_{\text{OP}}(x_0) \leq \mu^2(x_0) \leq \mu^2(x_0)$ $M_{\rm A}L^2$, where L is the Lipschitz constant of u. We also notice that

$$
\mu(x_0) = Q_{\text{op}}^{1/2}(x_0) \quad \text{whenever} \quad x_0 \in \partial \Omega_u \cap B_1^+.
$$

Moreover, for every point $x_0 \in \partial \Omega_u \cap B_1$, we define the half-plane

$$
H_{x_0} := \{ x \in \mathbb{R}^2 : x \cdot \nu_{x_0} > 0 \}.
$$

We first prove the following:

Lemma 5.1. Let u be as in Theorem [1.1.](#page-1-0) There are constants $C > 0$, $\gamma > 0$ and $r_0 > 0$ such that, for every $x_0 \in \partial \Omega_u \cap B_1$, we have

(5.1)
$$
\Omega_{x_0,r} \cap B_1 \supset \{x \in B_1 : x \cdot \nu_{x_0} > Cr^{\gamma}\}\
$$
 and $\Omega_{x_0,r} \cap \{x \in B_1 : x \cdot \nu_{x_0} < -Cr^{\gamma}\} = \emptyset$,
for every $r \in (0,r_0)$, where $\Omega_{x_0,r} := \{u_{x_0,r} > 0\}$.

Proof. The first part of [\(5.1\)](#page-11-1) follows by the uniform convergence of the blow-up sequence $u_{x_0,r}$ (Proposition [4.7,](#page-10-2) equation [\(4.7\)](#page-10-1)) and the form of the blow-up limit u_{x_0} . The second part of [\(5.1\)](#page-11-1) follows again by [\(4.7\)](#page-10-1), the fact that $u_{x_0} \equiv 0$ on $B_1 \setminus H_{x_0}$ and by the non-degeneracy of u, which can be written as

If
$$
u_{x_0,r}(y_0) > 0
$$
, then $||u_{x_0,r}||_{L^{\infty}(B_s(y_0))} \ge Cs$, for every $s \in (0,1)$,

for some $C > 0$.

Lemma 5.2. Let u be as in Theorem [1.1.](#page-1-0) There are constants $R, \alpha \in (0,1)$ and $C > 0$ such that, for every $x_0, y_0 \in \partial \Omega_u \cap B_R$, we have

(5.2)
$$
|\nu_{x_0} - \nu_{y_0}| \le C|x_0 - y_0|^\alpha \quad and \quad |\mu(x_0) - \mu(y_0)| \le C|x_0 - y_0|^\alpha.
$$

Proof. Let $\gamma \in (0,1)$ be the exponent from Proposition [4.7](#page-10-2) and let $\alpha := \frac{\gamma}{1+1}$ $\frac{\gamma}{1+\gamma}$. Let $x_0, y_0 \in$ $B_R \cap \partial \Omega_u$, where we choose R such that $(2R)^{1-\alpha} \leq r_0$, where r_0 is the constant from Proposition [4.7.](#page-10-2) We set $r := |x_0 - y_0|^{1-\alpha}$. Recall that u is Lipschitz continuous and set $L = ||\nabla u||_{L^{\infty}}$. Then, for every $x \in B_1$, we have

$$
|u_{x_0,r}(x) - u_{y_0,r}(x)| = \frac{1}{r}|u(x_0 + rx) - u(y_0 + rx)| \le L\frac{|x_0 - y_0|}{r} = L|x_0 - y_0|^\alpha.
$$

and then, by an integration on B_1 , we get

$$
||u_{x_0,r} - u_{y_0,r}||_{L^2(B_1)} \le |B_1|^{1/2} L |x_0 - y_0|^\alpha.
$$

On the other hand, by the choice of R, we have that $r \le r_0$; applying Proposition [4.7,](#page-10-2) we get

$$
||u_{x_0,r} - u_{x_0}||_{L^2(B_1)} \le Cr^{\gamma}
$$
 and $||u_{y_0,r} - u_{y_0}||_{L^2(B_1)} \le Cr^{\gamma}$.

Thus, by the triangular inequality and the fact that $r^{\gamma} = |x_0 - y_0|^{\alpha}$, we obtain

(5.3)
$$
||u_{x_0} - u_{y_0}||_{L^2(B_1)} \le (|B_1|^{1/2}L + 2C)|x_0 - y_0|^{\alpha}.
$$

The conclusion now follows by a general argument. Indeed, for any pair of vectors $v_1, v_2 \in \mathbb{R}^2$, we have

(5.4)
\n
$$
|v_1 - v_2| = \left(\frac{2}{\pi} \int_{B_1} |v_1 \cdot x - v_2 \cdot x|^2 dx\right)^{1/2}
$$
\n
$$
\leq \left(\int_{B_1} |(v_1 \cdot x)_{+} - (v_2 \cdot x)_{+}|^2 dx\right)^{1/2} + \left(\int_{B_1} |(v_1 \cdot x)_{-} - (v_2 \cdot x)_{-}|^2 dx\right)^{1/2}
$$
\n
$$
= 2 \left(\int_{B_1} |(v_1 \cdot x)_{+} - (v_2 \cdot x)_{+}|^2 dx\right)^{1/2}.
$$

Applying the above estimate to $v_1 = \mu(x_0)\nu_{x_0}$ and $v_2 = \mu(y_0)\nu_{y_0}$, and using [\(5.3\)](#page-12-0), we get [\(5.2\)](#page-12-1). \Box **Proof of Theorem [1.1.](#page-1-0)** We first claim that, for every $\varepsilon > 0$, there exists $\rho > 0$ such that, for $x_0 \in \partial \Omega_u \cap B_\rho$ we have

(5.5)
$$
u > 0 \text{ on } C^+(x_0, \varepsilon) \cap B_{\rho}(x_0) \quad \text{and} \quad u = 0 \text{ on } C^-(x_0, \varepsilon) \cap B_{\rho}(x_0),
$$

where

$$
C^{\pm}(x_0,\varepsilon):=\left\{x\in\mathbb{R}^2\backslash\{0\}: \pm\nu_{x_0}\cdot(x-x_0)\geq\varepsilon|x-x_0|\right\}.
$$

Indeed, the flatness estimate [\(5.1\)](#page-11-1) implies [\(5.5\)](#page-12-2) by taking ρ such that $C\rho^{\gamma} \leq \varepsilon$, where C and γ are the constants from Lemma [5.1.](#page-11-2)

We now fix $x_0 \in B_1 \cap \partial \Omega_u$. Without loss of generality we can suppose that $x_0 = 0$ and $H_{x_0} = \{(s, t) \in \mathbb{R}^2 : t > 0\}$. Now, let $\varepsilon \in (0, 1)$ and $\rho > 0$ as in [\(5.5\)](#page-12-2) and set $\delta = \rho\sqrt{1-\varepsilon^2}$. By [\(5.5\)](#page-12-2) we have for every $s \in (-\delta, \delta)$

- the set $\mathcal{S}_{+}^{s} := \{t \in (-\delta, \delta) \; : \; u(s,t) > 0\}$ contains the interval $(\rho \varepsilon, \delta)$;
- the set $S_0^s := \{t \in (-\delta, \delta) : u(s, t) = 0\}$ contains the interval $(-\delta, -\rho \varepsilon)$.

This implies that the function

$$
g(s) := \max\{t \in \mathbb{R} : u(s, t) > 0\}
$$

is well defined and such that

 $SQ_{\delta} \cap \Omega_u = \{(s, t) \in SQ_{\delta} : g(s) < t\}$ and $SQ_{\delta} \setminus \Omega_u = \{(s, t) \in SQ_{\delta} : g(s) \ge t\},$

where $SQ_{\delta} = (-\delta, \delta) \times (-\delta, \delta)$. Now, the flatness condition [\(5.1\)](#page-11-1) implies that g is differentiable on $(-\delta, \delta)$. Furthermore, since ν is Hölder continuous, we deduce that g is a function of class $C^{1,\alpha}$. This concludes the proof. $C^{1,\alpha}$. This concludes the proof.

6. Regularity of the two-phase free boundaries. Proof of Theorem [1.5](#page-2-0)

Let u be as in Theorem [1.5.](#page-2-0) Then, by Proposition [4.7,](#page-10-2) at every point $x_0 \in \partial \Omega_u \cap B_1$ there is a unique blow-up limit u_{x_0} given by

$$
u_{x_0}(x) = \mu_+(x_0) \max\{0, x \cdot \nu_{x_0}\}, \quad \text{if } x_0 \in \Gamma_+ := (\partial \Omega_u^+ \backslash \partial \Omega_u^-) \cap B_1;
$$

$$
u_{x_0}(x) = \mu_-(x_0) \min\{0, x \cdot \nu_{x_0}\}, \quad \text{if } x_0 \in \Gamma_- := (\partial \Omega_u^- \backslash \partial \Omega_u^+) \cap B_1;
$$

$$
u_{x_0}(x) = \mu_+(x_0) \max\{0, x \cdot \nu_{x_0}\} + \mu_-(x_0) \min\{0, x \cdot \nu_{x_0}\}, \quad \text{if } x_0 \in \Gamma_{\text{TP}} := \partial \Omega_u^+ \cap \partial \Omega_u^- \cap B_1,
$$

where $\nu_{x_0} \in \mathbb{R}^2$ is of the form $A_{x_0}^{-1/2}[\nu]$, for some $\nu \in \partial B_1$, and $\mu_+(x_0)$ and $\mu_-(x_0)$ are positive and such that $Q_{\text{TP}}^{\pm}(x_0) \leq \mu_{\pm}^2(x_0) \leq M_{\text{A}}L^2$, where $L = ||\nabla u||_{L^{\infty}(B_2)}$ is the Lipschitz constant of u, and

$$
\mu_{\pm}^{2}(x_{0}) = Q_{\text{TP}}^{\pm}(x_{0}), \text{ if } x_{0} \in \Gamma_{\pm}
$$

$$
\mu_{+}^{2}(x_{0}) - \mu_{-}^{2}(x_{0}) = Q_{\text{TP}}^{+}(x_{0}) - Q_{\text{TP}}^{-}(x_{0}), \text{ if } x_{0} \in \Gamma_{\text{TP}}.
$$

Notice that Corollary [1.3](#page-1-1) already implies that the one-phase free boundaries Γ_+ and Γ_- are $C^{1,\alpha}$ regular. Thus, it remains to prove that $\partial\Omega_u^+$ and $\partial\Omega_u^-$ are smooth in a neighborhood of Γ_{TP} .

Lemma 6.1 (Flatness of the free boundary at the two-phase points). Let u be as in Theorem [1.5.](#page-2-0) There are constants $C > 0$, $\gamma > 0$ and $r_0 > 0$ such that, for every $x_0 \in \partial \Gamma_{\rm rp}$, we have

$$
(6.1) \quad \Omega_{x_0,r}^+ \cap B_1 \supset \{x \in B_1 : x \cdot \nu_{x_0} > Cr^{\gamma}\} \quad and \quad \Omega_{x_0,r}^- \cap B_1 \supset \{x \in B_1 : x \cdot \nu_{x_0} < -Cr^{\gamma}\},
$$

for every $r \in (0,r_0)$, where $\Omega_{x_0,r}^+ := \{u_{x_0,r} > 0\}$ and $\Omega_{x_0,r}^- := \{u_{x_0,r} < 0\}.$

Proof. Both the inclusions of [\(6.1\)](#page-13-1) follow by the uniform convergence of $u_{x_0,r}$ (Proposition [4.7,](#page-10-2) equation [\(4.8\)](#page-11-3)) to the blow-up limit u_{x_0} . .

Lemma 6.2. Let u be as in Theorem [1.5.](#page-2-0) There are constants $R, \alpha \in (0,1)$ and $C > 0$ such that, for every $x_0, y_0 \in \partial \Gamma_{\text{TP}} \cap B_R$, we have

(6.2)
$$
|\nu_{x_0} - \nu_{y_0}| \le C|x_0 - y_0|^\alpha \quad and \quad |\mu_{\pm}(x_0) - \mu_{\pm}(y_0)| \le C|x_0 - y_0|^\alpha.
$$

Proof. The proof follows step by step the one of Lemma [5.2.](#page-12-3)

Reasoning as in the one-phase case, and using Lemma [6.1](#page-13-2) and Lemma [6.2,](#page-13-3) one can prove that the two-phase free boundary Γ_{TP} is *contained* in a $C^{1,\alpha}$ curve. Unfortunately, this result by itself is not sufficient to deduce that $\partial\Omega_u^{\pm}$ are smooth. We now prove that the function u_+ (resp. u_-) is a solution of the one-phase free boundary problem

(6.3)
$$
-\text{div}(A\nabla u_+) = f_+
$$
 in Ω_u^+ , $|A_{x_0}^{1/2}\nabla u_+|(x_0) = \mu_+(x_0)$ for every $x_0 \in \partial\Omega_u^+$

where the boundary equation is understood in a classical sense. This is an immediate consequence of the following lemma which states that u_+ is differentiable in Ω_u^+ up to the boundary.

Lemma 6.3 (Differentiability at points of the free boundary). Let u be as in Theorem [1.5.](#page-2-0) We consider two cases.

(OP). For every $x_0 \in (\partial \Omega_u^+ \setminus \partial \Omega_u^-) \cap B_1$, u_+ is differentiable at x_0 and there is $r(x_0) > 0$ such that

$$
|u_+(x) - \mu_+(x_0)(x - x_0) \cdot \nu_{x_0}| \le C|x - x_0|^{1+\gamma} \quad \text{for every} \quad x \in B_{r(x_0)}(x_0) \cap \Omega_u^+.
$$

$$
\mathbf{r}^{\prime}
$$

(TP). There exists a universal constant $r_0 > 0$ such that for every $x_0 \in \partial \Omega_u^+ \cap \partial \Omega_u^- \cap B_1$, u_+ is differentiable in $B_{r_0}(x_0) \cap \Omega_u^+$, and

(6.4)
$$
\left| u_{+}(x) - \mu_{+}(x_{0})(x-x_{0}) \cdot \nu_{x_{0}} \right| \leq C|x-x_{0}|^{1+\gamma} \quad \text{for every} \quad x \in B_{r_{0}}(x_{0}) \cap \Omega_{u}^{+}.
$$

In particular, for every $x_0 \in \partial \Omega_u^+ \cap B_1$, we have $\nabla u_+(x_0) = \mu_+(x_0)\nu_{x_0}$.

Proof. The two cases are analogous. We will prove (TP). By Proposition [4.7,](#page-10-2) for every $r < r_0$, we have

$$
\|\max\{0, u_{x_0,r}\} - \max\{0, u_{x_0}\}\|_{L^{\infty}(B_1)} \le \|u_{x_0,r} - u_{x_0}\|_{L^{\infty}(B_1)} \le Cr^{\gamma}.
$$

Thus, using the flatness of the free boundary (Lemma [6.1\)](#page-13-2), we get for every $x \in B_1 \cap \{u_{x_0,r} > 0\}$

$$
|\max\{0, u_{x_0,r}(x)\} - \mu_+(x_0)x \cdot \nu_{x_0}| \leq |\max\{0, u_{x_0,r}(x)\} - \max\{0, u_{x_0}(x)\}| + \mu_+(x_0)|\min\{0, x \cdot \nu_{x_0}\}| \leq Cr^{\gamma}.
$$

Now, taking $r = |x - x_0|$ and rescaling the above inequality, we obtain [\(6.4\)](#page-14-0)

We notice that at the two-phase free boundary point the estimate (6.4) holds in a ball whose radius does not depend on the point. Moreover, on the two-phase free boundary the gradient has a universal modulus of continuity (see Lemma [6.2\)](#page-13-3). We next show that μ_{+} is Hölder continuous on $\partial Ω_u⁺$.

Lemma 6.4. The function $\mu_+ : \partial \Omega_u^+ \to \mathbb{R}$ is (locally) Hölder continuous.

Proof. We infer that lemma [6.4](#page-14-1) is a consequence of the following claim: if $(x_n)_{n>1}$ is a sequence of one-phase points, $x_n \in \Gamma_+$, converging to a two-phase point $y_0 \in \Gamma_{\text{TP}}$, then $\mu_+(y_0) = Q_+^{1/2}(y_0)$, where we set $Q_+ := Q_{\text{TP}}^+$. Indeed, we we notice that:

• on the set Γ_+ , we have $\mu_+ = Q_+^{1/2}$.

• for every $y_1, y_2 \in \Gamma_{\text{TP}}$, we have $|\mu_+(y_1) - \mu_+(y_2)| \le C |y_1 - y_2|^{\alpha}$.

By the first bullet μ_+ is Hölder continuous on the open subset Γ_+ of $\partial\Omega_u \cap B_1$. To prove that μ_+ is Hölder continuous on Γ_{TP} , let $y_0 \in \Gamma_{\text{TP}}$ and $x, y \in \partial \Omega_u^+ \cap B_{r_0}(y_0)$; we have to show that $|\mu_+(x) - \mu_+(y)| \leq C|x-y|^{\alpha}$. By the two bullets, it is obvious if either $x, y \in \Gamma_+$ or $x, y \in \Gamma_{\text{TP}}$. Then, assume that $x \in \Gamma_+$ and $y \in \Gamma_{\text{TP}}$ and denotes by y_1 the projection of y on the closed set Γ , where Γ is the set of points z in Γ_{TP} such that every neighborhood of z has non-empty intersection with Γ_+ . Note that, by definition of y_1 , we have $|y_1 - y| \le |x - y|$ and then $|x - y_1| \le 2|x - y|$. Therefore, using the triangular inequality and the claim we get

$$
|\mu_+(x) - \mu_+(y)| \leq |Q_+^{1/2}(x) - Q_+^{1/2}(y_1)| + |\mu_+(y_1) - \mu_+(y)| \leq C|x - y|^{\alpha}.
$$

We now prove the claim. Up to a linear change of coordinates we may suppose that $A_{y_0} = Id$. Denote by y_n the projection of x_n on the closed set $\partial \Omega_u^+ \cap \partial \Omega_u^-$ and set $r_n := |x_n - y_n|$. Since u is Lipschitz continuous, up to a subsequence, $u_n := u^+_{x_n,r_n}$ converges locally uniformly to some function u_{∞} . The absence of two-phase points in $B_{r_n}(x_n)$ implies that u_n is a solution of

$$
-\text{div}(A_n \nabla u_n) = r_n f_n \quad \text{in} \quad \{u_n > 0\} \cap B_1, \qquad |\nabla u_n| = q_n \quad \text{on} \quad \partial \{u_n > 0\} \cap B_1,
$$

where $A_n(x) := A(x_n + r_n x)$, $f_n(x) := f_+(x_n + r_n x)$ and $q_n(x) = Q_+^{1/2}(x_n + r_n x)|\nu_{x_n + r_n x}|$, where we recall that $\nu_{x_n+r_n}$ is of the form $A_{x_n}^{1/2}$ $x_{n+rn}^{1/2}[v]$, for some $\tilde{v} \in \partial B_1$. Passing to the limit as $n \to \infty$, we obtain that u_{∞} is a viscosity solution to

$$
-\Delta u_{\infty} = 0 \text{ in } \{u_{\infty} > 0\} \cap B_1, \qquad |\nabla u_{\infty}| = Q_+^{1/2}(y_0)|\nu_{y_0}| \text{ on } \partial\{u_{\infty} > 0\} \cap B_1.
$$

On the other hand, for every $\xi \in B_1$, we have

$$
u_{x_n,r_n}(\xi) = u_{y_n,r_n}(\xi + \xi_n), \text{ where } \xi_n := \frac{x_n - y_n}{r_n} \in \partial B_1,
$$

and, up to a subsequence, we can assume that ξ_n converges to some $\xi_\infty \in \partial B_1$. Since $y_n \in$ $\partial \Omega_u^+ \cap \partial \Omega_u^-$, Lemma [6.3](#page-13-4) implies that, for every $x \in B_{2r_n}(y_n) \cap \{u > 0\}$, we have

$$
|u(x) - \mu_+(y_n) \max\{0, (x - y_n) \cdot \nu_{y_n}\}| \le C|x - y_n|^{1+\gamma} \le Cr_n^{1+\gamma}.
$$

After rescaling, this gives

$$
|u_{y_n,r_n}(\xi+\xi_n)-\mu_+(y_n)\max\{0,(\xi+\xi_n)\cdot\nu_{y_n}\}|\leq Cr_n^{\gamma}\quad\text{for every}\quad\xi\in B_1\cap\{u_{x_n,r_n}>0\}.
$$

Moreover, by the continuity of μ_+ on $\partial \Omega_u^+ \cap \partial \Omega_u^-$, we have that, for every $\xi \in B_1$,

$$
\lim_{n \to \infty} \left| \mu_+(y_n) \max\{0, (\xi + \xi_n) \cdot \nu_{y_n}\} - \mu_+(y_0) \max\{0, (\xi + \xi_\infty) \cdot \nu_{y_0}\} \right| = 0.
$$

Therefore, it follows that $u_{x_n,r_n}(\xi) = u_{y_n,r_n}(\xi + \xi_n)$ converges to

$$
u_{\infty}(\xi) = \mu_+(y_0) \max \left\{ 0, (\xi + \xi_{\infty}) \cdot \nu_{y_0} \right\} \text{ for every } \xi \in B_1.
$$

Next we claim that $\xi_{\infty} \cdot \nu_{y_0} = 0$. Indeed, if $\xi_{\infty} \cdot \nu_{x_0} > 0$, then $u_{\infty}(0) > 0$ which is in contradiction with the uniform convergence of u_n ; on the other hand, if $\xi_\infty \cdot e_{x_0} < 0$, then $u_\infty \equiv 0$ in a neighborhood of zero, which is in contradiction with the non-degeneracy of u_n . Thus, we get

$$
u_{\infty}(\xi) = \mu_+(y_0) \max\left\{0, \xi \cdot \nu_{y_0}\right\} \text{ for every } \xi \in B_1.
$$

Now since $|\nabla u_{\infty}| = \mu_+(y_0)$, we get that $\mu_+(y_0) = Q_+^{1/2}(y_0)$.

Theorem [1.5](#page-2-0) is now a consequence of [\(6.3\)](#page-13-5), the Lemma [6.4](#page-14-1) and a general result (Theorem [A.1\)](#page-18-0) on the regularity of the one-phase flat free boundaries, which is due to De Silva (see [\[13\]](#page-20-11)). In the appendix we state Theorem [A.1](#page-18-0) in its full generality, for viscosity solutions of the problem (6.3) , but in our case the function u_+ is a classical solution, differentiable everywhere on $\overline{\Omega}_u^+$ $_{u}^{-}.$

7. Proof of Theorem [1.8](#page-2-1)

7.1. Preliminary results. In this subsection, we briefly recall the known results on the problem (1.4) . The existence of a solution of (1.4) in the class of the almost-open subsets of \mathcal{D} can be proved by a general variational argument (we refer to [\[7\]](#page-20-12) and to the book [\[5\]](#page-20-13) for more details). In the context of open sets, the existence of an optimal n-uple was proved in [\[3\]](#page-20-4).

From now on, $(\Omega_1, \ldots, \Omega_n)$ will be a solution of (1.4) and $u_i : \mathbb{R}^d \to \mathbb{R}$, for $i = 1, \ldots, n$, will denote the first normalized eigenfunction of the Dirichlet Laplacian on Ω_i , that is,

$$
-\Delta u_i = \lambda_1(\Omega_i)u_i \quad \text{in} \quad \Omega_i\,, \qquad u_i = 0 \quad \text{on} \quad \mathbb{R}^2 \setminus \Omega_i\,, \qquad \int_{\Omega_i} u_i^2 \, dx = 1,
$$

where, for every $i = 1, \ldots, n$,

$$
\lambda_1(\Omega_i) = \min_{u \in H_0^1(\Omega_i) \setminus \{0\}} \frac{\int_{\Omega_i} |\nabla u|^2 dx}{\int_{\Omega_i} u^2 dx} = \frac{\int_{\Omega_i} |\nabla u_i|^2 dx}{\int_{\Omega_i} u_i^2 dx}
$$

,

where $H_0^1(\Omega_i) = \{u \in H^1(\mathbb{R}^d) : u = 0 \text{ on } \mathbb{R}^2 \setminus \Omega_i\}.$ In particular, $u_i \ge 0$ on \mathbb{R}^2 and $\Omega_i = \{u_i > 0\}.$

Lipschitz continuity. The functions $u_i : \mathbb{R}^2 \to \mathbb{R}$ are Lipschitz continuous on \mathbb{R}^2 , that is, there is a universal constant $L > 0$ such that $\|\nabla u_i\|_{L^{\infty}(\mathbb{R}^2)} \leq L$, for every $i = 1, \ldots, n$. We refer to [\[7\]](#page-20-12) for the general case and to [\[3\]](#page-20-4) for a simplified version in dimension two.

Absence of triple points. For every $1 \leq i < j < k \leq n$, we have that $\partial \Omega_i \cap \partial \Omega_j \cap \partial \Omega_k = \emptyset$ (see [\[7\]](#page-20-12) and [\[3\]](#page-20-4) for a simpler proof in dimension two).

Absence of two-phase points on the boundary of the box. For every $1 \leq i \leq j \leq n$, we have that $\partial\Omega_i \cap \partial\Omega_j \cap \partial\mathcal{D} = \emptyset$ (see [\[3\]](#page-20-4)).

As a consequence of the above properties, we have that, for every $i \in \{1, \ldots, n\}$, the boundary $\partial\Omega_i$ can be decomposed as follows:

$$
\partial\Omega_i = \bigcup_{k\neq i} (\partial\Omega_i \cap \partial\Omega_k) \cup (\partial\Omega_i \cap \partial\mathcal{D}) \cup \Gamma_{\text{op}}(\Omega_i),
$$

where $\Gamma_{\text{op}}(\Omega_i)$ is the one-phase free boundary of Ω_i , determined by:

$$
x_0 \in \Gamma_{\text{op}}(\Omega_i) \Leftrightarrow
$$
 there exists $r > 0$ such that $B_r(x_0) \cap \left((\mathbb{R}^2 \setminus \mathcal{D}) \cup \bigcup_{k \neq i} \Omega_k \right) = \emptyset$.

We notice that already using the the regularity result of Briançon and Lamboley [\[4\]](#page-20-10), the onephase free boundary (lying inside the open set \mathcal{D}) is locally a $C^{1,\alpha}$ curve. Thus, in order to prove Theorem [1.8,](#page-2-1) it will be sufficient to show that $\partial\Omega_i$ is $C^{1,\alpha}$ in a neighborhood of the points of $\partial\Omega_i \cap \partial \mathcal{D}$ (Subsection [7.2\)](#page-16-0) and $\partial\Omega_i \cap \partial\Omega_k$ (Subsection [7.3\)](#page-18-1).

7.2. One-phase points at the boundary of the box. Let $1 \leq i \leq n$ and $x_0 \in \partial D \cap \partial \Omega_i$. Then, there is a neighborhood U of x_0 such that $U \cap \Omega_i = \emptyset$, for every $j \neq i$. For the sake of simplicity, in this subsection, we will set

$$
\Omega = \Omega_i
$$
, $u = u_i$, $x_0 = 0$ and $D = \Omega_i \cup (\mathcal{D} \cap \mathcal{U})$.

It is well known that the eigenvalues of the Dirichlet Laplacian are variationally characterized by

$$
\lambda_1(\Omega) = \int_{\Omega} |\nabla u|^2 dx = \min \left\{ \int_{\Omega} |\nabla v|^2 dx : v \in H_0^1(\Omega), \int_{\Omega} v^2 dx = 1 \right\}.
$$

Moreover, $\{u > 0\} = \Omega$ and u is a solution of the following minimization problem:

(7.1)
$$
\min \left\{ \int_D |\nabla v|^2 dx + \Lambda |\{v > 0\}| : v \in H_0^1(D), \int_D v^2 dx = 1 \right\}.
$$

We will show that the solution u of (7.1) is an almost-minimizer of the one-phase functional J_{op} . A result in the same spirit was proved given in a more general case in [\[16,](#page-20-5) Proposition 2.1].

Lemma 7.1 (Almost-minimality of the eigenfunction). Let $u : \mathbb{R}^d \to \mathbb{R}$ be a Lipschitz continuous function, $L = \|\nabla u\|_{L^{\infty}}$ be the Lipschitz constant of u and $\lambda_1(\Omega_u) = \int_D |\nabla u|^2 dx$. If u is a solution of the minimization problem [\(7.1\)](#page-16-1), then there exists $r_0 > 0$ such that u satisfies the following almost-minimality condition:

For every $r \in (0, r_0)$ and $x_0 \in \partial \Omega_u$,

$$
\int_{B_r(x_0)} |\nabla u|^2 dx + \Lambda |\Omega_u \cap B_r(x_0)| \le (1 + C_1 r^{d+2}) \int_{B_r(x_0)} |\nabla v|^2 dx + \Lambda |\Omega_v \cap B_r(x_0)| + C_2 r^{d+2},
$$

for every $v \in H_0^1(D)$ such that $u = v$ on $D \setminus B_r(x_0)$, where $C_1 = 2L^2$ and $C_2 = \lambda_1(\Omega_u)2L^2$.

Proof. Let $x_0 \in \partial \Omega_u$, $r > 0$ and $v \in H_0^1(D)$ be such that $u = v$ on $D \setminus B_r(x_0)$. Then, define the renormalization $w = ||v||_{L^2}^{-1} v \in H_0^1(D)$ and notice that we have

$$
\int_{D} |\nabla w|^2 dx = \left(\int_{D} v^2 dx \right)^{-1} \int_{D} |\nabla v|^2 dx \le \left(1 - \int_{B_r(x_0)} u^2 dx \right)^{-1} \int_{D} |\nabla v|^2 dx
$$

$$
\le \frac{1}{1 - L^2 r^{d+2}} \int_{D} |\nabla v|^2 dx \le \left(1 + 2L^2 r^{d+2} \right) \int_{D} |\nabla v|^2 dx,
$$

where for the last inequality, we choose r_0 such that $2L^2 r_0^{d+2} \leq 1$ and we use the inequality 1 $\frac{1}{1-X} \leq 1+2X$, for every $X \leq 1/2$, with $X = L^2r^{d+2}$. Now use w as a test function in [\(7.1\)](#page-16-1) to get that

(7.2)
$$
\int_{D} |\nabla u|^2 dx + \Lambda |\{u > 0\}| \le (1 + 2L^2 r^{d+2}) \int_{D} |\nabla v|^2 dx + \Lambda |\{v > 0\}|,
$$

from which the claim easily follows since \int $\int_{D\setminus B_r(x_0)}|\nabla v|^2\,dx=\int$ $\int_{D\setminus B_r(x_0)}|\nabla u|^2\,dx\leq \lambda_1(\Omega_u). \quad \Box$

We now notice that the C^2 regularity of $\partial \mathcal{D}$ implies that there is a constant $\delta > 0$ and a function $g: (-\delta, \delta) \to \mathbb{R}$ such that

$$
\mathcal{D}\cap SQ_{\delta}=\big\{(x_1,x_2)\in SQ_{\delta}\,:\,g(x_1)
$$

where $SQ_{\delta} = (-\delta, \delta) \times (-\delta, \delta)$. Moreover, up to a rotation of the plane, we can assume that $g'(0) = 0$. Let $\psi : SQ_{\delta} \subset \mathbb{R}^2 \to \mathbb{R}^2$ be the function that straightens out the boundary of D and let $\phi = \psi^{-1} : \psi(SQ_{\delta}) \subset \mathbb{R}^2 \to \mathbb{R}^2$ be its inverse:

$$
\psi(x_1, x_2) = (x_1, x_2 - g(x_1)), \quad \phi(x_1, x_2) = (x_1, x_2 + g(x_1)).
$$

We define the matrix-valued function $A = (a_{ij})_{ij} : SQ_\delta \to M^2(\mathbb{R})$ by

$$
A_x := \begin{pmatrix} a_{11}(x) & a_{12}(x) \\ a_{21}(x) & a_{22}(x) \end{pmatrix} = \begin{pmatrix} 1 & -g'(x_1) \\ -g'(x_1) & 1 + (g'(x_1))^2 \end{pmatrix} \text{ for every } x = (x_1, x_2) \in SQ_\delta.
$$

We recall that $H = \{(x_1, x_2) \in \mathbb{R}^2 : x_2 > 0\}$. By an elementary change of coordinates, we obtain the following result.

Lemma 7.2. Let u and A be as above. There exist constants $C_1, C_2 > 0$ and $r_0 > 0$ such that $B_{2r_0} \subset \psi(SQ_\delta)$ and the function $\tilde{u} := u \circ \phi$ satisfies the following almost-minimality condition:

For every $x_0 \in \partial \Omega_{\tilde{u}} \cap B_{r_0}$ and $r \in (0, r_0)$ we have

$$
\int_{B_r(x_0)} a_{ij}(x) \frac{\partial \tilde{u}}{\partial x_i} \frac{\partial \tilde{u}}{\partial x_j} dx + \Lambda |\Omega_{\tilde{u}} \cap B_r(x_0)|
$$
\n
$$
\leq (1 + C_1 r^{d+2}) \int_{B_r(x_0)} a_{ij}(x) \frac{\partial \tilde{v}}{\partial x_i} \frac{\partial \tilde{v}}{\partial x_j} dx + \Lambda |\Omega_{\tilde{v}} \cap B_r(x_0)| + C_2 r^{d+2},
$$

for every $\tilde{v} \in H^1(B_{2r_0})$ such that $\tilde{u} = \tilde{v}$ on $B_{2r_0} \setminus B_r(x_0)$ and $\Omega_{\tilde{v}} \subset H$.

Proof. Let $x_0 \in B_{r_0}$, $r \in (0, r_0)$ and \tilde{v} such that $\tilde{u} = \tilde{v}$ on $B_{2r_0} \setminus B_r(x_0)$. Then, use $v \in H_0^1(D)$ defined by $v = \tilde{v} \circ \psi$ in $\psi^{-1}(B_{2r_0})$ and $v = u$ otherwise, as a test function in Lemma [7.1](#page-16-2) to get

$$
\int_{B_{c_{\phi}r}(y_0)} |\nabla u|^2 \, dx + \Lambda |\Omega_u \cap B_{c_{\phi}r}(y_0)| \le (1 + C_1 r^{d+2}) \int_{B_{c_{\phi}r}(y_0)} |\nabla v|^2 \, dx + \Lambda |\Omega_v \cap B_{c_{\phi}r}(y_0)| + Cr^{d+2},
$$

where c_{ϕ} is a positive constant depending only on ϕ such that $\phi(B_r(x_0)) \subset B_{c_{\phi}r}(y_0)$ and $y_0 =$ $\phi(x_0)$. Now, with a change of coordinates and noticing that $u = v$ on $\phi(B_r(x_0))$ we have

$$
\int_{B_r(x_0)} a_{ij}(x) \frac{\partial \tilde{u}}{\partial x_i} \frac{\partial \tilde{u}}{\partial x_j} dx + \Lambda |\Omega_{\tilde{u}} \cap B_r(x_0)| = \int_{\phi(B_r(x_0))} |\nabla u|^2 dx + \Lambda |\Omega_u \cap \phi(B_r(x_0))|
$$

\n
$$
\leq (1 + C_1 r^{d+2}) \int_{\phi(B_r(x_0))} |\nabla v|^2 dx + \Lambda |\Omega_v \cap \phi(B_r(x_0))| + C_2 r^{d+2}
$$

\n
$$
= (1 + C_1 r^{d+2}) \int_{B_r(x_0)} a_{ij}(x) \frac{\partial \tilde{v}}{\partial x_i} \frac{\partial \tilde{v}}{\partial x_j} dx + \Lambda |\Omega_{\tilde{v}} \cap B_r(x_0)| + C_2 r^{d+2},
$$

where $C_2 = \lambda_1(\Omega_u)C_1 + C$. This concludes the proof.

Proof of Theorem [1.8](#page-2-1) (the one-phase boundary points). We are now in position to conclude the regularity of the free boundary $\partial\Omega_i$ in a neighborhood of any one-phase boundary point $x_0 \in$ $\partial\Omega_i \cap \partial\mathcal{D}$. Indeed, we may assume that $x_0 = 0$ and that $\partial\mathcal{D}$ is the graph of a function g. Reasoning as above, we have that $\tilde{u}_i(x_1, x_2) = u_i(x_1, x_2 + g(x_1))$ satisfies the almost-minimality condition from Lemma [7.2](#page-17-0) in a neighborhood of the origin. On the other hand, it is immediate to check that \tilde{u}_i is still Lipschitz continuous. Thus, we can apply Theorem [1.1](#page-1-0) obtaining that, in a neighborhood of zero, $\partial\Omega_i$ is the graph of a $C^{1,\alpha}$ function.

$$
\Box
$$

7.3. Two-phase points. Let Ω_i and Ω_j be two different sets from the optimal n-uple $(\Omega_1, \ldots, \Omega_n)$, solution of [\(1.4\)](#page-2-2). Let u_i and u_j be the first normalized eigenfunctions, respectively on Ω_i and Ω_j . Finally, let $x_0 \in \partial \Omega_i \cap \partial \Omega_j$. We know that there is a neighborhood $\mathcal{U} \subset \mathcal{D}$ of x_0 such that $U \cap \Omega_k = \emptyset$, for every $k \notin \{i, j\}$. Setting $D := \Omega_i \cup \Omega_j \cup \mathcal{U}$, we get that the function $u := u_i - u_j$ is the solution of the two-phase problem

(7.3)
$$
\min\Big\{\int_D|\nabla v|^2\,dx+q_i|\Omega_v^+|+q_j|\Omega_v^-| \;:\; v\in H_0^1(D),\;\int_D v_+^2\,dx=\int_D v_-^2\,dx=1\Big\}.
$$

We next show that the solutions of [\(7.3\)](#page-18-2) satisfy a almost-minimality condition.

Lemma 7.3. Let $D \subset \mathbb{R}^d$, $u \in H_0^1(D)$ be a Lipschitz continuous function on \mathbb{R}^d and L its Lipschitz constant. Suppose that u is a solution of the minimization problem (7.3) . Then, there is some $r_0 > 0$ such that u satisfies the following almost-minimality condition: For every $r \in (0, r_0)$ and $x_0 \in \partial \Omega_u$,

$$
\int_{B_r(x_0)} |\nabla u|^2 dx + q_i |\Omega_u^+ \cap B_r(x_0)| + q_j |\Omega_u^- \cap B_r(x_0)|
$$
\n
$$
\leq (1 + C_1 r^{d+2}) \int_{B_r(x_0)} |\nabla v|^2 dx + q_i |\Omega_v^+ \cap B_r(x_0)| + q_j |\Omega_v^- \cap B_r(x_0)| + C_2 r^{d+2},
$$

for every $v \in H_0^1(D)$ such that $u - v \in H_0^1(B_r(x_0))$, where $C_1 = 2L^2$ and $C_2 = C_1$ $\int_D |\nabla u|^2 dx.$

Proof. Follows precisely as in Lemma [7.1.](#page-16-2)

We are now in position to complete the proof of Theorem [1.8.](#page-2-1)

Proof of Theorem [1.8](#page-2-1) (the two-phase free boundary). We only need to notice that in a neighborhood of any two-phase point $x_0 \cap \partial \Omega_i \cap \partial \Omega_j \cap \mathcal{D}$, Lemma [7.3](#page-18-3) implies that u is a almost-minimizer of J_{TP} , where the matrix A is the identity, $Q_+ = q_i$ and $Q_- = q_i$. Thus, it is sufficient to apply Theorem [1.5.](#page-2-0) \Box

APPENDIX A. THE FLAT ONE-PHASE FREE BOUNDARIES ARE $C^{1,\alpha}$

In this section we discuss a regularity theorem for viscosity solutions of the one-phase problem (without constraint). We fix $f : B_2 \to \mathbb{R}$ to be a bounded and continuous function and $A : B_2 \to$ Sym_d^+ to be a matrix-valued coercive and bounded function with Hölder continuous coefficients, as in the Introduction. Before we state the result, we recall that the continuous function $u : \mathbb{R}^d \supset$ $B_1 \to \mathbb{R}, u \geq 0$ is a viscosity solution of

(A.1)
$$
-\text{div}(A\nabla u) = f \text{ in } \Omega_u \cap B_1, \qquad |A^{1/2}[\nabla u]| = g \text{ on } \partial \Omega_u \cap B_1,
$$

if the first equation holds in the open set Ω_u and if, for every $x_0 \in \partial \Omega_u$ and every $\varphi \in C^{\infty}(\mathbb{R}^d)$ touching $u \circ F_{x_0}$ from above (below) at zero, we have that $|\nabla \varphi|(0) \ge g(x_0)$ (resp. $|\nabla \varphi|(0) \le g(x_0)$). Recall that touching from above (below) means that $\varphi(0) = 0$ and $\varphi \geq u \circ F_{x_0}$ (resp. $\varphi \leq u \circ F_{x_0}$) in $\Omega_u \cap B_1$. Moreover, we suppose that g is Hölder continuous and that there are constants $\eta_g > 0$, $C_q > 0$ and $\delta_g > 0$ such that

(A.2)
$$
\begin{cases} |g(x) - g(y)| \le C_g |x - y|^{\delta_g} & \text{for every} \quad x, y \in \partial \Omega_u \cap B_1, \\ \eta_g \le g(x) & \text{for every} \quad x \in \partial \Omega_u \cap B_1. \end{cases}
$$

The following result follows immediately from the results proved in [\[13\]](#page-20-11).

Theorem A.1 (Flat free boundaries are $C^{1,\alpha}$). Suppose that $u : B_1 \to \mathbb{R}$ is a viscosity solution of [\(A.1\)](#page-18-4) and that $g: \partial \Omega_u \to \mathbb{R}$ satisfies [\(A.2\)](#page-18-5). Then, there exist $\varepsilon > 0$ and $\rho > 0$ such that if $x_0 \in \partial \Omega_u \cap B_1$ and u is such that

 $g(x_0) \max\{0, x \cdot \nu - \varepsilon \rho\} \le u \circ F_{x_0}(x) \le g(x_0) \max\{0, x \cdot \nu + \varepsilon \rho\}$ for every $x \in B_\rho$, then $\partial \Omega_u$ is $C^{1,\alpha}$ in $B_{\rho/2}(x_0)$.

$$
\Box
$$

Remark A.2. Notice that since in dimension two all the blow-up limits of u_+ (given by Theorem [1.5\)](#page-2-0) are half-plane solutions (Proposition [4.3\)](#page-9-1), we have that the flatness assumption of the above Theorem is satisfied at every point of the free boundary $\partial\Omega_u^+$. We also notice that, in our case, we have $g = \mu_{+}$, which is Hölder continuous by Lemma [6.4.](#page-14-1)

Definition A.3 (Flatness). Let $u : B_1 \to \mathbb{R}$ be continuous, $u \geq 0$ and $u \in H^1(B_1)$. We say that u is (ε, ν) -flat, if there are a matrix-valued $A : B_1 \to Sym_d^+$ with Hölder continuous coefficients, and a continuous $f : B_1 \to \mathbb{R}$ such that:

$$
(A.3) \t\t -div (A\nabla u) = f \t in \t \Omega_u \cap B_1;
$$

(A.4)
$$
||f||_{L^{\infty}(B_1)} \leq \varepsilon \quad and \quad ||a_{ij} - \delta_{ij}||_{L^{\infty}(B_1)} \leq \varepsilon^2 \quad for \ every \quad 1 \leq i, j \leq d;
$$

(A.5)
$$
1 - \varepsilon^2 \leq |\nabla u| \leq 1 + \varepsilon^2 \quad on \quad \partial \Omega_u \cap B_1;
$$

(A.6)
$$
\max\{0, x \cdot \nu - \varepsilon\} \le u(x) \le \max\{0, x \cdot \nu + \varepsilon\} \text{ for every } x \in B_1.
$$

Remark A.4. The condition [\(A.5\)](#page-19-0) is intended in a viscosity sense, that is, for any $\varphi \in C^{\infty}(B_1)$, we have:

- if
$$
\varphi(x_0) = u(x_0)
$$
 for some $x_0 \in \partial\Omega_u \cap B_1$ and $\varphi^+ \ge u$ in $\Omega_u \cap B_1$, then $|\nabla \varphi(x_0)| \ge 1 - \varepsilon^2$;
- if $\varphi(x_0) = u(x_0)$ for some $x_0 \in \partial\Omega_u \cap B_1$ and $\varphi \le u$ in $\Omega_u \cap B_1$, then $|\nabla \varphi(x_0)| \le 1 + \varepsilon^2$.

In order to prove Theorem [A.1](#page-18-0) one has to show that the flatness improves at lower scales, that is, if u is (ε, ν) -flat, then a rescaling u_r of u is $(\varepsilon/2, \nu')$ -flat for some ν' , which is close to ν . Of course, the essential (and hardest) part of the proof is to show the improvement of the geometric flatness [\(A.6\)](#page-19-1). This was proved by De Silva in [\[13,](#page-20-11) Lemma 4.1].

Lemma A.5 (Improvement of the geometric flatness). There are universal constants $C > 0$, $r_0 > 0$ and $\varepsilon_0 > 0$ such that if u is ε -flat in the direction ν , for some $\varepsilon \in (0, \varepsilon_0)$ and $\nu \in \partial B_1$, then, for every $r \in (0, r_0)$ there is some $\nu' \in \partial B_1$ such that $|\nu - \nu'| \leq C \varepsilon^2$ and

$$
\max\left\{0, x \cdot \nu' - \frac{\varepsilon}{2}\right\} \le u_r(x) \le \max\left\{0, x \cdot \nu' + \frac{\varepsilon}{2}\right\} \quad \textit{for every} \quad x \in B_1,
$$

where $u_r : B_1 \to \mathbb{R}$ is the one-homogeneous rescaling $u_r(x) = \frac{u(rx)}{r}$.

Proof of Theorem [A.1.](#page-18-0) We will first prove that the flatness condition $(A.3)-(A.6)$ $(A.3)-(A.6)$ improves at smaller scales. We fix $x_0 \in \partial \Omega_u \cap B_1$ and we consider the function $\tilde{u} = \frac{1}{g(x)}$ $\frac{1}{g(x_0)}u \circ F_{x_0}$ (recall that $F_{x_0}(x) = x_0 + A_{x_0}^{1/2}[x]$. Let ε and r_0 be the constants from Lemma [A.5.](#page-19-3) We will prove that there is $r_1 \leq r_0$ such that: if \tilde{u} is (ε, ν) -flat, then for every $r \leq r_1$, \tilde{u}_r is $(\varepsilon/2, \nu')$ -flat, for ν' given again by Lemma [A.5.](#page-19-3) It is sufficient that the conditions [\(A.3\)](#page-19-2), [\(A.4\)](#page-19-4) and [\(A.5\)](#page-19-0) are satisfied for \tilde{u}_r with the flatness parameter $\varepsilon/2$. We notice that \tilde{u} is a viscosity solution of

(A.7)
$$
-\text{div}(\tilde{A}\nabla\tilde{u}) = \tilde{f} \quad \text{in} \quad \Omega_{\tilde{u}}, \qquad |\tilde{A}^{1/2}[\nabla\tilde{u}]| = \tilde{g} \quad \text{on} \quad \partial\Omega_{\tilde{u}},
$$

where $\tilde{A}_x = A_{x_0}^{-1/2} A_{F_{x_0}(x)} A_{x_0}^{-1/2}, \ \tilde{f} = \frac{1}{g(x_0)}$ $\frac{1}{g(x_0)} f \circ F_{x_0}, \; \tilde{g} \, = \, \frac{1}{g(x_0)}$ $\frac{1}{g(x_0)}g \circ F_{x_0}$ and $\tilde{A}_x^{1/2} = A_{F_x}^{1/2}$ $\frac{1}{2}F_{x_0}(x) \circ A_{x_0}^{-1/2}.$ Notice that $0 \in \partial \Omega_{\tilde{u}}$ and set $\tilde{u}_r(x) := \frac{\tilde{u}(rx)}{r}$. Thus, for small enough $r > 0$, \tilde{u}_r is a viscosity solution of

(A.8)
$$
-\text{div}(\tilde{A}_r \nabla \tilde{u}_r) = \tilde{f}_r \quad \text{in} \quad \Omega_{\tilde{u}} \cap B_1, \qquad \left| \tilde{A}_r^{1/2} [\nabla \tilde{u}_r] \right| = \tilde{g}_r \quad \text{on} \quad \partial \Omega_{\tilde{u}} \cap B_1,
$$

where $\tilde{A}_r(x) := \tilde{A}(rx)$, $\tilde{f}_r(x) = r\tilde{f}(rx)$, $\tilde{g}_r(x) = \tilde{g}(rx)$ and $\tilde{A}_r^{1/2}(x) = A_{F_r}^{1/2}$ $_{F_{x_0}(rx)}^{1/2} \circ A_{x_0}^{-1/2}$. Now, if \tilde{u} is (ε, ν) -flat, then the Hölder continuity of the coefficients a_{ij} and the boundedness of f imply that [\(A.4\)](#page-19-4) holds with $\epsilon/2$ and \tilde{u}_r , for any $r \leq r_1$, where $r_1 \leq r_0$, is a universal constant depending on the Hölder norm of a_{ij} . Now, in order to get [\(A.5\)](#page-19-0) for $\varepsilon/2$ and \tilde{u}_r , we suppose that $\varphi \in C^{\infty}(B_1)$ touches \tilde{u}_r from below at a point $y_0 \in B_1 \cap \partial {\tilde{u}_r} > 0$. Thus, we have that

$$
|A_{F_{x_0}(ry_0)}^{1/2} \circ A_{x_0}^{-1/2}[\nabla \varphi(y_0)]| \leq \frac{g(F_{x_0}(ry_0))}{g(x_0)},
$$

and so, if $\|\cdot\| = \|\cdot\|_{\mathcal{L}(\mathbb{R}^d)}$ stands for the space of $d \times d$ matrices, we have

$$
|\nabla \varphi(y_0)| \leq \|A_{x_0}^{1/2} \circ A_{F_{x_0}(ry_0)}^{-1/2} \| \frac{g(F_{x_0}(ry_0))}{g(F_{x_0}(0))}.
$$

Now, by the Hölder continuity (and the uniform boundedness from below) of g, we can choose r_1 such that

$$
\frac{g(F_{x_0}(ry_0))}{g(F_{x_0}(0))} \le 1 + \frac{\varepsilon^2}{10}.
$$

On the other hand, there are universal constants C and $\delta > 0$, depending only on the Hölder exponent δ_A and the norm C_A , of the matrix-valued function A, such that

$$
||A_{x_0}^{1/2} \circ A_{F_{x_0}(ry_0)}^{-1/2} - Id|| \leq ||A_{x_0}^{1/2} - A_{F_{x_0}(ry_0)}^{1/2}|| ||A_{F_{x_0}(ry_0)}^{-1/2}|| \leq C |ry_0|^{\delta} \leq Cr_1^{\delta}.
$$

Choosing r_1 such that $Cr_1^{\delta} \leq \frac{\varepsilon^2}{10}$ and using the triangular inequality, we get

$$
|\nabla \varphi(y_0)| \leq \|A_{x_0}^{1/2} \circ A_{F_{x_0}(ry_0)}^{-1/2} \|\frac{g(F_{x_0}(ry_0))}{g(F_{x_0}(0))} \leq \left(1 + \frac{\varepsilon^2}{10}\right)^2 \leq 1 + (\varepsilon/2)^2,
$$

which completes the proof of the improvement of flatness for \tilde{u} , the case when φ touches from above being analogous. Now, the claim follows by a standard argument, similar to the one we used in Section [5.](#page-11-0)

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